

Topological Vector Spaces

Let's see if we can approximate a continuous function

$$X \xrightarrow{f} Y \quad \text{with a linear function}$$

$$X \xrightarrow{L} Y \quad L \in \text{Lin}(X; Y).$$

For this to be meaningful, X, Y should have linear structure as well as a topology. We'll assume that X and Y are real, finite-dimensional "topological vector spaces."

This is just a Hausdorff space where $(x, x') \mapsto x + x'$, $(a, x) \mapsto ax$ and $x \mapsto -x$ are continuous. Geroch proves that any such space is isomorphic to the euclidean topology in \mathbb{R}^n given by the metric

$$d : (x, x') \mapsto |x - x'|$$

where $\|\cdot\| : X \rightarrow \mathbb{R}$ is a norm i.e.

$$|ax| = |a||x| \quad a \in \mathbb{R}$$

$$|x + x'| \leq |x| + |x'|$$

$$|x| \geq 0 \quad \text{and } |x| = 0 \Rightarrow x = 0.$$

Exercise: Prove that d , as defined, is a metric.

We can now use the norm to measure how close a linear map is to a continuous map.

Aside: Is \mathbb{R}^n boring? How about this? Are all manifolds over topological space \mathbb{R}^n isomorphic?

Answer: Yes, except for $n=4$ when there are an infinite number of solutions. Donaldson, 1986 Fields Medal.

$$X \xrightarrow{f} Y \quad f \in \text{Mor}(X, Y)$$

$$X \xrightarrow{L} Y \quad L \in \text{Lin}(X; Y)$$

For a fixed $x \in X$, let R be defined by

$$f(x+h) = f(x) + L(h) + R(x, h)$$

$$\text{If } h \mapsto \begin{cases} |R(x, h)| / |h| & \text{if } h \neq 0 \\ 0 & \text{if } h = 0 \end{cases} \text{ is continuous}$$

then f is said to be differentiable at x . In this case, L is called the differential of f at x , often denoted $df_x \in \text{Lin}(X; Y)$.

The differentials have various nice properties, e.g.

Theorem : Differentials are unique.

Proof. Suppose that L and L' are both differentials of f at x and L and L' differ at some nonzero $h \in X$.

$$\text{Then } \frac{L(h) - L'(h)}{|h|} = \frac{R'(x, h)}{|h|} - \frac{R(x, h)}{|h|} \quad \text{if } h \neq 0$$

$$\Rightarrow |(L-L')(h/|h|)| = \left| \frac{R'(x, h)}{|h|} - \frac{R(x, h)}{|h|} \right| \leq \frac{|R'(x, h)|}{|h|} + \frac{|R(x, h)|}{|h|}.$$

Since the quantity on the left is constant on the line $a \cdot h$ $a \in \mathbb{R}^+$, any neighbourhood of 0 in X contains such a point (why?)

$$\text{and so } h \mapsto \begin{cases} \frac{|R'(x, h)|}{|h|} + \frac{|R(x, h)|}{|h|} & \text{if } h \neq 0 \\ 0 & \text{if } h = 0 \end{cases} \text{ is not}$$

continuous $\Rightarrow L$ or L' are not differentials of f at $x \Rightarrow L = L'$.

exercise: Fill in the details of this proof.

Besides uniqueness, the most important fact about differentials is

Theorem (The chain rule): If $X \xrightarrow{f} Y$ and $Y \xrightarrow{g} Z$ are differentiable, then $X \xrightarrow{g \circ f} Z$ is differentiable and, at any $x \in X$, $d(g \circ f)_x = dg_{f(x)} \circ df_x$.

Proof. Given $x \in X$, let $f(x+h) = f(x) + df_x(h) + R_f(x, h)$.

Then $g(f(x+h)) = g(f(x)) + dg_{f(x)} \circ df_x(h) + R_g(f(x), h)$

where $R(x, h) = dg_{f(x)}(R_f(x, h)) + R_g(f(x), df_x(h) + R_f(x, h))$.

Since $h \mapsto \frac{|R(x, h)|}{|h|}$ if $h \neq 0$: if $h=0$ is continuous, we're done.

Exercise: Fill in the details of the proof.

Notice that these two results give us a functor

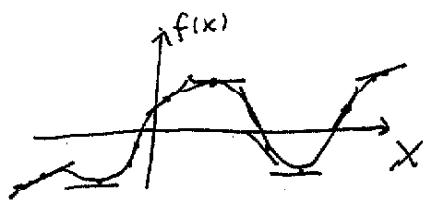
$$(X, x) \xrightarrow{f} (Y, y) \quad \text{Pointed topological spaces}$$

$$T(X, x) \xrightarrow{Tf} T(Y, y) \quad \text{Real vector spaces, } Tf = df_x$$

T is called the "Tangent-space" functor.

Notice how the functorial condition $T(g \circ f) = Tg \circ Tf$ is exactly the chain rule.

It's often useful to look at the differential as a "section" of a fiber bundle.



$$\text{Let } TX = \{(x, v) : x \in X, v \in T(x, x)\}$$

This is the "tangent bundle" with projection $TX \xrightarrow{\pi} X$, $\pi : (x, v) \mapsto x$.

Then, the differential of f is a particular section of TX , i.e. a $TX \xleftarrow{d} X$ such that $\pi \circ d = i_x$ i.e. a choice of $v \in T(X, x)$ for every $x \in X$.

In general the category of bundles over X is

Objects: $E \xrightarrow{\pi} X$, $F \xrightarrow{\nu} X$, etc.

Morphisms: $(E, \pi) \xrightarrow{g} (F, \nu)$ s.t. $\begin{array}{ccc} E & \xrightarrow{\pi} & F \\ g \downarrow & \swarrow & \nu \\ X & & \end{array}$ commutes.

Bundles are a convenient way to define, e.g., vector fields. A vector field on manifold M is a section of $TM \xrightarrow{\pi} M$ i.e. for each $p \in M$, a vector

$$v(p) = a_1(p) \epsilon_1^p + a_2(p) \epsilon_2^p + \dots + a_m(p) \epsilon_m^p$$

where $\{\epsilon_1^p, \epsilon_2^p, \dots, \epsilon_m^p\}$ are a basis of $T(M, p) = "T_p(M)"$.

Differential Forms

Let's consider differentiable functions $U \xrightarrow{f} \mathbb{R}$ where U is an open subset of \mathbb{R}^n . The simplest such functions are projections

$$\text{"}\chi_i\text{"}: (a_1, a_2, \dots, a_i, \dots, a_n) \mapsto a_i$$

Since $\chi_i(x+h) = \chi_i(x) + \chi_i(h) + 0$, χ_i is differentiable with $d\chi_i: h \mapsto h_i$. Notice that $d\chi_i \in \text{Lin}(\mathbb{R}^n; \mathbb{R}) = (\mathbb{R}^n)^*$, and $\{dx_1, dx_2, \dots, dx_n\}$ are a basis for $(\mathbb{R}^n)^*$.

Thus, in general,

$$df_x(h) = \sum_{j=1}^n a_j(x) dx_j(h)$$

The $a_j(x)$ are called partial derivatives, usually written $a_j(x) = \frac{\partial f}{\partial x_j}(x)$.

Exercise: Consider $f(x+\lambda e_j)$ and show that $a_j(x)$ is the partial derivative of f defined in the usual way.

Given any vector space V , (e.g. $(\mathbb{R}^n)^*$), we have already defined $V \wedge V$, $V \wedge V \wedge V$, etc., meaning that

$$dx_1 dy + dy_1 dz \in \Lambda^3 V \text{ etc. are defined.}$$

We would like to extend calculus to these geometric objects as well.

If we define a "k-upla" $I = (i_1, i_2, \dots, i_k)$ to be an increasing list of k integers between 1 and n , then

$$dx_I = dx_{i_1} \wedge dx_{i_2} \wedge \dots \wedge dx_{i_k}$$

is a basis of $\Lambda^k(\mathbb{R}^n)^*$. Thus, the general k -form can be written

$$\omega = \sum_I a_I dx_I \quad a_I \in \mathbb{R}$$

note that a 0-form is a real constant. We can extend this to $U \subset \mathbb{R}^n$ by letting

$$\omega(x) = \sum_I a_I(x) dx_I$$

where $a_I(x)$ are smooth functions from U to \mathbb{R} . In this context, ω is called a differentiable k -form.

Notice that differentiable 0-forms are simply real valued functions on U . This means that we already have

$$\Lambda^0(\mathbb{R}^n)^* \xrightarrow{d} \Lambda^1(\mathbb{R}^n)^*$$

i.e. the differential. We can extend this to $\Lambda^k(\mathbb{R}^n)^* \xrightarrow{d} \Lambda^{k+1}(\mathbb{R}^n)^*$ by

$$d\omega = \sum_I da_I \wedge dx_I$$

d is called the "exterior derivative".

The exterior derivative has these properties

- If $f \in \Lambda^0(\mathbb{R}^n)^*$, $d(f)$ is the differential
- $d(\omega_1 + \omega_2) = d\omega_1 + d\omega_2$
- $d(d\omega) = 0$
- $d(\omega \wedge \varphi) = d\omega \wedge \varphi + (-1)^k \omega \wedge d\varphi$ where ω is a k -form and φ is an s -form.

Exercise. Prove a), b), d) above.

To prove c), let $f : \mathbb{R}^n \rightarrow \mathbb{R}$ be a 0-form. Then

$$\begin{aligned} d(df) &= d\left(\sum_{j=1}^n \frac{\partial f}{\partial x_j} dx_j\right) = \sum_{j=1}^n d\left(\frac{\partial f}{\partial x_j}\right) \wedge dx_j \\ &= \sum_{j=1}^n \sum_{k=1}^n \frac{\partial^2 f}{\partial x_j \partial x_k} dx_k \wedge dx_j = 0 \quad \text{since } \frac{\partial^2 f}{\partial x_j \partial x_k} = \frac{\partial^2 f}{\partial x_k \partial x_j} \end{aligned}$$

for smooth f . Using d), this is true for any k -form.

If $d\omega = 0$, ω is called closed, i.e. $\omega \in \text{Ker } d$.

If $\omega = d\varphi$ for some form φ , ω is called exact,

i.e. $\omega \in \text{Im } d$.

Since $d^2 = 0$, every exact form is closed.

Theorem (Poincaré'): If $U \subset \mathbb{R}^n$ is contractible, then a K -form on U is closed iff it is exact.

Example: \mathbb{R}^3

For $f \in \Lambda^0(\mathbb{R}^3)^*$,

$$df = \frac{\partial f}{\partial x} dx + \frac{\partial f}{\partial y} dy + \frac{\partial f}{\partial z} dz$$

For $f_1 dx + f_2 dy + f_3 dz \in \Lambda^1(\mathbb{R}^3)^*$

$$d(f_1 dx + f_2 dy + f_3 dz) =$$

$$\left(\frac{\partial f_1}{\partial y} - \frac{\partial f_2}{\partial x} \right) dy \wedge dx - \left(\frac{\partial f_1}{\partial z} - \frac{\partial f_3}{\partial x} \right) dx \wedge dz + \left(\frac{\partial f_2}{\partial z} - \frac{\partial f_3}{\partial y} \right) dy \wedge dz$$

For $\omega = f_1 dy \wedge dz + f_2 dz \wedge dx + f_3 dx \wedge dy \in \Lambda^2(\mathbb{R}^3)^*$

$$\begin{aligned} d\omega &= \frac{\partial f_1}{\partial x} dx \wedge dy \wedge dz + \frac{\partial f_2}{\partial y} dy \wedge dz \wedge dx + \frac{\partial f_3}{\partial z} dz \wedge dx \wedge dy \\ &= \left(\frac{\partial f_1}{\partial x} + \frac{\partial f_2}{\partial y} + \frac{\partial f_3}{\partial z} \right) dx \wedge dy \wedge dz \end{aligned}$$

$d(0\text{-form}) \cong \text{Gradient}$

$d(1\text{-form}) \cong \text{Curl}$

$d(2\text{-form}) \cong \text{divergence}$

Unlike tensor fields, differential forms can be pulled back across smooth maps.

$$\mathbb{R}^n \xrightarrow{f} \mathbb{R}^m \quad f \text{ smooth}$$

If we have a k -form ω on \mathbb{R}^m , we can define

$$\omega = \sum_I a_I dy_{i_1} \wedge dy_{i_2} \wedge \cdots \wedge dy_{i_k}$$

$$f^* \omega = \sum_I (a_I \circ f) d(y_{i_1} \circ f) \wedge d(y_{i_2} \circ f) \wedge \cdots \wedge d(y_{i_k} \circ f)$$

Pullbacks have these properties

$$(a) f^*(\omega + \omega') = f^*\omega + f^*\omega'$$

$$(b) f^*(g\omega) = f^*(g)f^*(\omega)$$

$$(c) f^*(g_1 g_2 \wedge \cdots \wedge g_k) = f^*(g_1) \wedge \cdots \wedge f^*(g_k)$$

for 1-forms g_i .

They also interact nicely with the exterior derivative

$$d(f^*\omega) = f^*(d\omega).$$

~~For example, f^* , f^* corresponds to~~

Example. Let $U = \{(r, \theta) : r > 0, 0 < \theta < 2\pi\} \subset \mathbb{R}^2$

$\varphi : U \rightarrow \mathbb{R}^2$ be $\varphi : (r, \theta) \mapsto (r \cos \theta, r \sin \theta)$.

Then the pullback of the volume form $dx \wedge dy$ is given by

$$\varphi^*(dx \wedge dy) = d(x \circ \varphi) \wedge d(y \circ \varphi) = r dr \wedge d\theta$$

Exercise: Fill in the steps

Exercise: Let $\omega = \frac{-y}{x^2+y^2} dx + \frac{x}{x^2+y^2} dy$ be a 1-form in $\mathbb{R}^2 - \{0\}$. Show that $\varphi^*\omega = d\theta$.

Exercise: Let $f : (r, \theta, y) \mapsto (r \sin \theta \cos \varphi, r \sin \theta \sin \varphi, r \cos \theta)$.

Show that $f^*(dx \wedge dy \wedge dz) = r^2 \sin \theta \ dr \wedge d\theta \wedge dy$.

Notice that the pullback of an isomorphism is "change of variables".

The sequence

$$\Lambda^0(\mathbb{R}^n)^* \xrightarrow{d} \Lambda^1(\mathbb{R}^n)^* \xrightarrow{d} \dots \xrightarrow{d} \Lambda^n(\mathbb{R}^n)^*$$

is called "the de Rham complex": $\Lambda^*(\mathbb{R}^n)^* = \bigoplus_k \Lambda^k(\mathbb{R}^n)^*$

The direct sum of vector spaces $H^k = \text{Ker } d / \text{Im } d$

$\bigoplus_k H^k$ is called "the de Rham cohomology."

de Rham showed that this is a topological invariant.

More generally, a direct sum $\bigoplus_k V^k$ of vector spaces with

$$V^0 \xrightarrow{d} V^1 \xrightarrow{d} V^2 \xrightarrow{d} \dots$$

$d^2 = 0$, is called a differential complex. $\bigoplus_k \text{Ker } d / \text{Im } d$ is called the cohomology of V . A morphism between differential complexes A and B is defined to be a sequence of morphisms $f: A^k \rightarrow B^k$ s.t

$$\begin{array}{ccccccc} A^0 & \xrightarrow{d} & A^1 & \xrightarrow{d} & A^2 & \xrightarrow{d} & \dots \\ f \downarrow & & \downarrow f & & \downarrow f & & \downarrow f \\ B^0 & \xrightarrow{d} & B^1 & \xrightarrow{d} & B^2 & \xrightarrow{d} & \dots \end{array}$$

commutes. In this language, Λ^* is a cofunctor from the category of Euclidean spaces \mathbb{R}^n to the category of differential complexes

$$\mathbb{R}^n \xrightarrow{f} \mathbb{R}^m \quad \Lambda^* \mathbb{R}^n \xleftarrow{f^*} \Lambda^* \mathbb{R}^m$$

where f^* is the pullback.

Hodge Star

An inner product $\langle \cdot, \cdot \rangle : V \times V \rightarrow \mathbb{R}$ on a real finite dimensional vector space V is a bilinear symmetric product which is "nondegenerate" meaning that

$$\langle v, x \rangle = 0 \text{ for all } x \in V \Rightarrow v = 0.$$

Lemma. For $f \in V^*$, $f = (v \mapsto \langle v_f, v \rangle)$ for some unique $v_f \in V$.

Proof. $\Psi: x \mapsto (v \mapsto \langle x, v \rangle)$ from V to V^* is one-to-one (using the nondegenerate property) $\Rightarrow \Psi$ is an isomorphism.

Let $\dim(V) = n$, $V \wedge V \wedge \cdots \wedge V = \Lambda^n$. Fix $\lambda \in \Lambda^n$.

Then the mapping

$m \mapsto \lambda \wedge m$ from Λ^{n-k} to $\Lambda^n \cong \mathbb{R}$ is

$\lambda \wedge m = \langle v_\lambda, m \rangle \sigma$ for some fixed basis vector $\sigma \in \Lambda^n$ for some special unique $v_\lambda \in \Lambda^{n-k}$. This v_λ is denoted $*\lambda$, the "Hodge Star" of λ . Notice that $*\lambda$ depends both on the inner product and on the choice of basis $\sigma \in \Lambda^n$.

For example, let $n=3$, $\epsilon_x, \epsilon_y, \epsilon_z$ be an orthonormal basis in \mathbb{R}^3 , $\langle \cdot, \cdot \rangle$ the standard euclidean inner product, let $\delta = \epsilon_x \wedge \epsilon_y \wedge \epsilon_z$.

$$\text{Then } \epsilon_x \wedge \epsilon_y \wedge (\sum_i a_i \epsilon_i) = \langle \star(\epsilon_x \wedge \epsilon_y), \sum_i a_i \epsilon_i \rangle \epsilon_x \wedge \epsilon_y \wedge \epsilon_z$$

$$\text{is solved by } \star(\epsilon_x \wedge \epsilon_y) = \epsilon_z.$$

Similarly,

$$\star \epsilon_x = \epsilon_y \wedge \epsilon_z \quad \star(\epsilon_x \wedge \epsilon_y) = \epsilon_z$$

$$\star \epsilon_y = \epsilon_z \wedge \epsilon_x \quad \star(\epsilon_y \wedge \epsilon_z) = \epsilon_x \quad \star(\epsilon_x \wedge \epsilon_y \wedge \epsilon_z) = 1.$$

$$\star \epsilon_z = \epsilon_x \wedge \epsilon_y \quad \star(\epsilon_z \wedge \epsilon_x) = \epsilon_y$$

Exercise. Show that $\hat{A} \times \hat{B} = \star(A \wedge B)$.

The Hodge star has this property: $\alpha \wedge \star \beta = \star \beta \wedge \star \alpha$

Exercise: For $f: \mathbb{R}^n \rightarrow \mathbb{R}$, show that $d \star(df) = \Delta f \mu$ where μ is the volume form and Δ is the laplacian.

Example. With $F_{\mu\nu}$ the EM field strength, $F = \frac{1}{2} F_{\mu\nu} dx_\mu \wedge dx_\nu$ is the electromagnetic field. Maxwell's equations are

$$\begin{aligned} dF &= 0 & \rightarrow F \text{ is closed. If it is defined} \\ \star d \star F &= \star j & \text{on a contractible subset of } \mathbb{R}^4, \\ && \text{then } F = dA \text{ by Poincaré's lemma.} \end{aligned}$$