

prediction of factorization :

fragmentation is indep of rest of event (i.e. cares only about momentum of quark)

def: $z_i \equiv \frac{p_i}{P}$ $\leftarrow$ individual hadron
 $\leftarrow$ quark jet $(0 \leq z_i \leq 1)$

fragmentation functions: $f_{q \rightarrow i}(z)$

e.g. $f_{u \rightarrow \pi^+}(z)$ probability of u quark to produce π^+ of momentum fraction $z = \frac{p_{\pi^+}}{p_u}$

(note: indep of p_i, P separately (follows from factorization and dim'l analysis, no other scale available))

f are universal.

Plot
Fig 12

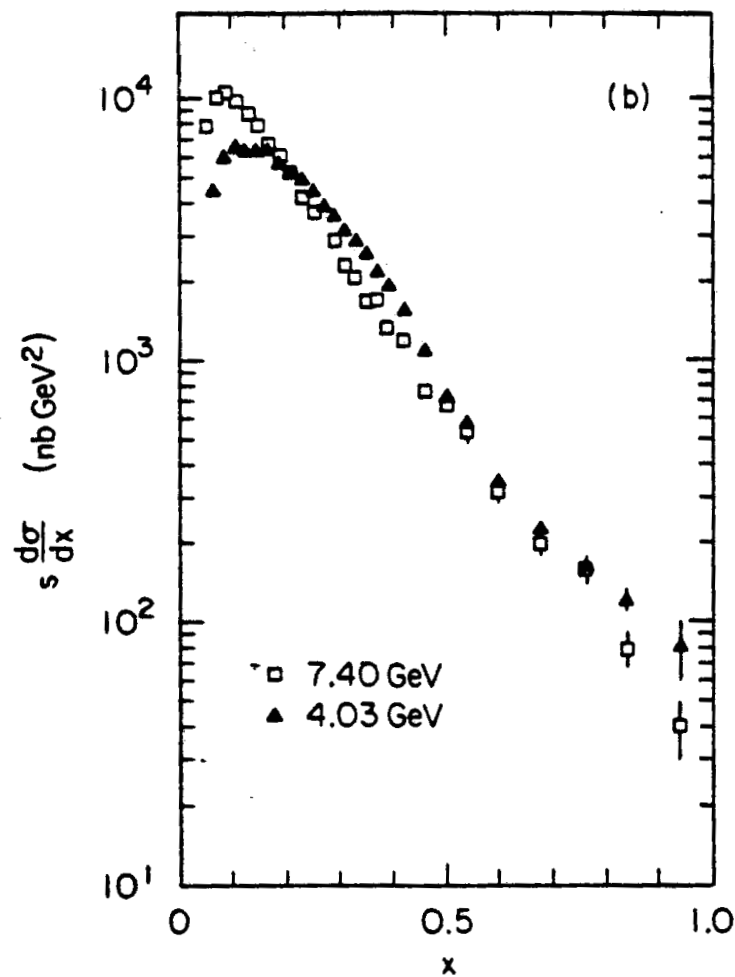
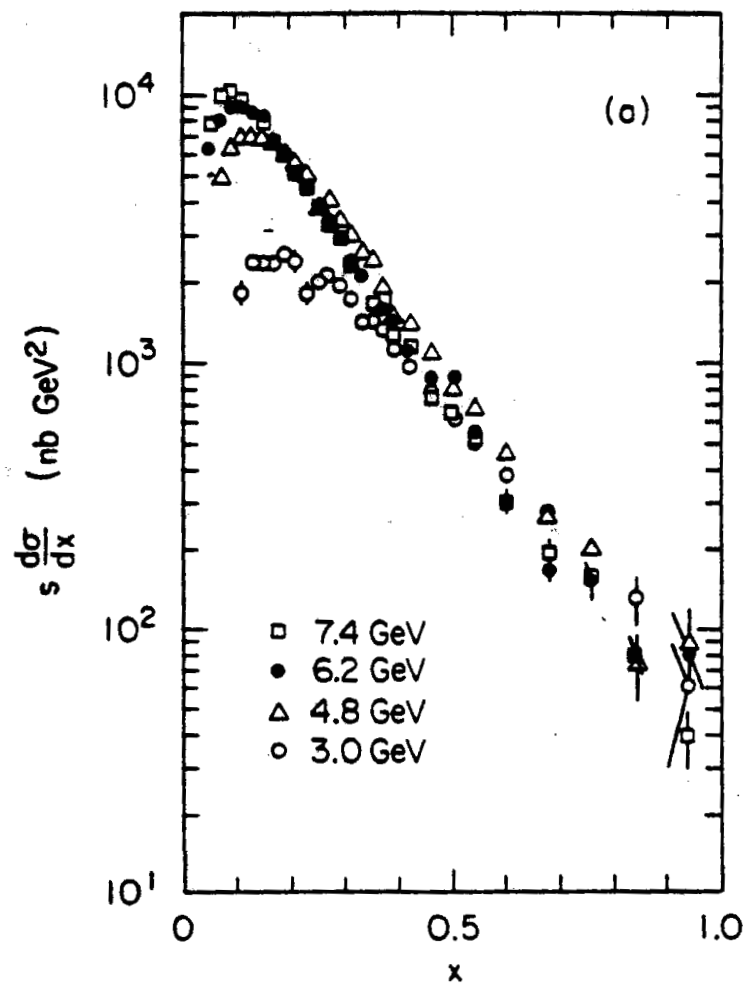
$$X = \underbrace{\frac{\hat{n} \cdot \vec{p}_{\pi}}{E_{cm}/2}}_{X\text{-axis}} \sim \frac{p_{\pi}}{P} = z$$

$$\frac{sd\sigma_{\pi}}{dx} \underset{\substack{\uparrow \\ E_{cm}^2}}{\sim} \underbrace{S \sigma_{q\bar{q}}}_{const} \cdot f_{q \rightarrow \pi}(x)$$

Fragmentation functions

J.L. Siegrist
et al
(Mark I)

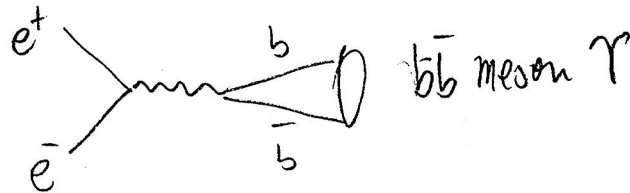
Phys. Rev. D26,
969
(1980)



$$x = z = \frac{\hat{n} \cdot \vec{p}_\pi}{(\frac{1}{2}E_{cm})}$$

$\hat{n}$ is the jet axis

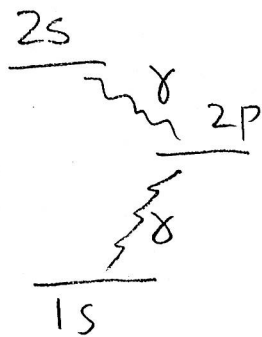
PLOT 19 $e^+e^- \rightarrow \text{hadrons near } 2m_b$



$1s, 2s, 3s, 4s$ non-relativistic bound states

$$r \sim \alpha \frac{1}{sm_b}$$

$$\Delta E_r \sim \alpha_s^2 \frac{1}{m_b}$$



spectroscopy to measure all energy levels

PLOT 20 $c\bar{c}$ E_γ continuum from $\pi^0 \rightarrow \gamma\gamma$ decays
spikes from "atomic" transitions

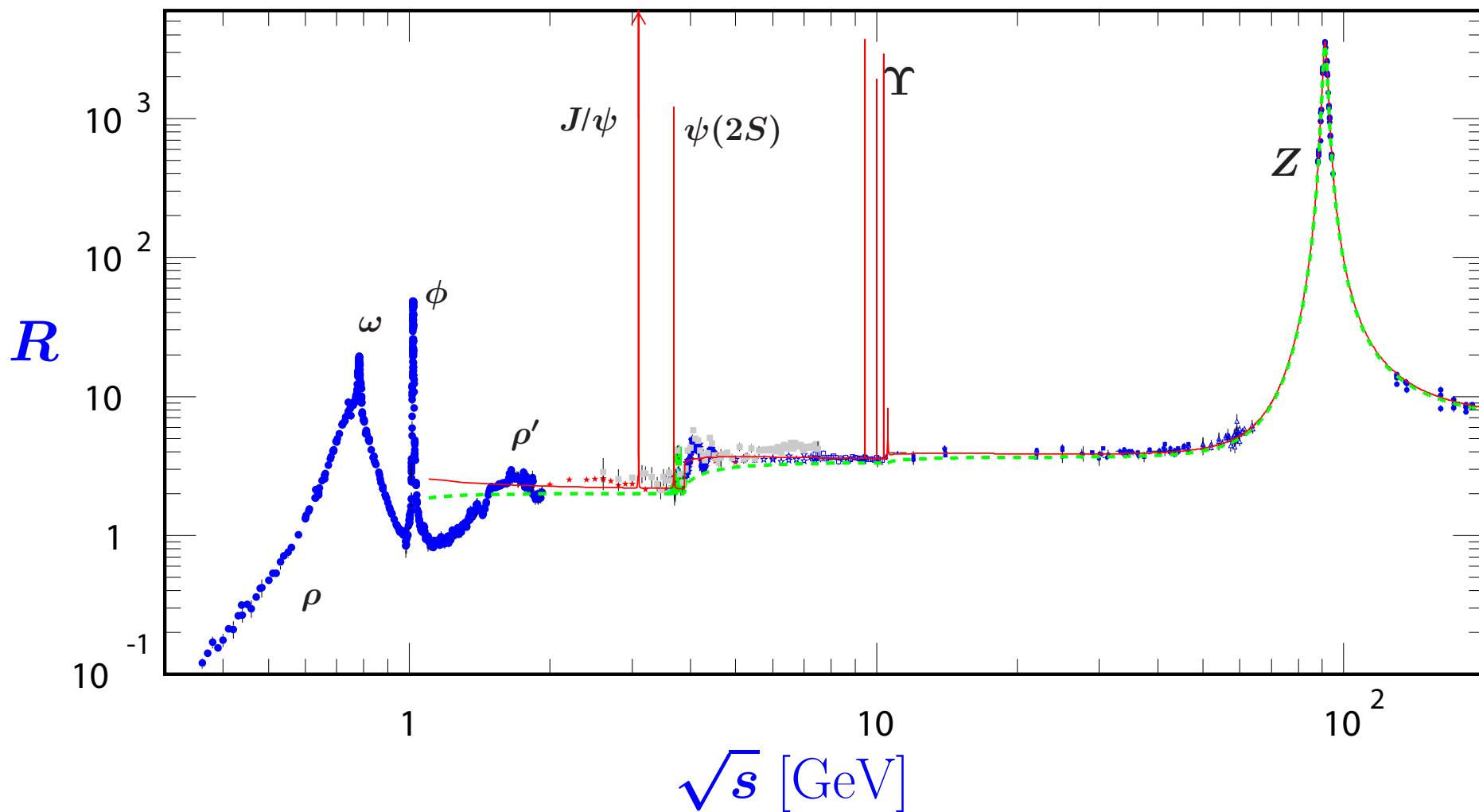
charmonium spectrum: Glashow-Lane

PLOT 24 $V(r)$ phenomenological models.

dual nature • hard scattering (QED-like)

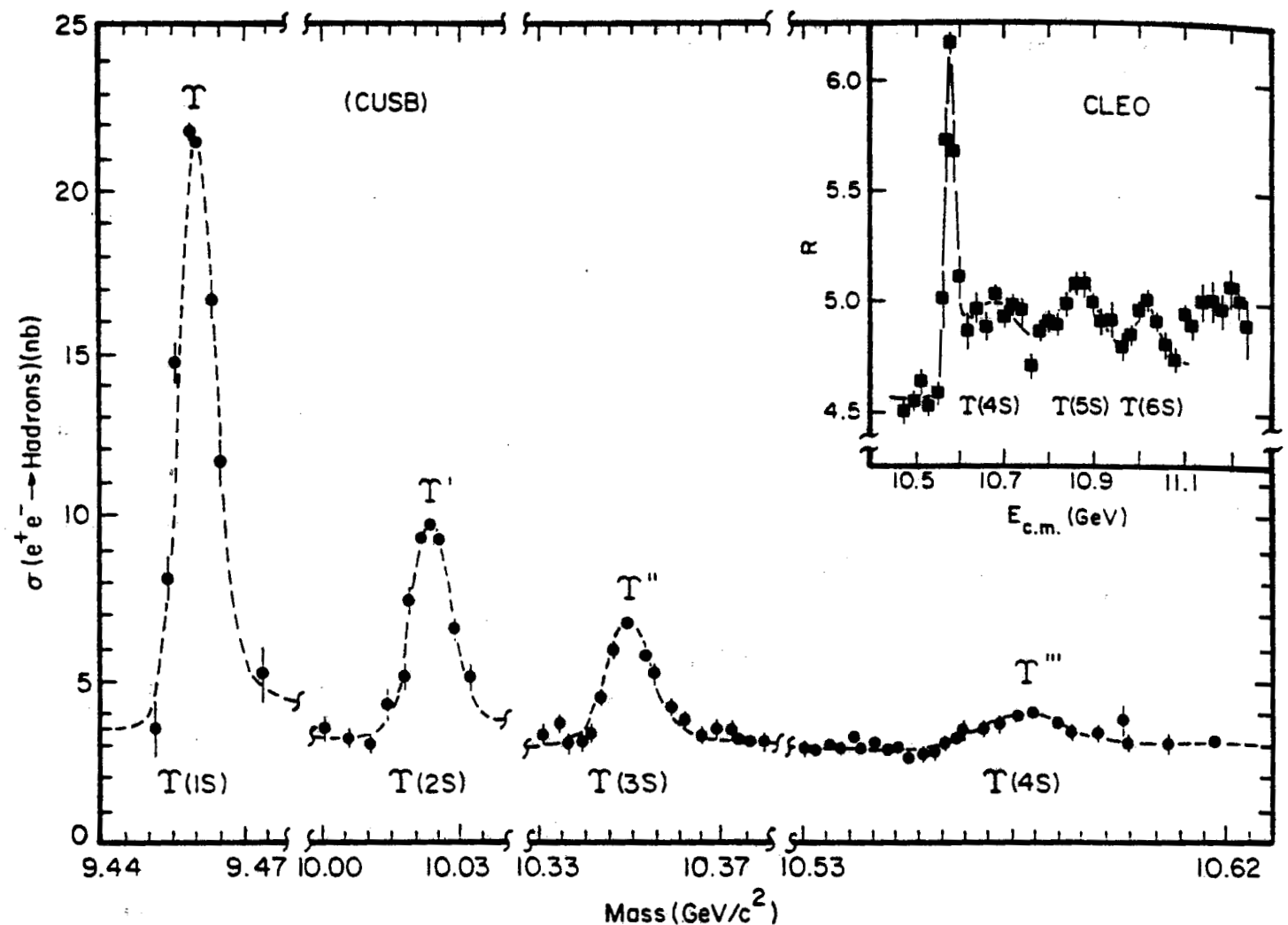
of QCD: • soft fragmentation, hadronization, bound states

R in e^+e^- Collisions



Particle Data Group compilation (2018)

Upsilon $\sim b\bar{b}^*$
bound state
resonances
near $e^+e^- \rightarrow b\bar{b}^*$
threshold

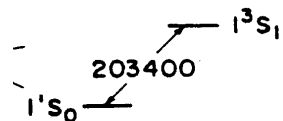
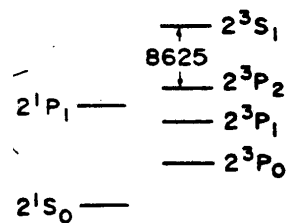


D. Besson and T. Skwarnicki

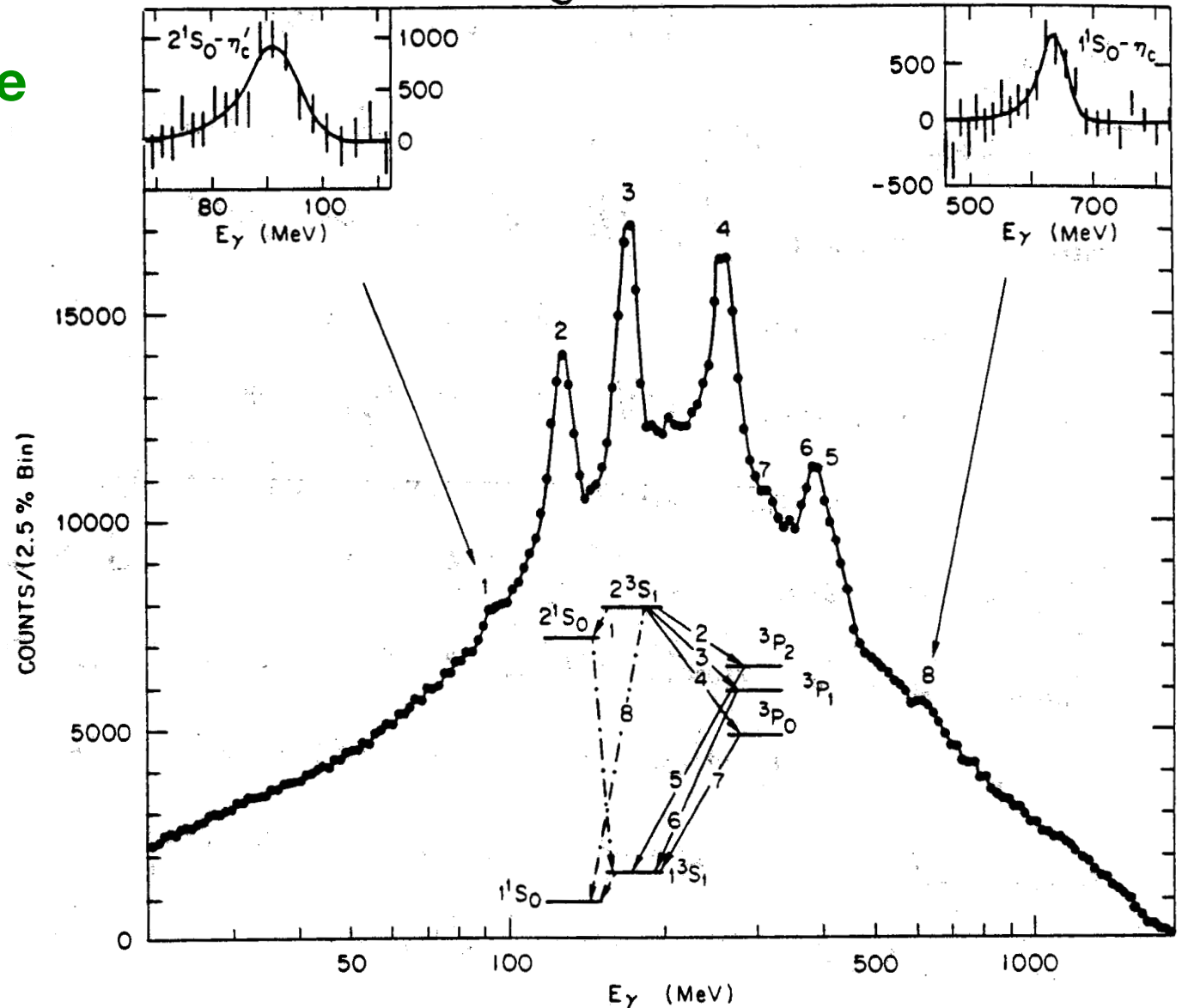
Ann Rev. Nucl. Part. Sci 43, 333 (1993)

c c* bound state "charmonium" spectroscopy

Crystal Ball at $E_{cm} = m(\psi')$

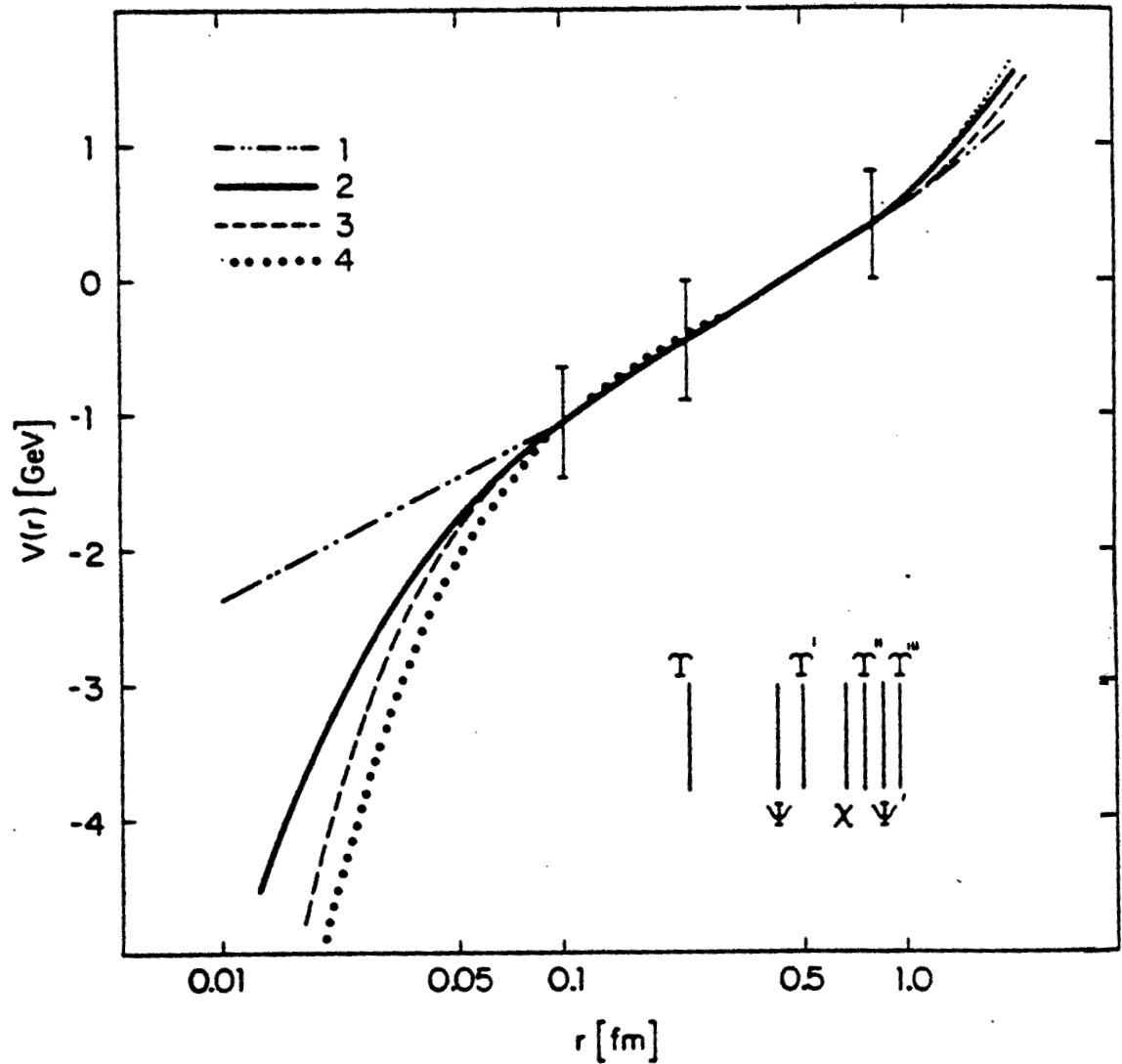


Positronium



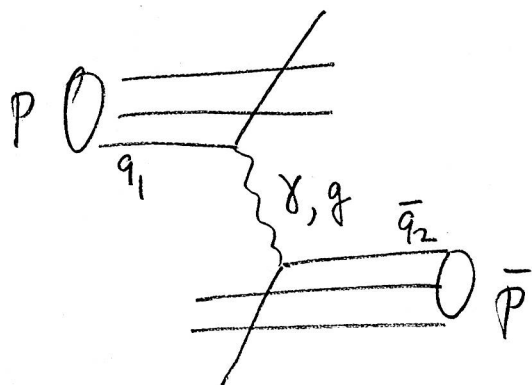
E. Bloom and C E Peck,
Ann. Rev. Nucl. Part. Sci. 33, 143 (1983)

Phenomenological
non-relativistic QM
potentials $V(r)$
(models)



W. Buchmüller and S.-H. H. Tye
Phys. Rev. D 24, 132 (1981)

$P\bar{P}$ scattering $E_{cm} \gg m_p$ (in $P\bar{P}$ -CM)

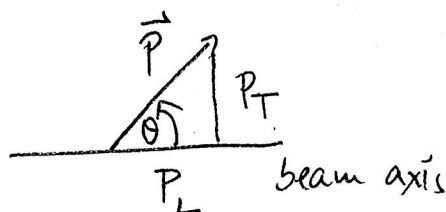


definitions:

1. $x = \frac{P_q}{P}$ momentum fraction carried by quark

$x_1 \neq x_2$ in general $\Rightarrow$ "partonic CM" has unknown boost

2. $P_T = p \sin \theta$
 $P_L = p \cos \theta$



P_L is not invariant under "boosts" $\Rightarrow$ not a good measure

P_T is boost invariant $\Rightarrow$ better measure of "hardness" of particle.

P_L not boost inv. $\Rightarrow P_L$ large when "nothing" happened



2.b. for $E \gg m$ $p \sim E$

(E measured with calorimeters)

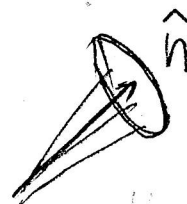
$\Rightarrow$ define $E_{Ti} = E_i \sin \theta$ (i -th particle)

$E_{T_{\text{even}}} = \sum_i E_{Ti}$ hardness of event

PLOT 15 most hadrons are soft $\star 10^6 \sim \frac{N(1 \text{ GeV})}{N(8 \text{ GeV})} \ll 546 \text{ GeV}$

PLOT 16 events can be hard
two jets $E(\theta, \varphi)$ "Lego plot"

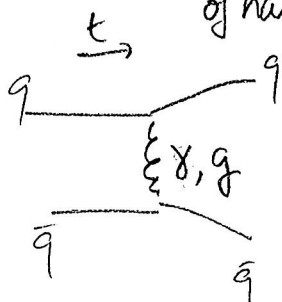
PLOT 17 $H_2 = \frac{E_T \text{ of leading jet}}{E_{T \text{ total}}}$



$$K = \sqrt{2 + \Delta\phi^2}$$

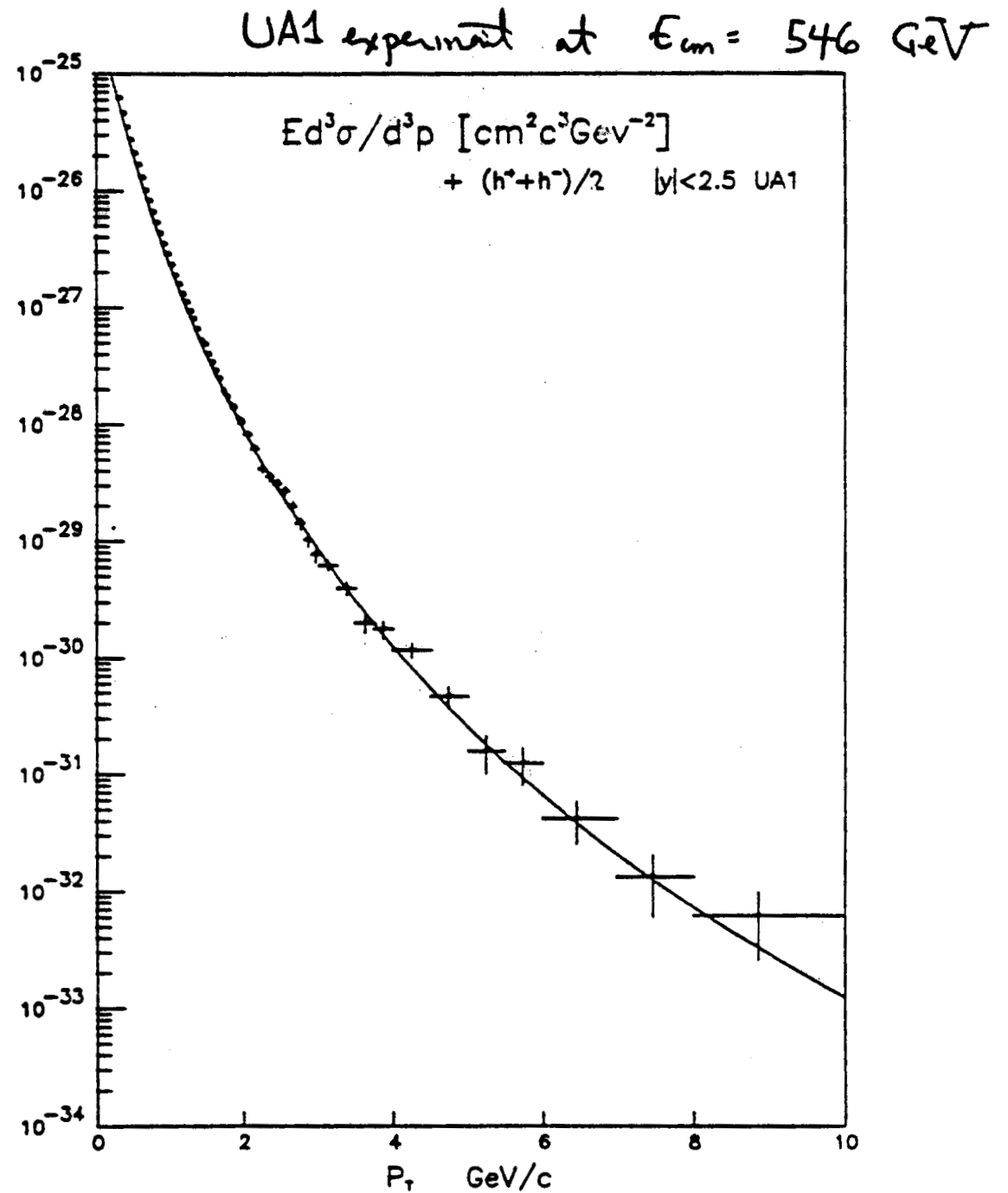
hardest events are 2-jet like.

PLOT 18 angular distribution? of hard jets $\frac{d\sigma}{d\cos\theta_{\text{cm}}} \sim \frac{1}{\sin^4 \theta/2}$ Coulomb scattering!

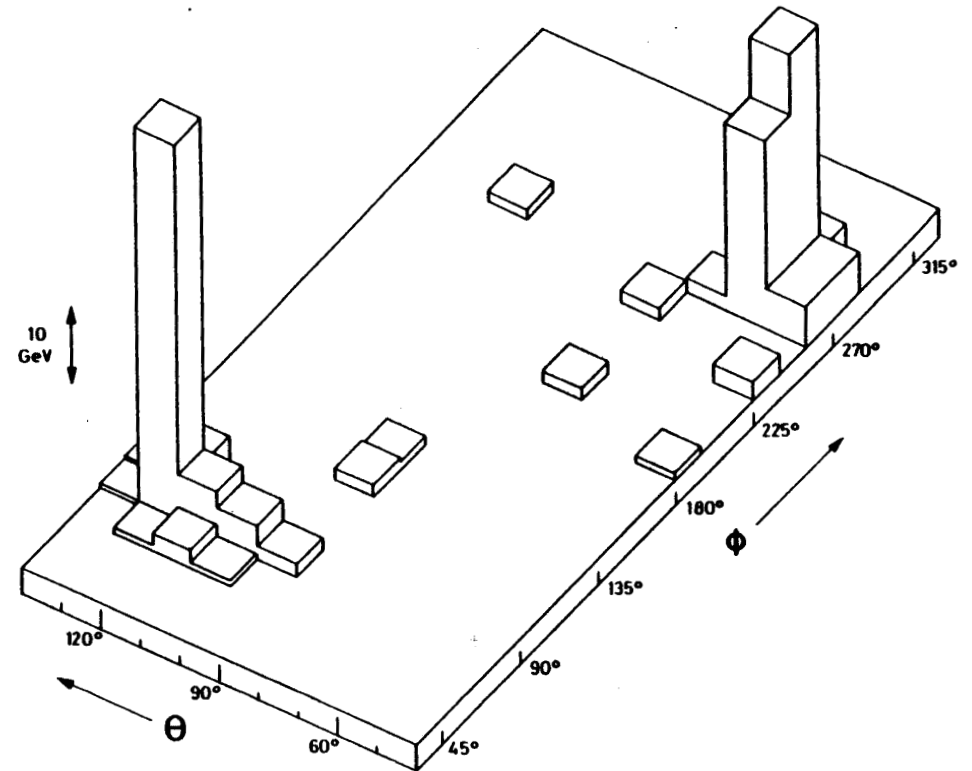
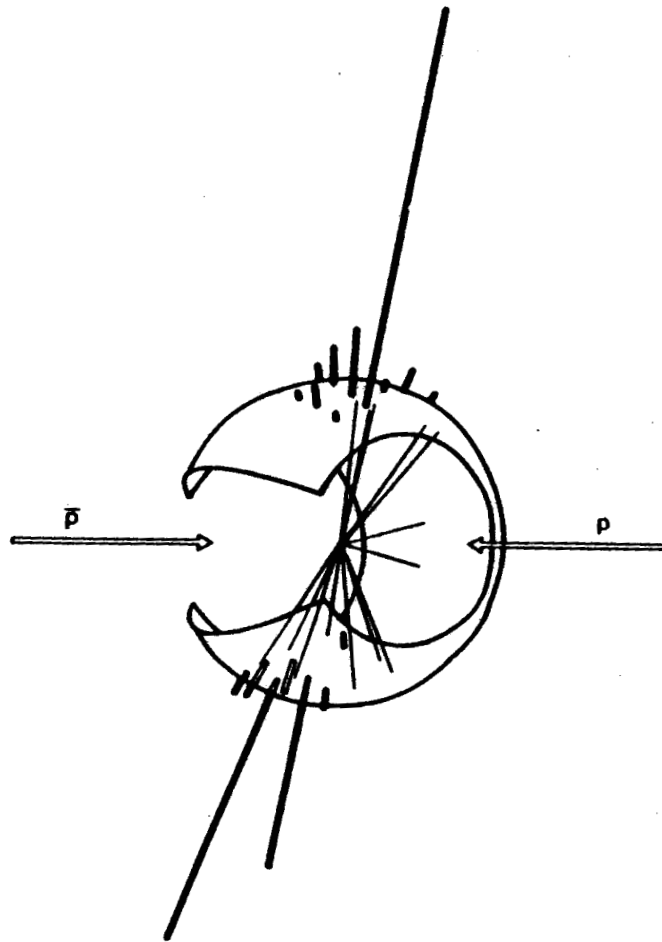


Most particles
produced are "soft":
 $p_T \ll E_{CM}$

G. Altarelli and L. DiLella,
Photon-Antiproton Collider Physics



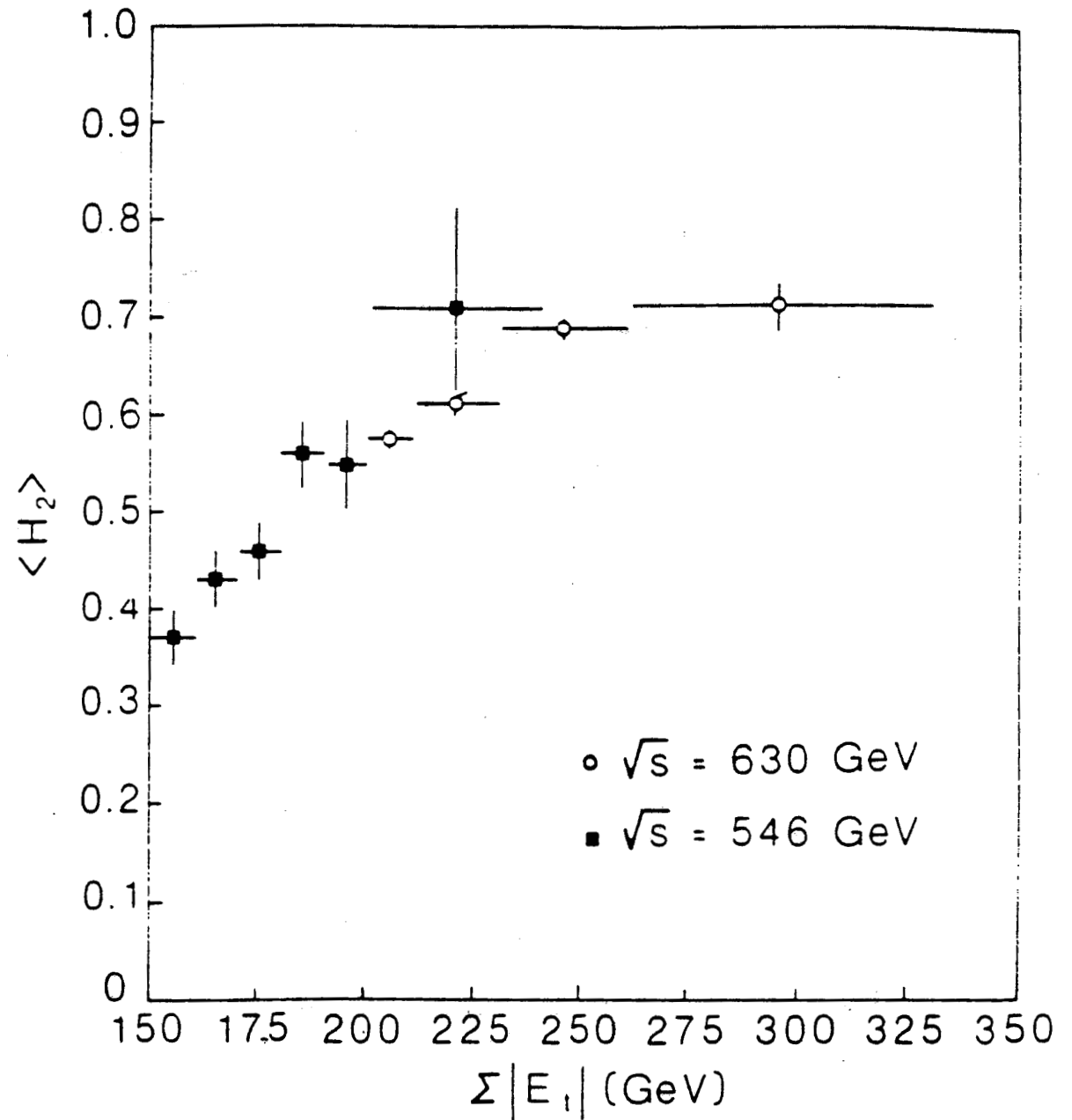
Jets can be “hard”



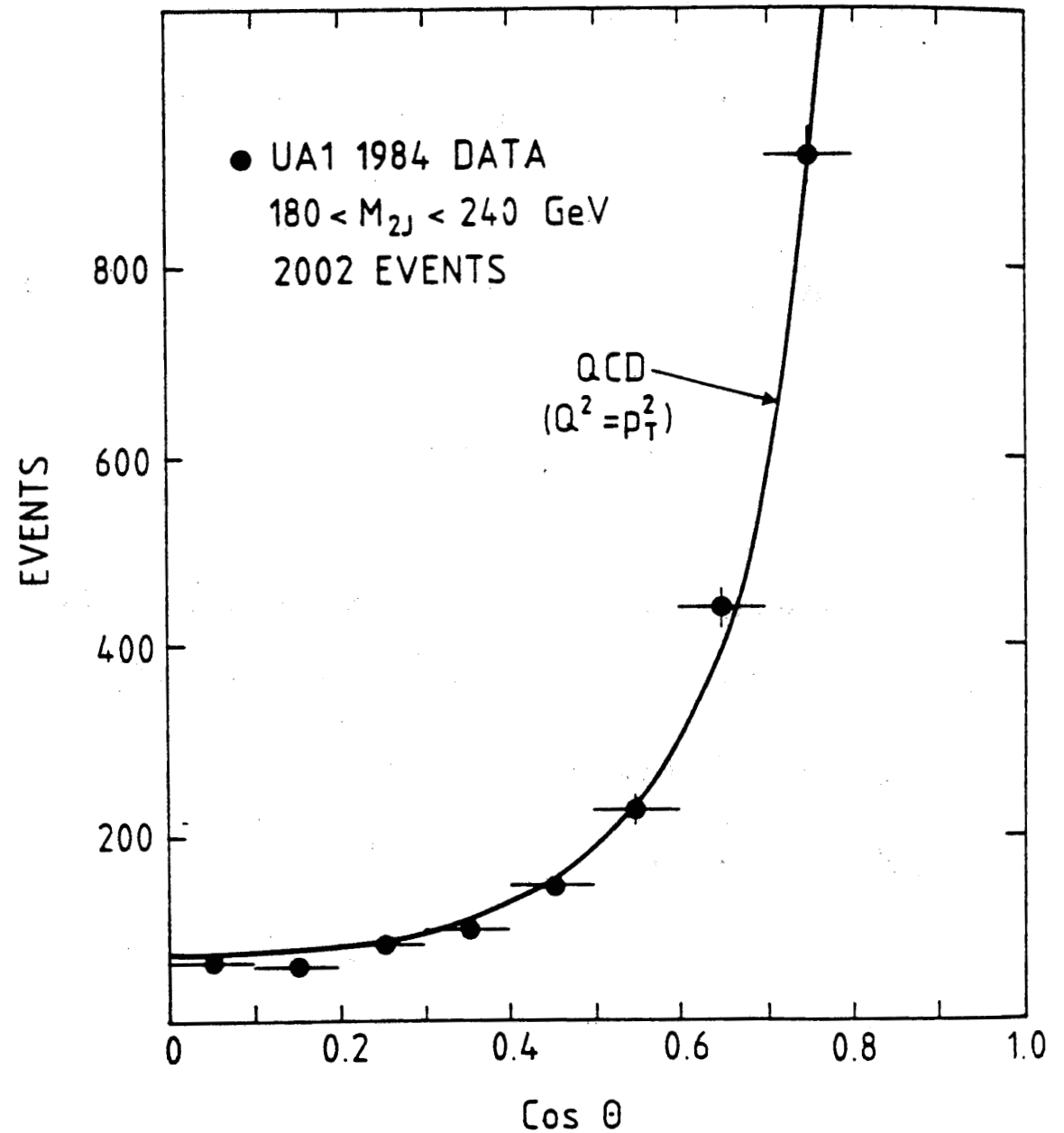
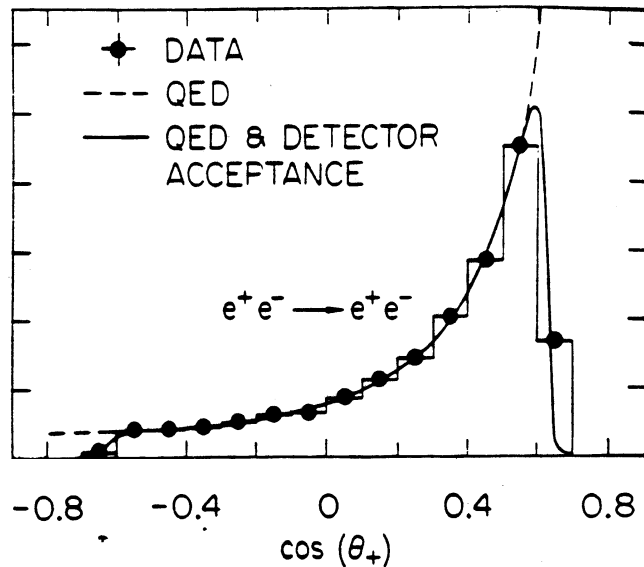
M. Banner et al (UA2)
Phys. Lett. 118B, 203 (1982)

Fraction of E_T contained in the leading two jets

C. Albajar et al (UA1)
Z Phys. C36, 33(1987)



Similarity of angular distributions of $e^+e^- \rightarrow e^+e^-$ versus $q q^* \rightarrow q q^*$,
Coulomb forward-scattering singularity



G. Arrison et al (UA1)

Phys. Lett. B177, 244 (1986)