

As we were saying last time an inner product on a vector space  $V$  induces a natural isomorphism  $V \cong V^*$ . Let  $\langle \cdot, v \rangle$  denote the element of  $V^*$  defined by  $x \mapsto \langle x, v \rangle$ . Let

$$\# : v \mapsto \langle \cdot, v \rangle, \quad \# : V \rightarrow V^*$$

so that the nondegeneracy of  $\langle \cdot, \cdot \rangle$  is equivalent to  $\ker \# = \{0\}$ .

$$V \begin{array}{c} \xrightarrow{\#} \\ \xleftarrow{\#} \end{array} V^* \quad \text{"natural isomorphism"} \\ \text{(finite dim. only.)}$$

These then help to define the "adjoint" of operator  $L : V \rightarrow V$ .

If  $L^+ : V \rightarrow V$  satisfies

$$\langle Lv, w \rangle = \langle v, L^+w \rangle \quad v, w \in V$$

$L^+$  is called an adjoint of  $L$ . Adjoints are unique in finite dimensional spaces because of the following. First note that

$(v \mapsto \langle Lv, w \rangle) \in V^*$  and must therefore be equal to  $\#(x_w)$

for some  $x_w \in V$ .  $\Rightarrow \#(x_w)(v) = \cancel{\langle Lv, w \rangle} \cancel{\langle v, x_w \rangle} = \langle Lv, w \rangle$

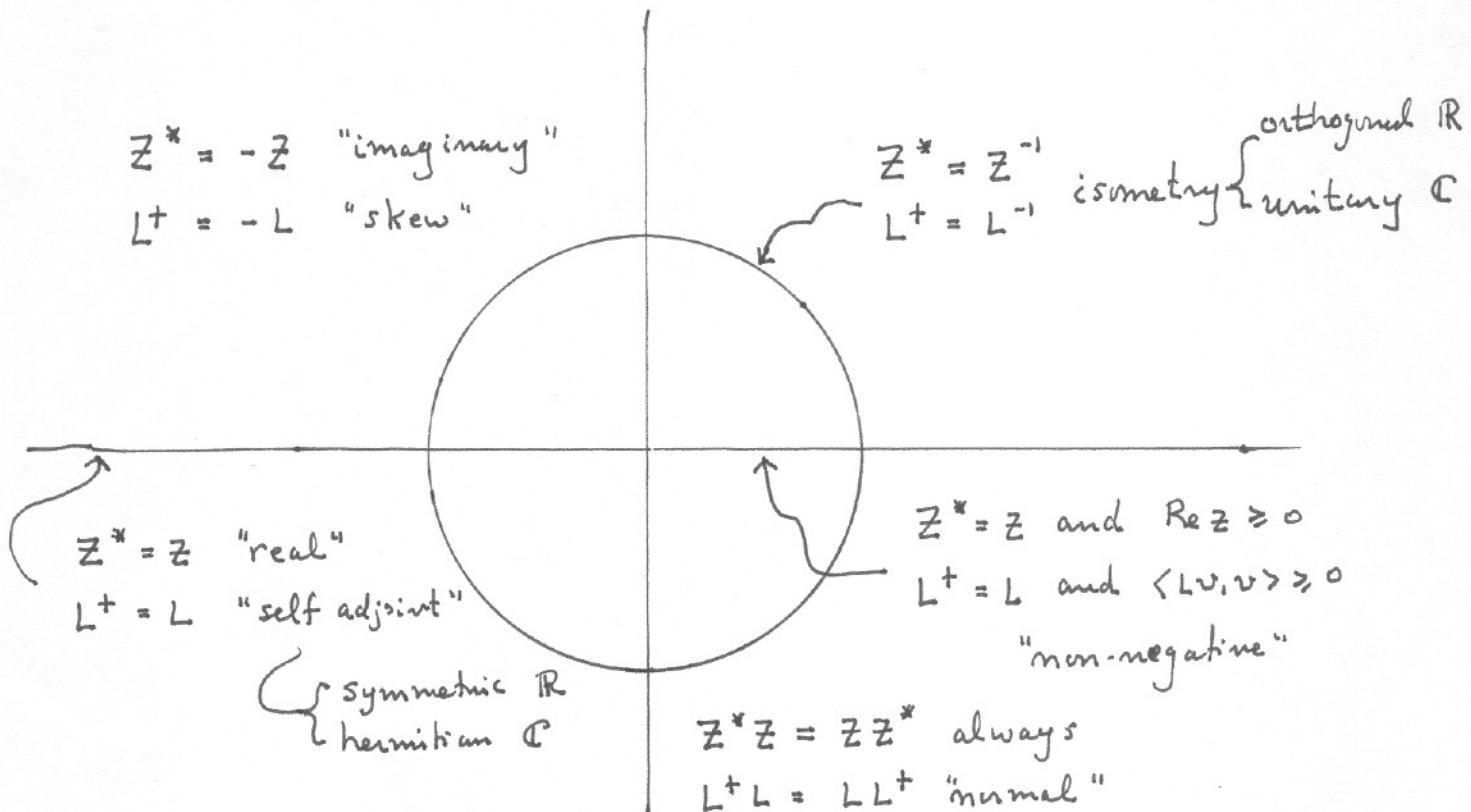
for all  $v, w \in V$ .  $\Rightarrow L^+w = x_w$  is the unique adjoint of  $L$ .

Adjoints have a number of nice and easily proved properties

$$(L^+)^+ = L \quad (\text{i.e. } + \text{ is an } \underline{\text{involution}})$$

$$1^+ = 1 \quad (LM)^+ = M^+L^+$$

$L$  has an inverse iff  $L^+$  has an inverse ...



Complex spectral theorem:  $V$  has an orthonormal basis of eigenvectors iff  $L$  is normal.

Real spectral theorem:  $V$  has an orthonormal basis of eigenvectors iff  $L$  is self adjoint.

ex:

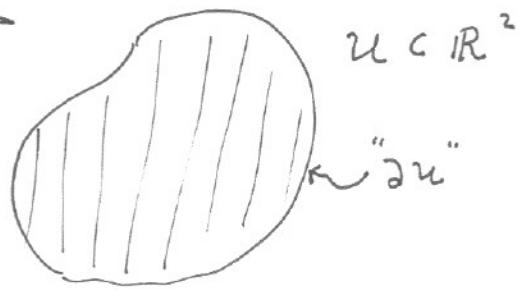
Theorem: Any operator  $L$  can be written

$$L = S(L^+L)^{1/2}$$

for some isometry  $S$ .

See Axler for proof.

## Application: The Numerical Laplacian



$\Delta \varphi = 0$  Laplace eq.

$\Delta \varphi = f$  Poisson eq.

$\partial_t \varphi = \Delta \varphi$  Diffusion eq.

$\partial_t^2 \varphi = c^2 \Delta \varphi$  Wave eq.

$i\hbar \partial_t \psi = -\frac{\hbar^2}{2m} \Delta \psi + V \psi$  Schrödinger eq.

$$\Delta = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2}$$

First let's proceed in the continuous case in hand waving fashion..

This can be done precisely but only later in the course.

We think of  $\Delta$  as an operator on some vector space  $\mathcal{C}(U)$  of smooth functions  $U \rightarrow \mathbb{R}$ , with inner product

$$\langle f, g \rangle = \int_U f(x) g(x) dx$$

We'll just assume that  $U$  and  $\mathcal{C}(U)$  are chosen so that this is true.

We'll also assume (Green)

$$\langle \Delta f, g \rangle = \int_U (\Delta f) g = - \int_U (\nabla f) \cdot (\nabla g) + \int_{\partial U} (\hat{n} \cdot \nabla f) g$$

If  $\partial u = \emptyset$ ,  $\Delta$  is self adjoint, so that if  $\varphi$  is a solution

to Laplace,  $\langle \Delta \varphi, \varphi \rangle = - \langle \nabla \varphi, \nabla \varphi \rangle = 0 \Rightarrow \nabla \varphi = 0 \Rightarrow \varphi = \text{const.}$  are the only solutions.

e.g. Suppose  $\Delta g = \rho$  and  $f$  is a solution to Laplace,

$$\Rightarrow \langle \Delta g, f \rangle = \langle \rho, f \rangle = \langle \rho, \Delta f \rangle = 0 \Rightarrow \int_u \rho(x) = 0.$$

e.g. Eigenvalues of  $\Delta$  are  $\leq 0$ . Suppose  $\Delta g = 2g$ ,  $g \neq 0$ .

$$\langle \Delta g, g \rangle = 2 \langle g, g \rangle = -\langle \nabla g, \nabla g \rangle \Rightarrow 2 \leq 0.$$

e.g.  $\Delta g = \lambda g$ ,  $\Delta g' = \lambda' g'$  eigenvectors, then

$$\langle \Delta g, g' \rangle = \lambda \langle g, g' \rangle = \lambda' \langle g, g' \rangle$$

$$\Rightarrow \lambda \neq \lambda' \Rightarrow \langle g, g' \rangle = 0.$$

Dropping the boundaryless assumption, we can get variational results. Suppose that we want  $g \in C(u)$  to satisfy  $g(x) = F(x)$  on  $x \in \partial u$  boundary points. Let

$$A = \{g \in C(u) : g(x) = F(x) \text{ for all } x \in \partial u\}$$

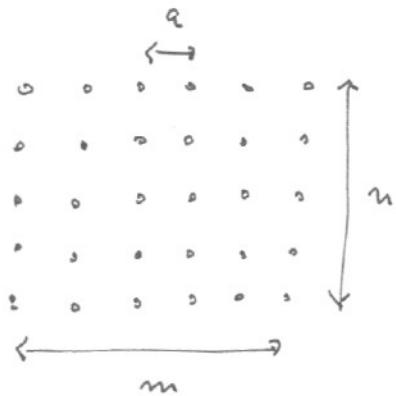
We claim that if  $g \in A$  minimizes  $E(g) = \langle \nabla g, \nabla g \rangle$ , then it is a solution to Laplace with the desired boundary conditions. To see this, let

$$A_0 = \{g \in C(u) : g(x) = 0 \text{ for all } x \in \partial u\}$$

Then if  $f \in A_0$ ,  $g + a \cdot f \in A$ .

$$E(g + af) = E(g) + a^2 E(f) + 2a \langle \nabla g, \nabla f \rangle \text{ for all } a \in \mathbb{R}$$

$$\Rightarrow \langle \nabla g, \nabla f \rangle = 0 \Rightarrow \langle \Delta g, f \rangle = 0 \text{ for any } f. \text{ "Weak solution"}$$



$$U = \{0, 1, \dots, m-1\} \times \{0, 1, \dots, n-1\}$$

$C(U)$  is some  $\infty$  dim.  
real vector space which  
we have not precisely  
defined.

$C(U)$  is precisely the  
free real vector space  
on the set  $U$ .

$$\langle f, g \rangle \equiv \int_U f(x) g(x) dx$$

We assumed that this  
was an inner product.

$$\langle f, g \rangle = \sum_u f_u \cdot g_u$$

This is an inner product.

$$\Delta = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2}$$

$$(\Delta_h f)_{ij} =$$

$$\{f_{i+1,j} + f_{i-1,j} + f_{i,j+1} + f_{i,j-1} - 4f_{ij}\}/a^2$$

Exact Laplacian

Numerical Laplacian correct to  $\mathcal{O}(a^2)$

Lets look at the numerical Laplacian in more detail.

$$(\Delta_L f)_{ij} = \frac{1}{a^2} (f_{i+1,j} + f_{i-1,j} + f_{i,j+1} + f_{i,j-1} - 4f_{ij})$$

can be written in terms of "shift" operators

$$(C_x f)_{ij} = f_{i+1,j} \quad (C_x^+ f)_{ij} = f_{i-1,j}$$

$$(C_y f)_{ij} = f_{i,j+1} \quad (C_y^+ f)_{ij} = f_{i,j-1}$$

Obviously,  $\langle C_x f, C_x f \rangle = \langle f, f \rangle$ , i.e.  $C_x, C_x^+, C_y, C_y^+$  are isometries.

$$\Rightarrow \Delta_L = \frac{1}{a^2} (C_x + C_x^+ + C_y + C_y^+ - 4I)$$

is self adjoint.

We can also define  $D_x = (C_x - I)/a$ ,  $D_y = (C_y - I)/a$

so that

$$-\Delta_L = D_x D_x^+ + D_y D_y^+$$

Now, the whole set of results that we suggested but did not prove in the continuous case follows precisely in the lattice case.

Continuous

$$\mathcal{U} \subset \mathbb{R}^2$$

$C(\mathcal{U})$  is an  $\infty$ -dim.  
real vector space

$$\langle f, g \rangle = \int_{\mathcal{U}} f(x)g(x) dx$$

$$\Delta = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2}$$

$$\left\{ \begin{array}{l} \langle \Delta f, g \rangle = \langle f, \Delta g \rangle \\ \langle \Delta f, g \rangle = -\langle \nabla f, \nabla g \rangle \\ \Delta y = 0 \Leftrightarrow \nabla y = 0 \end{array} \right.$$

$$\Delta y = p \Rightarrow \langle p, 1 \rangle = 0$$

$$E(y) = \langle \nabla y, \nabla y \rangle$$

Lattice

$$\mathcal{U} = \{0, 1, 2, \dots, m-1\} \times \{0, 1, \dots, n-1\}$$

$C(\mathcal{U})$  is the  $m \times n$  dim. free real  
vector space on  $\mathcal{U}$

$$\langle f, g \rangle = \sum_{u \in \mathcal{U}} f_u \cdot g_u$$

$$-\Delta_L = D_x D_x^+ + D_y D_y^+$$

$$\langle \Delta_L f, g \rangle = \langle f, \Delta_L g \rangle$$

$$\langle \Delta_L f, g \rangle = -\langle D_x f, D_x g \rangle - \langle D_y f, D_y g \rangle$$

$$\Delta_L y = 0 \Leftrightarrow D_x y = D_y y = 0$$

$$\Delta_L y = p \Rightarrow \langle p, 1 \rangle = 0$$

$$E_L(y) = \langle D_x y, D_x y \rangle + \langle D_y y, D_y y \rangle$$

Notice that if we had a basis of eigenvectors of  $D_x, D_y$ , this would also be a basis of  $D_x, D_y, D_x^+, D_y^+$  and  $\Delta$ .

Let's find this, but first, switch to  $\mathcal{U} = \{0, 1, 2, \dots, n-1\}$   
 $n = \text{even}$  for convenience ...

Since shift is an isometry the only possible eigenvalues are  $\pm 1$  (because  $\langle c\mathbf{y}, c\mathbf{y} \rangle = \langle \mathbf{y}, \mathbf{y} \rangle = \lambda^2 \langle \mathbf{y}, \mathbf{y} \rangle$ ), and all eigenvectors are proportional to

$$C \begin{pmatrix} 1 \\ 1 \\ \vdots \\ 1 \end{pmatrix} = 1 \cdot \begin{pmatrix} 1 \\ 1 \\ \vdots \\ 1 \end{pmatrix} \text{ or } C \begin{pmatrix} 1 \\ -1 \\ 1 \\ -1 \\ \vdots \\ -1 \end{pmatrix} = -1 \cdot \begin{pmatrix} 1 \\ -1 \\ 1 \\ -1 \\ \vdots \\ -1 \end{pmatrix}$$

and these two certainly don't span  $C(\mathcal{U})$  in general.

 But don't give up! We can try solving these problems in the free complex vector space on  $\mathcal{U}$  instead. Define  $C$  as a shift as before,  $D = \frac{1}{a}(C - I)$ ,  $\Delta = -DD^*$  as before.

$\langle f, g \rangle = \sum_{u \in \mathcal{U}} f_u g_u^*$  is the new inner product.

$C$  is still an isometry  $\Rightarrow \lambda = e^{i\theta_a}$ . Since  $C^n = 1$ , obviously  $\lambda^n = 1$  and the only possible eigenvalues are

$$\lambda \in \{z^0, z^1, z^2, \dots, z^{n-1}\} \quad z = e^{2\pi i / n}$$

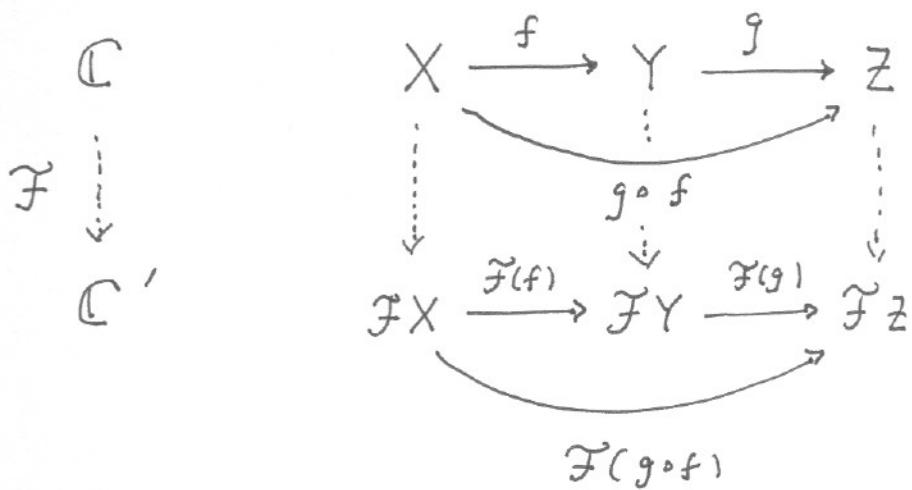
$$C \begin{pmatrix} z^0 \\ \vdots \\ z^0 \\ z^0 \\ z^0 \end{pmatrix} = z^0 \begin{pmatrix} z^0 \\ \vdots \\ z^0 \\ z^0 \\ z^0 \end{pmatrix} \quad C \begin{pmatrix} z^{n-1} \\ \vdots \\ z^2 \\ z^1 \\ z^0 \end{pmatrix} = z \begin{pmatrix} z^{n-1} \\ \vdots \\ z^2 \\ z^1 \\ z^0 \end{pmatrix} \quad C \begin{pmatrix} z^{2n-1} \\ \vdots \\ z^4 \\ z^2 \\ z^0 \end{pmatrix} = z^2 \begin{pmatrix} z^{2n-1} \\ \vdots \\ z^4 \\ z^2 \\ z^0 \end{pmatrix} \dots \quad C \begin{pmatrix} z^{(n-1)(n-1)-1} \\ \vdots \\ z^{n-1} \\ z^0 \end{pmatrix} = z^{n-1} \begin{pmatrix} z^{(n-1)(n-1)-1} \\ \vdots \\ z^{n-1} \\ z^0 \end{pmatrix}$$

$\Rightarrow$  we found eigenvectors for all  $n$  eigenvalues  $\Rightarrow$  this is a basis making the operation of  $C, D, \Delta$  trivial.

Changing to this basis is called the "Furier Transform."

[There is a very nice  $O(n \log n)$  FFT algorithm for this].

A functor from category  $\mathcal{C}$  to category  $\mathcal{C}'$  is a collection of mappings which preserves identities and commuting triangles:



For example:

$$\begin{array}{ccc}
 \text{Sets} & X & \xrightarrow{f} Y \xrightarrow{g} Z \\
 \downarrow & \downarrow & \downarrow \\
 \wp & \wp X & \xrightarrow{\wp f} \wp Y \xrightarrow{\wp g} \wp Z
 \end{array}$$

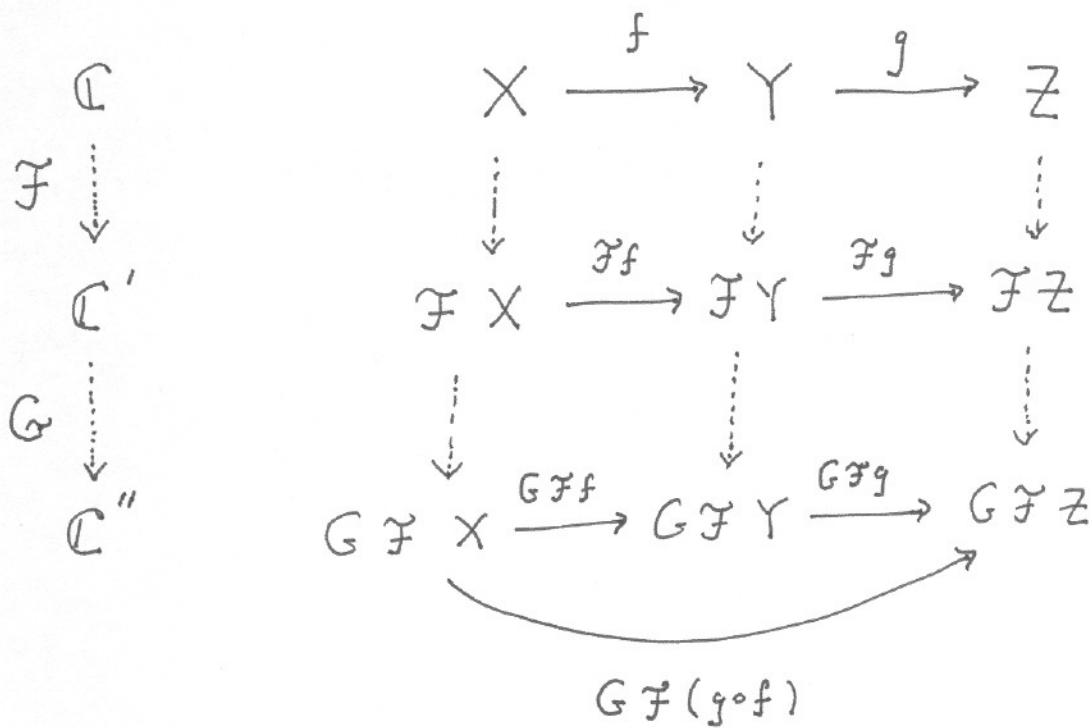
$$\wp X = \{ \text{the set of subsets of } X \}$$

$$\wp f: A \mapsto \{ f(a) : a \in A \} \in \wp Y$$

$$\wp(g \circ f) = \wp(g) \circ \wp(f)$$

$$\wp(i_X) = i_{\wp X}$$

This is a functor.



Composition of functors is defined by composing all the corresponding mappings.

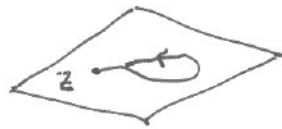
- Composition of functors is associative.
- For each category  $C$ , there is an identity functor (where all the mappings are identities) which obeys  $i_C \circ F = F$ ,  $F \circ i_{C'} = F$ .
- The categories of all categories is a category.

You can also easily verify that

Functors preserve isomorphisms.

Pointed

topological  
spaces

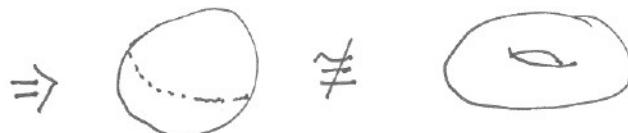


$$(X, x) \xrightarrow{f} (Y, y) \xrightarrow{g} (Z, z)$$

$\pi$   
Groups

$$\begin{array}{ccccc} & \downarrow & & \downarrow & \\ \pi(X, x) & \xrightarrow{\pi f} & \pi(Y, y) & \xrightarrow{\pi g} & \pi(Z, z) \\ \downarrow & & \downarrow & & \downarrow \\ \pi(g \circ f) & & & & \end{array}$$

$$\pi_f: [e] \mapsto [f \circ e]$$



"Homotopy"

Free constructions are all functors

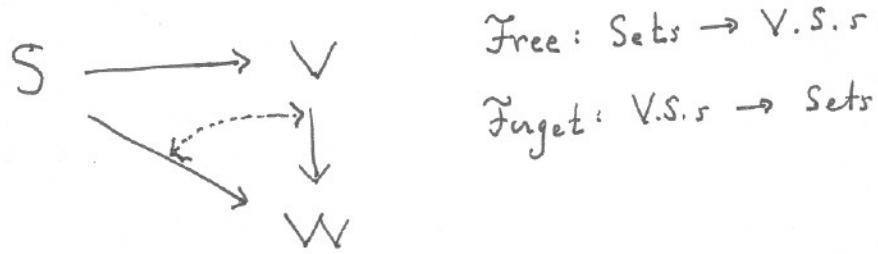
$$\begin{array}{ccc} S & \longrightarrow & \mathcal{F}(S) \\ f \downarrow & & \downarrow \mathcal{F}(f) \\ S' & \longrightarrow & \mathcal{F}(S') \end{array}$$

free vector space  
on  $S$

$$S \cong S' \Rightarrow \mathcal{F}(S) \cong \mathcal{F}(S')$$

Vector spaces with isomorphic bases are isomorphic.

Free constructions actually involve two functors



$$\text{Mor}_{\substack{\text{Vector Spaces} \\ (\text{Free } S, W)}} \cong \text{Mor}_{\text{Set}}(S, \text{Forget } W)$$

These are called "adjoint" functors.

## Associative Algebra

An associative algebra is just a vector space  $V$  with a bilinear product  $V \times V \rightarrow V$ .

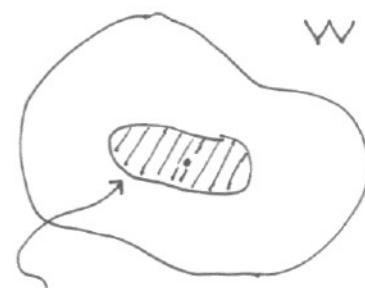
examples:

- Real valued functions on a set with pointwise multiplication.
- $\mathbb{R}^3$  with the usual cross product.
- Operators on a vector space with composition as multiplication.

Morphisms: Linear maps  $V \xrightarrow{\varphi} W$  with  $\varphi(vv') = \varphi(v)\varphi(v')$ .



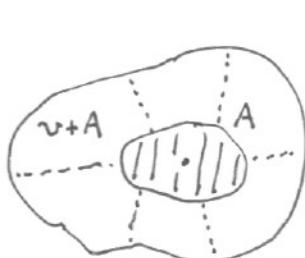
$$\text{Ker } \varphi = \{v : \varphi(v) = 0\}$$



$$\text{Im } \varphi = \{\varphi(v) : v \in V\}$$

are both subalgebras.

Suppose  $A$  is a subalgebra of  $V$



- cosets  $v+A$  cover  $V$  without overlapping
- cosets are isomorphic as sets
- try to make the cosets into an A.A.

$$(v+A) + (v'+A) \mapsto v+v'+A$$

$$(v+A) \cdot (v'+A) \mapsto vv'+A$$

This works if  $A$  is an Ideal

def: if  $v \cdot a \in A$  &  $a \in A$   
for all  $v \in V, a \in A$ .

This is a function if

$$(v+A, v'+A) = (w+A, w'+A)$$

$$\Rightarrow vv'+A = ww'+A$$

$$\Rightarrow v=w+a, v'=w'+a' \quad a, a' \in A$$

$$vv' = ww' + aw' + wa' + aa'$$

example  $V \equiv \{\text{Real functions on } \mathbb{R} \text{ with pointwise multiplication}\}$

- The subset of bounded functions in  $V$  is a subalgebra but not an ideal because a bounded function multiplied by a non-bounded function may not be bounded.
- The subset of functions in  $V$  which are zero except on  $[0, 1]$  is an ideal.

As with groups and normal subgroups,  $\text{Ker } \gamma$  is an ideal in the category of real or complex a.a.s.

$$\begin{array}{ccccc} V & \xrightarrow{\text{epi}} & V/\text{Ker } \gamma & \xrightarrow{\text{inv}} & \text{Im } \gamma \xrightarrow{\text{mon}} W \\ & & \downarrow \gamma & & \end{array}$$

commutes with natural mappings as usual.

Free associative algebra on  $V$  is  $V^0 \oplus V^1 \oplus V^2 \oplus \dots$

$$\begin{array}{ccc} S^0 \oplus S^1 \oplus S^2 \oplus \dots \oplus \dots & \longrightarrow & V^0 \oplus V^1 \oplus V^2 \oplus \dots \\ & \searrow \gamma & \downarrow \gamma \\ & & W \end{array}$$

This is much less scary than it looks.  $S^0 \oplus S^1 \oplus \dots$  is just the set of (possibly empty) lists of elements from the set  $S$  and  $V^0 \oplus V^1 \oplus \dots$  are linear combinations of the same with concatenation for multiplication.

# Topology

A topological space is a set  $X$  with a collection of "open" subsets satisfying

- (a)  $\emptyset$  and  $X$  are both open
- (b) If  $\mathcal{O}_\lambda$  are open,  $\bigcup \mathcal{O}_\lambda$  is open
- (c) If  $\mathcal{O}, \mathcal{O}'$  are open,  $\mathcal{O} \cap \mathcal{O}'$  is open

## examples

- Let  $X$  be any set. Let open sets be all subsets of  $X$ . This is the discrete topology on  $X$ .
- Let  $X$  be any set. Let only  $\emptyset$  and  $X$  be open sets. This is the indiscrete topology on  $X$ .
- Let  $X$  be a metric space and define  $B_x^\epsilon = \{y \in X : d(y, x) < \epsilon\}$ . Let a subset  $\mathcal{O}$  of  $X$  be open if, for each  $x \in \mathcal{O}$ , there is a  $B_x^\epsilon \subset \mathcal{O}$  with  $\epsilon > 0$ . This is a topology because
  - (a) Every  $x \in \emptyset$  vacuously satisfies the condition.
  - (b)  $B_x^\epsilon \subset X$  for any  $\epsilon > 0 \Rightarrow X$  is open.
  - (c) If  $\mathcal{O}_\lambda$  are all open and  $x \in \bigcup_\lambda \mathcal{O}_\lambda$ ,  $x \in \mathcal{O}_{\bar{\lambda}}$  for some  $\bar{\lambda} \Rightarrow B_x^\epsilon \subset \mathcal{O}_{\bar{\lambda}} \subset \bigcup_\lambda \mathcal{O}_\lambda$  for some  $\epsilon > 0$ .
  - (d) If  $\mathcal{O}, \mathcal{O}'$  are open, let  $x \in \mathcal{O}, \mathcal{O}'$ . Let  $B_x^\epsilon \subset \mathcal{O}$ ,  $B_x^{\epsilon'} \subset \mathcal{O}' \Rightarrow B_x^{\min(\epsilon, \epsilon')} \subset \mathcal{O} \cap \mathcal{O}' \Rightarrow \mathcal{O} \cap \mathcal{O}'$  is open.

This is the standard topology on  $\mathbb{R}^n$ .

Def: A set is called "closed" if it is the complement of an open set.

$\Rightarrow$  (a)  $\emptyset$  and  $X$  are closed

(b) If  $C_\lambda$  are closed, so is  $\bigcap C_\lambda$

(c) If  $C$  and  $C'$  are closed,  $C \cup C'$  is closed.

Note: Sets may be open, closed, both, neither.

### Characterizing open sets

Theorem: A subset  $A$  of a topological space  $X$  is open iff every  $a \in A$  has an open neighbourhood  $\mathcal{O} \subset A$ .

Proof: Suppose  $A$  is open. Then  $A \subset A$  is certainly an open neighbourhood of any  $a \in A$ . Conversely, suppose that the property holds, then  $A = \bigcup_a \mathcal{O}_a$  is open.

Def. Let  $A$  be a subset of topological space  $X$ .

$\text{Int}(A)$  = The union of all open subsets of  $A$

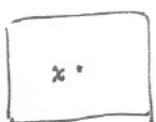
$\text{Cl}(A)$  = The intersection of all closed subsets of  $A$

$\text{Bdry}(A) = \text{Cl}(A) - \text{Int}(A)$

Theorem:  $A$  is open iff it contains no boundary points.

Theorem: Single points in a Hausdorff space are closed.

Proof:



$$\{x\}^c = \bigcup_y \mathcal{O}_y \text{ is open}$$

where  $x \notin \mathcal{O}_y$ , by Hausdorff metric.