

from gravitational effects of the DM "fluid" have:

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$$m \lesssim 10^{+57} \text{ eV} \quad \text{MACHO searches}$$

$$\left(\begin{array}{l} m_{\text{sun}} \sim 10^{66} \text{ eV} \\ m_{\text{earth}} \sim 10^{60} \text{ eV} \\ m_{\text{moon}} \sim 10^{58} \text{ eV} \end{array} \right)$$

$$m_{\text{boson}} \gtrsim 10^{-22} \text{ eV} \quad \text{dwarf galaxy}$$

$$m_{\text{fermion}} \gtrsim 10^3 \text{ eV} \quad \text{" "}$$

$$\text{background density: } \sim \frac{\text{GeV}}{\text{m}^3} \quad (\text{Cosmology})$$

in halos:

$$\rho_{\text{halo}} \sim \frac{1}{r^2}$$

$$\rho_{\text{at earth}} \sim 3 \cdot 10^5 \frac{\text{GeV}}{\text{m}^3}$$

$$\rho_{\text{water}} \sim 10^{30} \frac{\text{GeV}}{\text{m}^3}$$

$$\rho_{\text{solar system}} \sim \frac{M_{\text{sun}}}{R_{\text{solar system}}^3} \sim \frac{10^{57} \text{ GeV}}{(10^3 \text{ AU})^3} \sim \frac{10^{15} \text{ GeV}}{\text{m}^3}$$

$\Rightarrow \rho_{\text{DM}}$ not observable
in solar system gravity

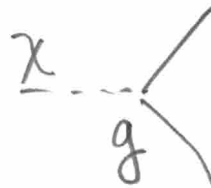
$$\uparrow \sim 10^{11} \text{ m}$$

$$V_{DM}^{\text{galaxy halo}} \sim 10^{-3} c$$

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Assume it is a "particle" $\Rightarrow M \leq M_{pl} \sim 10^{19} \text{ GeV}$

stability? massive particle χ



↑ negligible masses

$$\Gamma \sim \frac{g^2}{8\pi} m_\chi$$

$$\tau = \frac{1}{\Gamma} \gg \text{age of universe } 10^{10} \text{ yrs}$$

$$\Rightarrow \Gamma < \frac{1}{10^{10} \text{ yrs}} \frac{1}{\pi 10^7 \text{ s/yr}} \stackrel{t}{\approx} 7 \cdot 10^{-16} \text{ eVs} \approx 2 \cdot 10^{-33} \text{ eV}$$

$$\Rightarrow g^2 < 10^{-9} \left[\frac{10^{-22} \text{ eV}}{m_\chi} \right]$$

very small coupling,
even for smallest mass.

higher dimensional operator?

"weak decay"



← dimension 6 operator

$$\Gamma \sim \frac{m_\chi^5}{m_W^4}$$

$$\Rightarrow m_\chi \lesssim (\Gamma m_W^4)^{1/5} \sim (10^{-33} \text{ eV} (10^{11} \text{ eV})^4)^{1/5} \sim 100 \text{ eV}$$

dim 5 operator $\Gamma \sim \frac{m_\chi^3}{M^2}$

eg axion $M \sim 10^{10} \text{ GeV}$ $m_\chi \sim 10^{-10} \text{ GeV}$

$$\Rightarrow \Gamma \sim 10^{-50} \text{ GeV}$$

Conclusion: $g_{\text{eff}}(m_\chi) \ll 1$ or

stable due to quantum # (eg. e^- electric charge)
 P^+ baryon #)

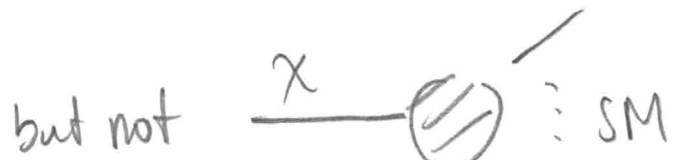
Simplest possibility: \mathbb{Z}_2 charge \leftrightarrow "DM parity"

$$\chi \rightarrow -\chi$$

$$SM \rightarrow +SM$$

χ lightest parity odd particle (e.g. SUSY R-parity)

consequence: vertex has even # of χ

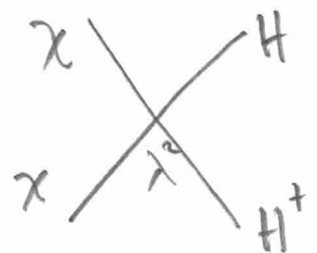


simplest: χ scalar $\mathcal{L} \sim \tilde{\lambda}^2 \chi^2 H^+ H^-$ Higgs portal DM

assuming $\tilde{\lambda}$ not too small ...

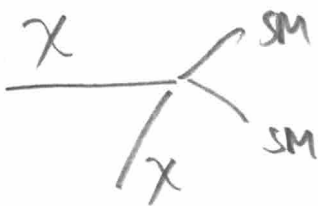
at high temperatures, densities (early universe)

$$T \gg m_\chi, m_H$$



$$\sigma \sim \frac{\alpha^2}{E_{\text{cm}}^2} \sim \frac{\alpha^2}{T^2} \quad \alpha = \frac{\lambda^2}{4\pi}$$

compare with expansion rate of universe $H \equiv \frac{\dot{R}}{R} \sim \frac{T^2}{\text{Mpc}}$

need rate of scattering 

from definition of flux $\sigma \sim \frac{1}{|v_1 - v_2|}$

$$\Gamma \sim n_\chi \langle \sigma v \rangle \sim T^3 \frac{\alpha^2}{T^2} \sim \alpha^2 T$$

↑ averaged over thermal distribution ↑ relativistic

→ scattering frequent when $\frac{\Gamma}{H} \sim \alpha^2 \frac{\text{Mpc}}{T} \gg 1$

eg. $\lambda \sim 1 \Rightarrow$ always frequent

non-relativistic $n \sim (MT)^{3/2} e^{-M/T}$

$$\Rightarrow \frac{\Gamma}{H} \sim e^{-M/T}$$

interactions become infrequent, "freeze-out"

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annihilations stop \rightarrow relic abundance of DM [Meissa's talk]

"WIMP miracle"

$$\Gamma = n_\chi \langle \sigma v \rangle \approx H$$

$$\langle \sigma v \rangle \sim \frac{\alpha^2}{m_\chi^2}$$

$$\text{get } \Omega_\chi = 0.1 \left(\frac{0.01}{\alpha} \right)^2 \left(\frac{m}{100 \text{ GeV}} \right)^2$$

want $\Omega_\chi \sim 0.25$

$$\left[\begin{array}{l} \alpha_w \sim 0.03 \\ \alpha_{em} \sim 0.01 \end{array} \right]$$

Indirect detection of DM

residual annihilation today:



rate for 1 to scatter $\Gamma \sim n \langle \sigma v \rangle$

rate for particle scatters per volume: $\frac{N}{V} \Gamma = n \cdot \Gamma = n^2 \langle \sigma v \rangle$

$$\text{but } n = \frac{\rho}{m_\chi}$$