

Measurement of missing transverse energy with the CMS detector at the LHC

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Received: 14 March 2006 / Revised version: 29 April 2006 /
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Abstract. The performance of the Compact Muon Solenoid detector for measuring missing transverse energy is evaluated using fully simulated pp collisions at a center-of-mass energy of 14 TeV at the Large Hadron Collider. For minimum bias events without pileup, a resolution of 6.1 GeV is computed, corresponding to a stochastic contribution of $0.63\sqrt{\Sigma E_T}$ GeV^{1/2}, where ΣE_T is the summed transverse energy in all calorimeter towers. When the contribution of pileup is included, the resolution degrades according to the overall deposited ΣE_T with the same stochastic coefficient. For QCD dijet events with event pileup corresponding to a luminosity $\mathcal{L} = 2 \times 10^{33}$ cm⁻²s⁻¹, we compute $\sigma = [(3.8 \text{ GeV})^2 + (0.97 \text{ GeV}^{1/2} \sqrt{\Sigma E_T})^2 + (0.012 \Sigma E_T)^2]^{1/2}$ resulting in a resolution of 45 GeV for jet events with reconstructed transverse momentum of 800 GeV/c. Monte Carlo samples of $t\bar{t}$ and W +jet events with high-momentum ($p_T > 20$ GeV/c) lepton decays leading to true missing transverse energy were used to determine the azimuthal angle resolution to be 0.1 radians (0.2 radians) for a reconstructed missing transverse energy of 200 GeV (100 GeV).

1 Introduction

The understanding of detector response to standard model physics from quantum chromodynamics (QCD) is a prerequisite to the search for new phenomena at the large hadron collider (LHC) [1]. Beginning with UA1 [2], all major detectors at hadron colliders have been designed to cover as much solid angle as practically possible with calorimetry. This was a major consideration in the design of calorimetry [3] for the Compact Muon Solenoid (CMS) [4]. The primary motivation is to provide as complete a picture as possible of the event, including the presence of one or more energetic neutrinos or other weakly-interacting stable particles though apparent missing energy. Energetic particles produced in the direction of the beam pipe make it impossible to directly measure missing energy longitudinal to the beam direction, however, the transverse energy balance can be measured with an accuracy good enough to help establish a physics signature involving one or more non-interacting particles. The W boson was discovered and its mass determined to 3% with just 6 events due to the ability of UA1 to infer the presence of 40-GeV neutrinos with a resolution of a few GeV [5]. Since the time of the W discovery, measurement of missing trans-

verse energy has been a major tool in the search for new phenomena at hadron colliders [6–9].

Measurement of the missing transverse energy vector ($\mathbf{E}_T^{\text{miss}}$) in events at the LHC will be complicated by the presence of pileup collisions. In CMS, measurement of $\mathbf{E}_T^{\text{miss}}$ will be further degraded by the difference between photon and pion response in the calorimeters and by the bending of tracks by the 4-T magnetic field. On the other hand, the excellent cell segmentation, hermeticity, and good forward coverage of CMS will help the measurement of $\mathbf{E}_T^{\text{miss}}$ [3].

2 Calculation of missing transverse energy in CMS

Readout cells in the CMS hadron calorimeter (HCAL) are arranged in a tower pattern in η, φ space, projective to the nominal interaction point. The HCAL cells have a segmentation of 0.087×0.087 in the central region ($|\eta| < 1.74$) and approximately 0.17×0.17 for $1.74 < |\eta| < 5$. Since the granularity of the electromagnetic crystal calorimeter (ECAL) [10] is much finer (0.017×0.017) than HCAL, calorimeter towers (ECAL plus HCAL) are formed by addition of signals in η, φ bins corresponding to the HCAL geometry. In forming the towers, individual calorimeter cells must pass energy threshold cuts of 500 MeV for HCAL

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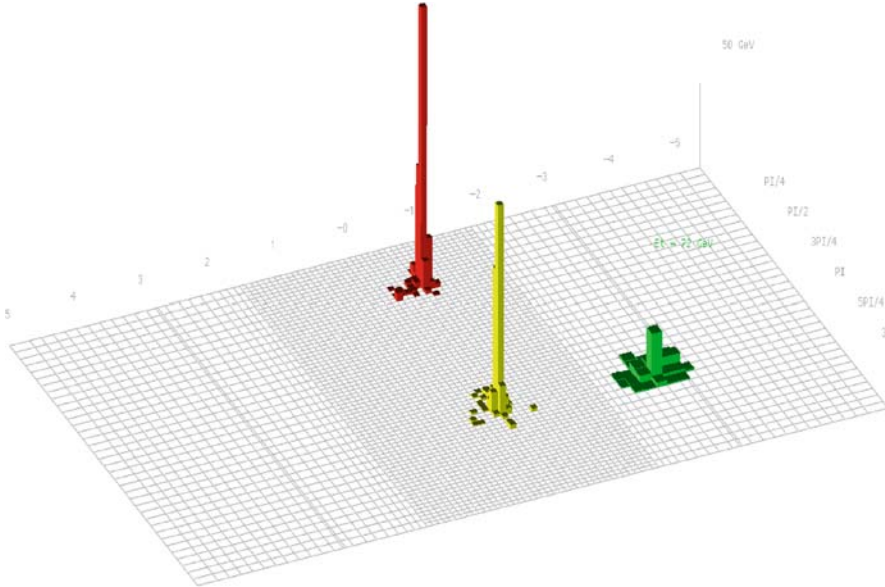


Fig. 1. The η, φ tower segmentation in CMS. The towers are defined to match the granularity of the hadron calorimeter

cells, 90 MeV for ECAL barrel ($|\eta| < 1.48$) and 450 MeV for ECAL endcap crystals ($1.48 < |\eta| < 3.00$). The thresholds correspond to roughly 2–3 standard deviations above the expected root-mean-square electronics noise. In total there are 4176 such towers, which when unfolded, may be represented in a familiar ‘lego’ plot (Fig. 1) [11].

The missing transverse energy vector is calculated by summing individual calorimeter towers having energy E_n , pseudorapidity η_n and azimuthal angle φ_n :

$$\begin{aligned} \mathbf{E}_T^{\text{miss}} &= - \sum \left(\frac{E_n \cos \varphi_n}{\cos h\eta_n} \hat{\mathbf{x}} + \frac{E_n \sin \varphi_n}{\cos h\eta_n} \hat{\mathbf{y}} \right) \\ &= E_x^{\text{miss}} \hat{\mathbf{x}} + E_y^{\text{miss}} \hat{\mathbf{y}}. \end{aligned}$$

Reconstructed muons are taken into account by replacing the expected calorimeter deposit (about 4 GeV) with the reconstructed track p_T . Section 8 discusses further corrections to the reconstructed missing transverse energy.

3 Expected performance in minimum-bias collisions

Accurate measurement of the missing transverse energy in individual events is a difficult experimental problem, because various detector factors are known to contribute in subtle ways. These factors include energy resolution, limited detector coverage, nonlinearity of calorimeter response, detector granularity, non-instrumented material, magnetic field and its effect on low p_T charged particles, quantization of detector readout, electronic noise, event pileup, and underlying event. In spite of all these detector subtleties, the $\mathbf{E}_T^{\text{miss}}$ resolution in CMS is expected to be dominated by calorimeter energy resolution.

The total scalar transverse energy ($\sum E_T$), defined as the scalar E_T sum of all calorimeter towers in an event, is a quantity highly associated with $\mathbf{E}_T^{\text{miss}}$ [12]. Many prop-

erties of reconstructed missing transverse energy can be expressed as a function of $\sum E_T$ because of its direct influence on the $\mathbf{E}_T^{\text{miss}}$ resolution via stochastic effects of calorimeter showers and the process of signal collection. The missing transverse energy resolution is normally determined by fitting the width of the measured E_x^{miss} or E_y^{miss} distribution in a sample of events in which no missing transverse energy is expected [5]. In minimum bias collisions at UA1, the resolution was observed to follow the form $\sigma = 0.4 \sqrt{\sum E_T} \text{GeV}^{1/2}$, where the constant 0.4 scales directly with the coefficient in the stochastic term $0.8 \sqrt{E_T} \text{GeV}^{1/2}$ in the calorimeter resolution [5, 12, 13] due to the contribution of hadronic showers in the measurement of $\mathbf{E}_T^{\text{miss}}$. The CDF experiment, which has a scintillating tile geometry similar to CMS (and a completely different magnetic field configuration compared to UA1, solenoid vs. dipole), observed a missing transverse energy resolution of $\sigma = 0.47 \sqrt{\sum E_T} \text{GeV}^{1/2}$ in Run I [14]. From the UA1 and CDF results and the measured CMS calorimeter resolution stochastic term from test beam of $1.15 \sqrt{E_T} \text{GeV}^{1/2}$, one may expect a missing transverse energy resolution in CMS of $\sigma \approx (0.6 - 0.7) \sqrt{\sum E_T} \text{GeV}^{1/2}$ for minimum bias events with no pileup, when dominated by the shower fluctuations.

4 Event samples

Previous studies of $\mathbf{E}_T^{\text{miss}}$ in CMS [15] were limited by the use of less sophisticated simulation and reconstruction tools. A more advanced understanding of $\mathbf{E}_T^{\text{miss}}$ in the CMS detector has required the large event samples available from the recent data challenge [16]. These Monte Carlo events with full-detector simulation have been used to study the $\mathbf{E}_T^{\text{miss}}$ performance of the CMS detector and to develop correction techniques, as well as to evaluate and optimize the trigger [17].

The samples were generated with PYTHIA [18] using the CMS software package CMKIN 3.0.0 [19], simulated with Geant4 [20] using the CMS software package OSCAR 2.4.5 [21], and digitized with the CMS object-oriented reconstruction code ORCA 7.6.1 [22]. The signal events are combined with low-luminosity ($\mathcal{L} = 2 \times 10^{33} \text{ cm}^{-2} \text{ s}^{-1}$) pileup corresponding to an average of 3.4 fully inelastic collisions per 25 ns beam crossing. The following samples were used in this study: 10^6 minimum bias events, 3.2×10^6 QCD events with parton transverse momenta $0 < \hat{p}_T < 4000 \text{ GeV}/c$, 5×10^5 W +jet events, and 4×10^5 $t\bar{t}$ events. Details about the data samples, generation and simulation parameters, and reconstruction description can be found in [24].

Event reconstruction and physics analysis was performed with ORCA 8.7.1 [22]. No additional thresholds on the calorimeter towers, beyond the noise suppression cuts already mentioned, are applied for jet reconstruction or the $\mathbf{E}_T^{\text{miss}}$ calculation. Muons reconstructed at level-3 (offline) [15] are included in the determination of $\mathbf{E}_T^{\text{miss}}$. No significant difference is observed in the $\mathbf{E}_T^{\text{miss}}$ resolution for the QCD samples between two different muon algorithms used in ORCA [23].

5 Performance in soft collisions

5.1 Minimum bias without pileup

Figure 2a shows the distribution of reconstructed $\sum E_T$ from a high-statistics sample of fully-simulated minimum bias events without pileup. The event simulation includes single and double diffraction. There is a minimum value of reconstructed $\sum E_T$, an offset, that is dependent on the calorimeter cell thresholds, corresponding to an on-line zero-suppression, that are used to form the calorimeter towers. For this sample, the offset is 142 GeV. The average value of 194 GeV corresponds to $\sum E_T$ of 52 GeV above the offset. This 52 GeV of event activity determines the size of the stochastic term in the $\mathbf{E}_T^{\text{miss}}$ resolution for this sample. Figure 2b shows the distribution of E_x^{miss} . The width of this distribution is $\sigma = 6.1 \text{ GeV}$ in agreement with expectation of addition in quadrature of a 3.8 GeV noise term and a $0.63\sqrt{\sum E_T - 142 \text{ GeV}} \text{ GeV}^{1/2}$ stochastic term. It is important to note that at this stage, the vector $\mathbf{E}_T^{\text{miss}} = E_x \hat{x} + E_y \hat{y}$ is made using towers that have been formed from ECAL cells calibrated for photons and HCAL cells calibration for hadrons, appropriate perhaps for understanding the detector response to first collisions. It is believed that making use of energy flow techniques such as the charged track corrections [25] will ultimately improve the $\mathbf{E}_T^{\text{miss}}$ resolution.

5.2 Soft collisions with pileup

The addition of pileup increases the observed $\sum E_T$ as shown in Fig. 3a for a sample of soft QCD collisions with $0 < \hat{p}_T < 15 \text{ GeV}/c$, where the same tower thresholds de-

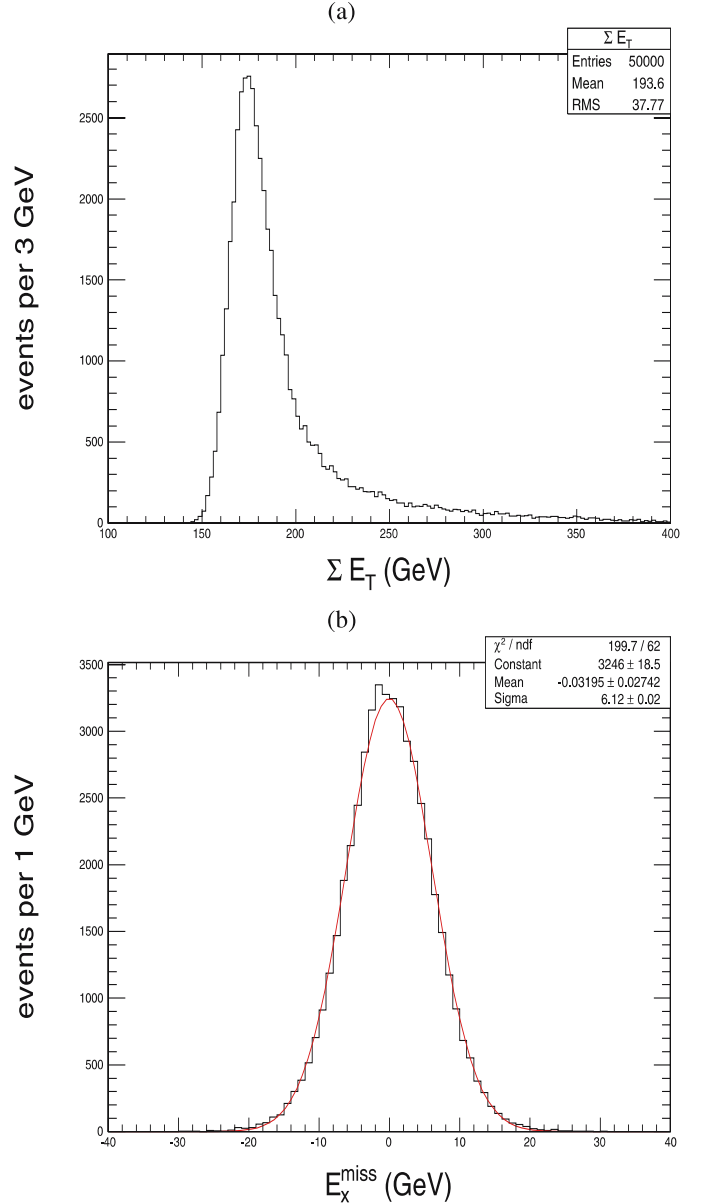


Fig. 2. **a** $\sum E_T$ for minimum bias events with no pileup. The reconstructed offset (minimum value of $\sum E_T$) of 142 GeV depends on calorimeter thresholds. **b** Distribution of E_x^{miss} in the same events. The resolution is fit to a Gaussian with $\sigma = 6.1 \text{ GeV}$, in agreement with expectations based on a stochastic term of $0.63\sqrt{\sum E_T - 142 \text{ GeV}} \text{ GeV}^{1/2}$ and a noise contribution of 3.8 GeV

scribed in Sect. 2 have been applied. The average value of $\sum E_T$ in this sample is 166 GeV greater than that of the minimum bias sample without pileup (Fig. 2a). The corresponding missing transverse energy resolution shown in Fig. 3b is 9.9 GeV. This increase from 6.1 GeV (without pileup) to 9.9 GeV (with pileup) is in agreement with expectations from a stochastic term of $0.63\sqrt{\sum E_T} \text{ GeV}^{1/2}$. Thus, the presence of pileup does not adversely affect the $\mathbf{E}_T^{\text{miss}}$ resolution beyond raising the overall activity ($\sum E_T$) of the event.

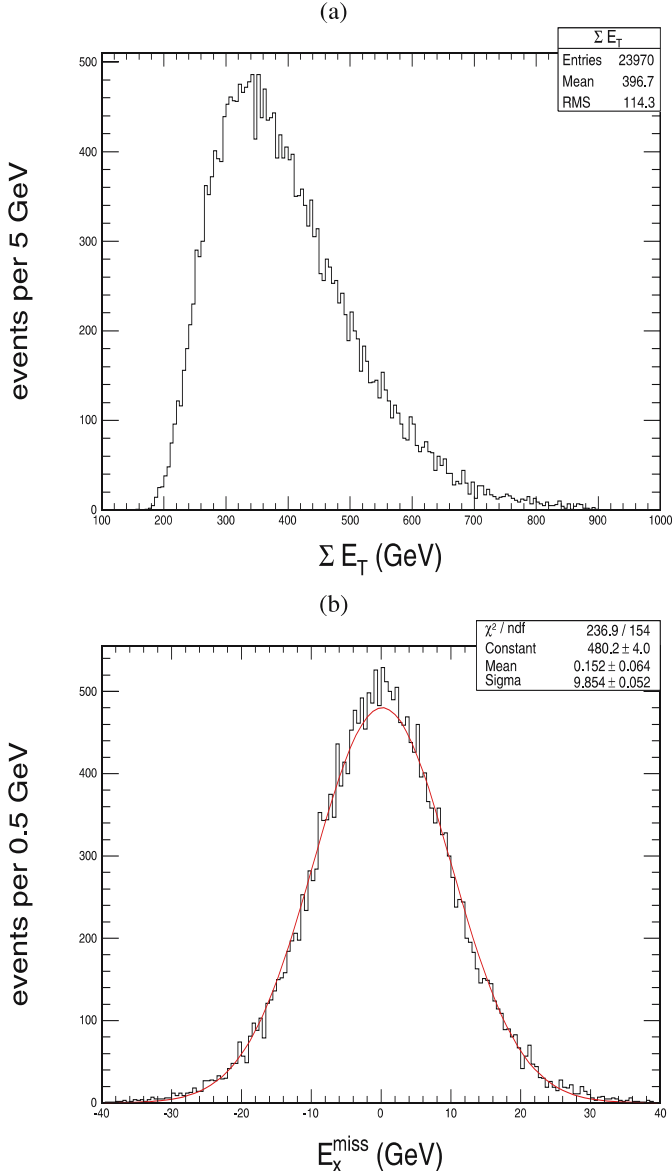


Fig. 3. **a** Distribution of $\sum E_T$ for soft QCD collisions ($0 < \hat{p}_T < 15$ GeV/c), including event pileup corresponding to a luminosity of $\mathcal{L} = 2 \times 10^{33} \text{ cm}^{-2} \text{ s}^{-1}$. **b** Distribution of E_x^{miss} for the same events. The resolution is fit to a Gaussian with $\sigma = 9.9$ GeV

A more detailed comparison of the reconstructed missing transverse energy resolution from minimum bias events and soft QCD events is shown in Fig. 4. The minimum bias and QCD samples of Fig. 4 both include event pileup corresponding to a luminosity of $\mathcal{L} = 2 \times 10^{33} \text{ cm}^{-2} \text{ s}^{-1}$. Figure 4 shows the $\mathbf{E}_T^{\text{miss}}$ resolution, determined from fits to reconstructed E_x^{miss} distributions like that of Fig. 3b, for minimum bias events (open circles) and soft QCD events, $0 < \hat{p}_T < 15$ GeV/c, (solid squares). The resolution as a function of $\sum E_T$ has parameterized with a fit (solid curve in Fig. 4) that gives

$$\sigma = [(4.0 \text{ GeV})^2 + (0.63 \text{ GeV}^{1/2} \sqrt{\sum E_T - 142 \text{ GeV}})^2]^{1/2}.$$

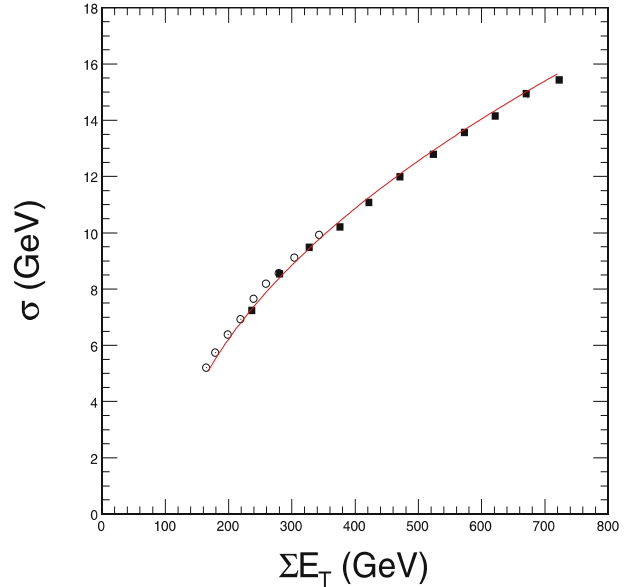


Fig. 4. Missing transverse energy resolution, as determined from Gaussian fits to distributions of reconstructed E_x^{miss} , vs. reconstructed $\sum E_T$ for soft QCD events, $0 < \hat{p}_T < 15$ GeV/c, (solid squares) and minimum bias events (open circles), including event pileup corresponding to a luminosity of $\mathcal{L} = 2 \times 10^{33} \text{ cm}^{-2} \text{ s}^{-1}$. The resolution (curve) has been fit to $\sigma = [(4.0 \text{ GeV})^2 + (0.63 \text{ GeV}^{1/2} \sqrt{\sum E_T - 142 \text{ GeV}})^2]^{1/2}$

The stochastic term is the same as observed in minimum bias collisions with no pileup (Fig. 2b).

6 Performance in hard collisions

6.1 Reconstructed missing transverse energy

For intrinsically balanced 2-to-2 scattering processes, the reconstructed $\mathbf{E}_T^{\text{miss}}$ is directly related to the $\mathbf{E}_T^{\text{miss}}$ resolution [6]. It is well known that missing transverse energy resolution measured in high momentum-transfer events will not fit the same stochastic term as soft collision [12]. This is due in part to fragmentation effects and contributions to the jet resolution from linear terms [26]. The UA1 experiment ($\sqrt{s} = 0.54$ TeV) measured a missing transverse energy, $\mathbf{E}_T^{\text{miss}} = \sqrt{(E_x^{\text{miss}})^2 + (E_y^{\text{miss}})^2}$, in jet events corresponding to a stochastic coefficient of 0.7 [6], while the CDF experiment ($\sqrt{s} = 1.8$ TeV) determined a stochastic coefficient near unity [27]. From these results, we may expect the distribution of $\mathbf{E}_T^{\text{miss}}$ in jet events for CMS to fit a stochastic coefficient in the range 1.1–1.3.

Figure 5 shows (a) the reconstructed $\sum E_T$ and (b) $\mathbf{E}_T^{\text{miss}}$ for QCD events generated with $80 \text{ GeV}/c < \hat{p}_T < 120 \text{ GeV}/c$. The observed $\mathbf{E}_T^{\text{miss}}$ distribution, by its construction, has a one-sided tail. A small \hat{p}_T interval can create a wide spectrum of $\mathbf{E}_T^{\text{miss}}$ at both generator and detector level. The reconstructed $\mathbf{E}_T^{\text{miss}}$ extends to large values (> 50 GeV) due to a small number of events con-

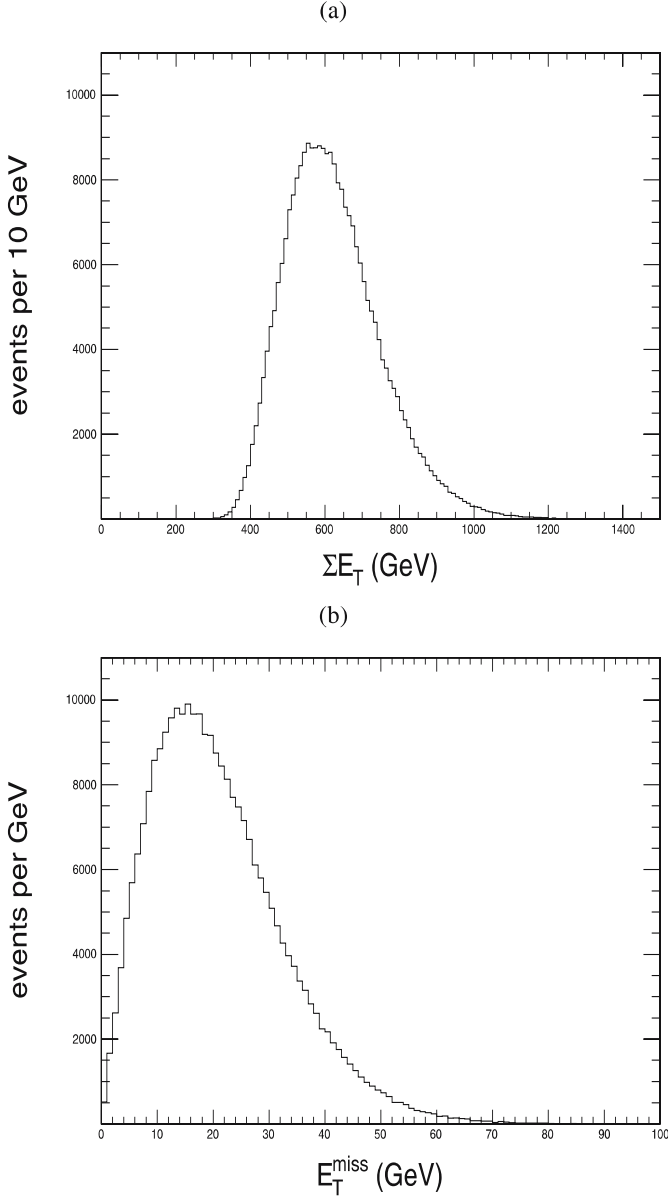


Fig. 5. **a** Reconstructed $\sum E_T$ in a QCD sample with $80 \text{ GeV}/c < \hat{p}_T < 120 \text{ GeV}/c$. The sample includes pileup corresponding to a luminosity of $\mathcal{L} = 2 \times 10^{33} \text{ cm}^{-2}\text{s}^{-1}$. **b** Reconstructed E_T^{miss} in the same sample

taining one or more energetic neutrinos or non-identified muons from heavy flavor (b and c) decays.

The average reconstructed E_T^{miss} in QCD events (with pileup) as a function of $\sum E_T$ is shown in Fig. 6. At $\sum E_T = 2100 \text{ GeV}$, which corresponds to jets with reconstructed $p_T \approx 800 \text{ GeV}/c$, an average E_T^{miss} of about 60 GeV is reconstructed. This number is consistent with measurement of such jets with a resolution of $1.25 \text{ GeV}^{1/2} \sqrt{E_T}$ [28] plus a small contribution from the underlying event. Due to the increased occupancy in these events, the offset from noise contributions in the reconstructed $\sum E_T$ distribution has shifted up to 350 GeV as shown in Fig. 5a. A fit to the reconstructed missing transverse energy (curve

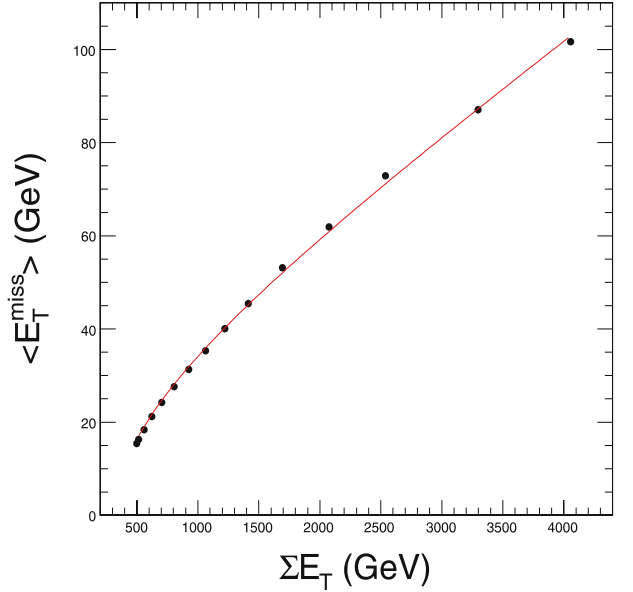


Fig. 6. Average reconstructed E_T^{miss} vs. $\sum E_T$ for QCD hard collisions, including event pileup corresponding to a luminosity of $\mathcal{L} = 2 \times 10^{33} \text{ cm}^{-2}\text{s}^{-1}$. The fit (curve) corresponds to $\langle E_T^{\text{miss}} \rangle = [(5.4 \text{ GeV})^2 + (1.23 \text{ GeV}^{1/2} \sqrt{\sum E_T - 350 \text{ GeV}})^2 + [0.019(\sum E_T - 350 \text{ GeV})]^2]^{1/2}$

in Fig. 6) gives

$$\langle E_T^{\text{miss}} \rangle = \left[(5.4 \text{ GeV})^2 + \left(1.23 \text{ GeV}^{1/2} \sqrt{\sum E_T - 350 \text{ GeV}} \right)^2 + \left[0.019 \left(\sum E_T - 350 \text{ GeV} \right) \right]^2 \right]^{1/2}$$

where the offset in $\sum E_T$ is subtracted as shown. This result is consistent with measurements from UA1 and CDF as noted above.

6.2 Scalar transverse energy sum

The distribution of $\sum E_T$ was investigated for QCD collisions with varying values of generator \hat{p}_T from $20\text{--}1000 \text{ GeV}$, including the effects of pileup corresponding to a luminosity of $\mathcal{L} = 2 \times 10^{33} \text{ cm}^{-2}\text{s}^{-1}$. As noted above, a relatively small \hat{p}_T interval corresponds to a wide range of $\sum E_T$ (see Fig. 5). Figure 7 shows the average reconstructed E_T of the leading jets vs. reconstructed $(\sum E_T)_r$. The points correspond to parton-level \hat{p}_T ranges of (from left to right) of $20\text{--}30$, $30\text{--}50$, $50\text{--}80$, $80\text{--}120$, $120\text{--}170$, $170\text{--}230$, $230\text{--}300$, $300\text{--}380$, $380\text{--}470$, $470\text{--}600$, $600\text{--}800$ and $800\text{--}1000 \text{ GeV}/c$. This allows us to interpret the $\sum E_T$ scale (e.g., Fig. 6) in terms of the dominating jet activity.

Figure 8 shows the ratio of reconstructed $(\sum E_T)_r$ to generated $(\sum E_T)_g$ summed transverse energy vs. $(\sum E_T)_g$. The larger value of $(\sum E_T)_r$ compared to $(\sum E_T)_g$ at small values of $(\sum E_T)_g$ is due to a threshold-dependent offset in the reconstructed $\sum E_T$ scale due to electronic noise.

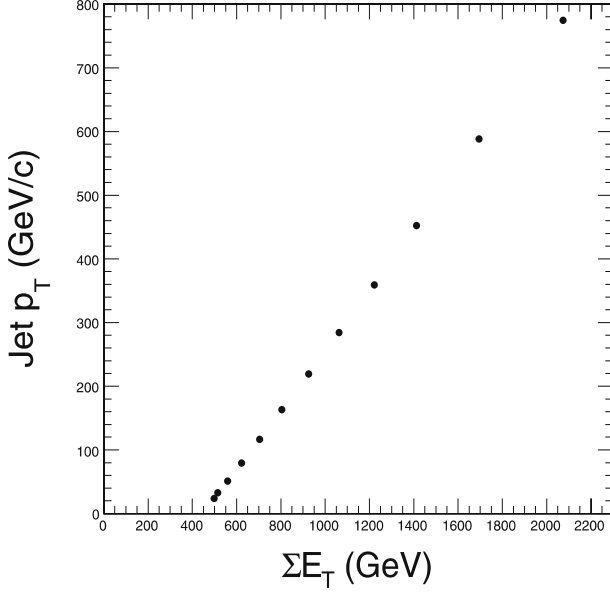


Fig. 7. Average reconstructed p_T of the leading jets vs. reconstructed ΣE_T . The points correspond to parton-level \hat{p}_T ranges of (from left to right) of 20–30, 30–50, 50–80, 80–120, 120–170, 170–230, 230–300, 300–380, 380–470, 470–600, 600–800 and 800–1000 GeV/c

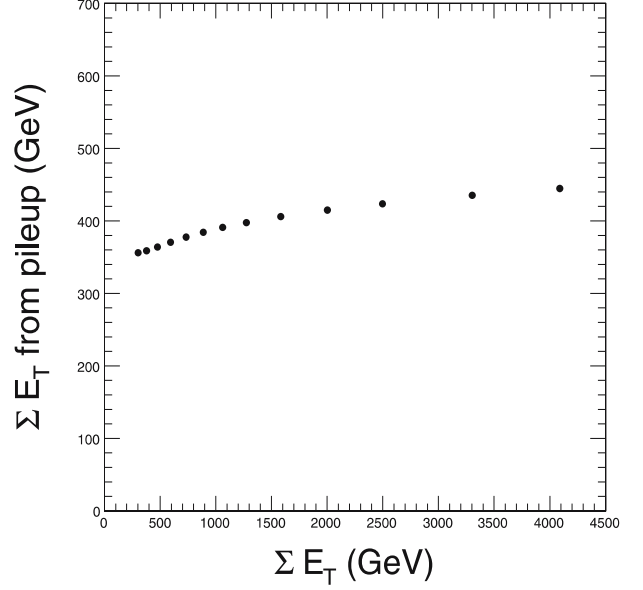


Fig. 9. Contribution of pileup to reconstructed ΣE_T as a function of generated ΣE_T . The slight dependence on generated ΣE_T is due to detector effects

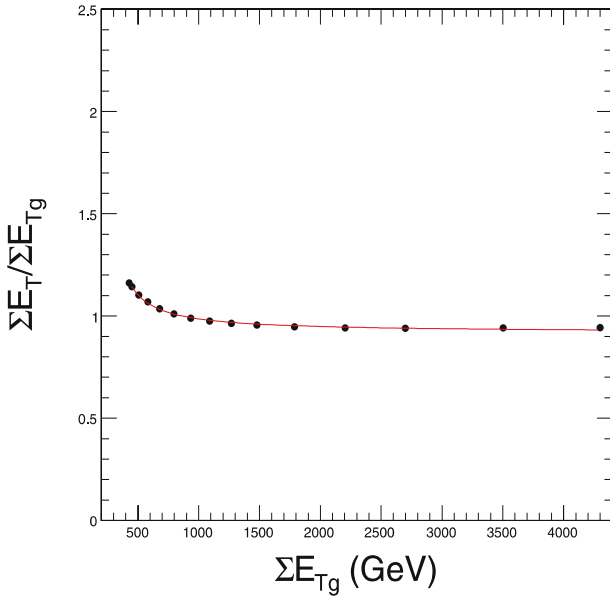


Fig. 8. Ratio of reconstructed $(\Sigma E_T)_r$ to generated $(\Sigma E_T)_g$ summed transverse energy vs. $(\Sigma E_T)_g$ from QCD events with pileup. The increase in the ratio at low values of $(\Sigma E_T)_g$ is due to a threshold-dependent offset in $(\Sigma E_T)_r$ caused by electronic noise. The ratio is fit (curve) to $(\Sigma E_T)_r/(\Sigma E_T)_g = 0.92 + 240 \text{ GeV}/[(\Sigma E_T)_g - 40 \text{ GeV}]$

Figure 9 shows the contribution of pileup to the reconstructed (ΣE_T) vs. generated (ΣE_T) in QCD events. The contribution of pileup to the reconstructed (ΣE_T) is about 350–450 GeV. The dependence of the pileup contribution on the generated (signal) ΣE_T comes from various

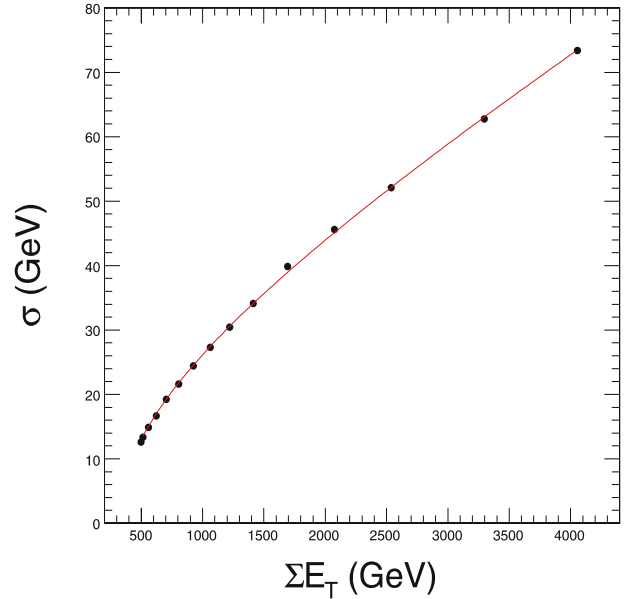


Fig. 10. Missing transverse energy resolution vs. ΣE_T for QCD hard collisions, including event pileup corresponding to a luminosity of $\mathcal{L} = 2 \times 10^{33} \text{ cm}^{-2} \text{ s}^{-1}$. The resolution (curve) has been fit to $\sigma = [(3.8 \text{ GeV})^2 + (0.97 \text{ GeV})^{1/2} \sqrt{\Sigma E_T - 350 \text{ GeV}} + [0.012(\Sigma E_T - 350 \text{ GeV})^2]^{1/2}]^{1/2}$

detector effects including nonlinear response, electronic noise, and the E_T threshold used in calorimeter hit reconstruction (0.5 GeV for HCAL).

6.3 Missing transverse energy resolution

Figure 10 shows the E_T^{miss} resolution, determined from fits to reconstructed E_x^{miss} distributions, for QCD hard colli-

sions, including event pileup corresponding to a luminosity of $\mathcal{L} = 2 \times 10^{33} \text{ cm}^{-2} \text{ s}^{-1}$. At low values of $\sum E_T$, the resolution (Fig. 10) agrees with that determined from soft collisions (Fig. 4). The resolution as a function of $\sum E_T$ is parameterized with a fit (curve in Fig. 10) that gives

$$\sigma = \left[(3.8 \text{ GeV})^2 + \left(0.97 \text{ GeV}^{1/2} \sqrt{\sum E_T - 350 \text{ GeV}} \right)^2 + \left[0.012 \left(\sum E_T - 350 \text{ GeV} \right) \right]^2 \right]^{1/2}.$$

This result is consistent with that measured in UA1 [6] and CDF [27] when scaled by the measured performance of the calorimetry in test beams. The large QCD dijet cross section at the LHC will allow us to measure distributions of E_x^{miss} and E_y^{miss} as a function of instantaneous luminosity in order to directly determine the missing transverse energy resolution under operating conditions.

7 Contribution of jets and unclustered towers

Since jets dominate the E_T contribution in QCD events, it is interesting to investigate out-of-cone radiation and underlying event effects by separating the contributions of reconstructed jets and unclustered towers to the reconstructed missing transverse energy. Two regions are defined in an event: the jet region (defined by the jet cones) and the unclustered region (all calorimeter towers that are outside jets). The value of E_T^{miss} can be calculated separately for each region. The two regions are defined by the reconstructed event at the detector level; at the generator

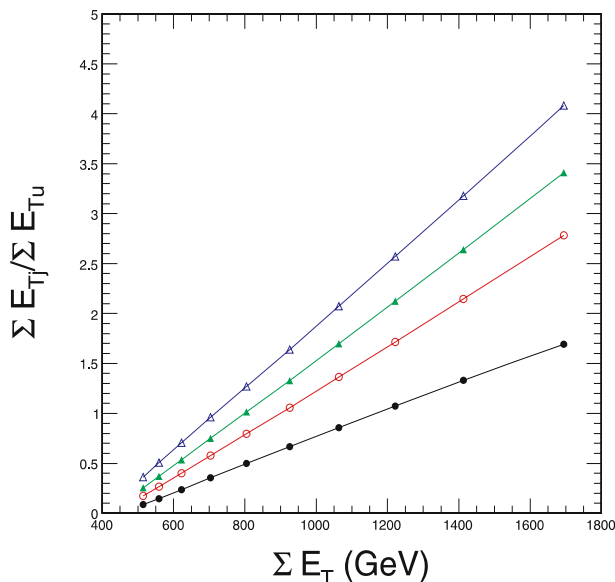


Fig. 11. Ratio of reconstructed $\sum E_T$ in the jet region to that in the unclustered region vs. $\sum E_T$ for jet cones of size $R = 0.2$ (solid circles), 0.4 (open circles), 0.6 (solid triangles), and 0.8 (open triangles)

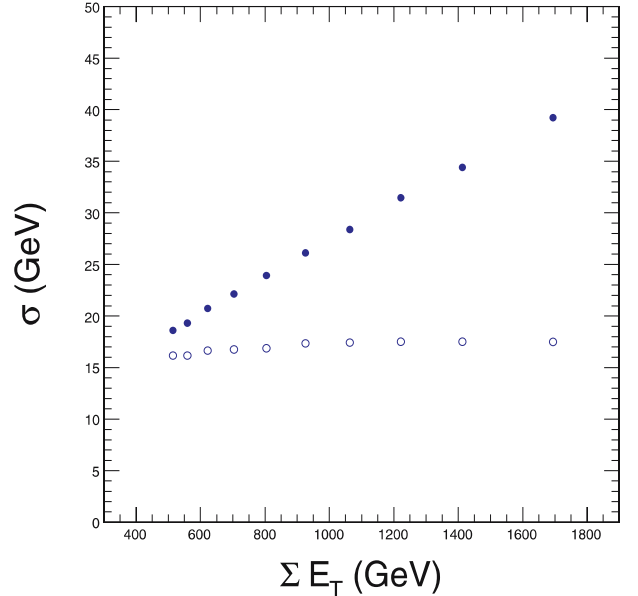


Fig. 12. Contributions of deposits in the jet (solid circles) and unclustered regions (open circles) to the missing transverse energy resolution. vs. reconstructed $\sum E_T$

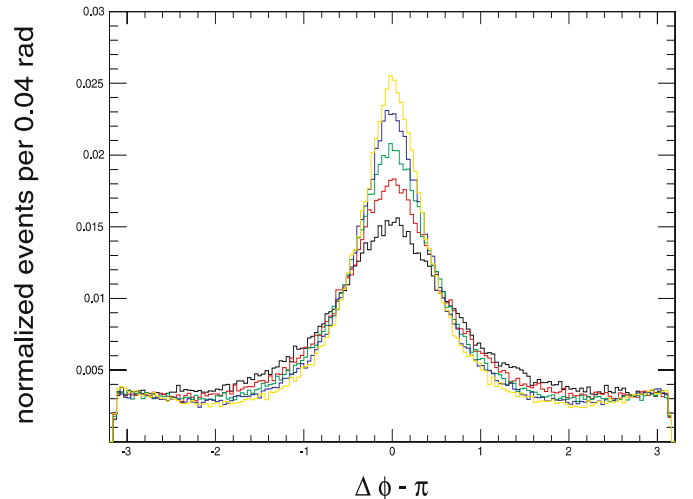


Fig. 13. The φ separation between the direction of the highest E_T jet and the reconstructed vector $\mathbf{E}_T^{\text{miss}}$. Five p_T^{hat} QCD samples are plotted: $50\text{--}80 \text{ GeV}/c$ (black), $80\text{--}120 \text{ GeV}/c$ (red), $120\text{--}170 \text{ GeV}/c$ (green), $170\text{--}230 \text{ GeV}/c$ (blue) and $230\text{--}300 \text{ GeV}/c$ (yellow) GeV/c

level, particles are associated with one of the two regions based on their direction. The definition of regions is performed for every event.

Jets are reconstructed by the iterative cone (IC) algorithm implemented in ORCA with cone size of $R = 0.5$ and a minimum jet threshold of $p_T > 20 \text{ GeV}/c$ [28]. The towers that do not contribute to jets are collected as unclustered towers. The jet cone size was varied to investigate the effect on the region definition and $\mathbf{E}_T^{\text{miss}}$ quantities. We use the axis of IC jets with a $R = 0.5$ cone as input to rebuild the jets with new cone sizes of $R = 0.2, 0.4, 0.6$ and 0.8 . Figure 11

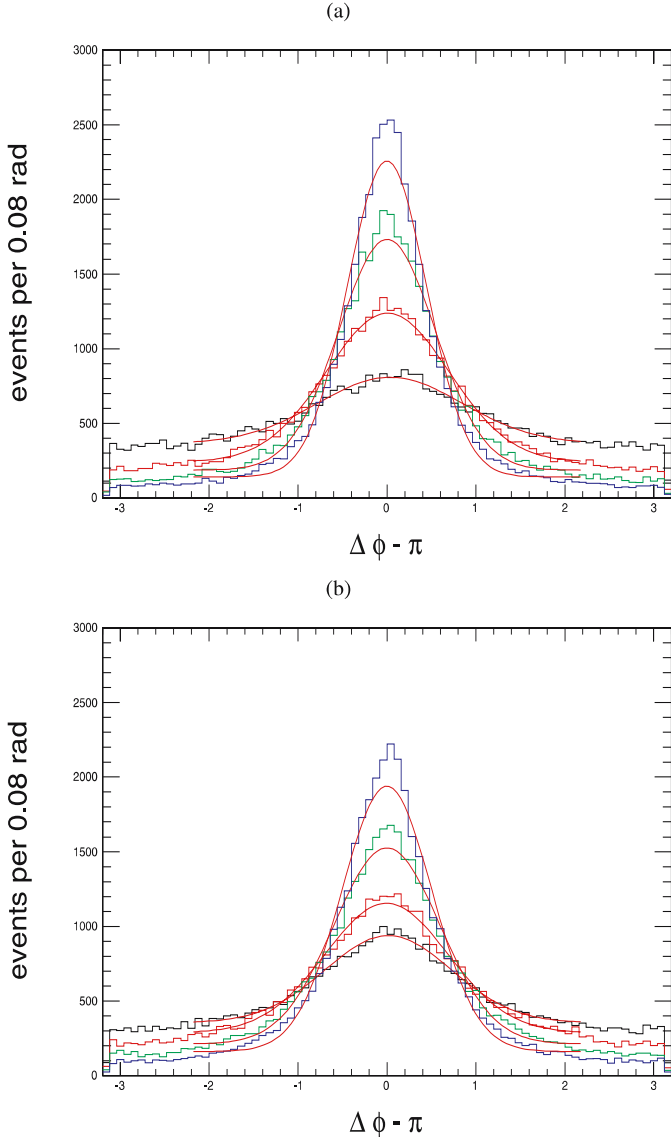


Fig. 14. The φ correlation between the missing transverse energy vectors reconstructed in the jet and unclustered regions for QCD samples with **a** p_T of 30–50 and **b** 50–80 GeV for four cone radii $R = 0.2, 0.4, 0.6$ and 0.8 , from low to high peak height, respectively. The distributions show that the vectors $\mathbf{E}_{Tj}^{\text{miss}}$ and $\mathbf{E}_{Tu}^{\text{miss}}$ are back-to-back

shows the ratio of $\sum E_T$ in the jet region to that in the unclustered region, as a function of reconstructed $\sum E_T$ for the different cone sizes. A cone size of $R = 0.2$ contains only the core of the jet [13]. At $\sum E_T = 1700$ GeV which corresponds to jets with average transverse momenta of approximately 530 GeV/c, about 3/4 of the reconstructed $\sum E_T$ is inside the $R = 0.5$ jet cones. As expected, the jets play a dominant role in defining the reconstructed $\mathbf{E}_T^{\text{miss}}$ in the event. Figure 12 shows the contributions to the missing transverse energy resolution determined separately for the jet and unclustered regions.

For QCD events, it is expected that the direction of reconstructed missing transverse energy be opposite in φ

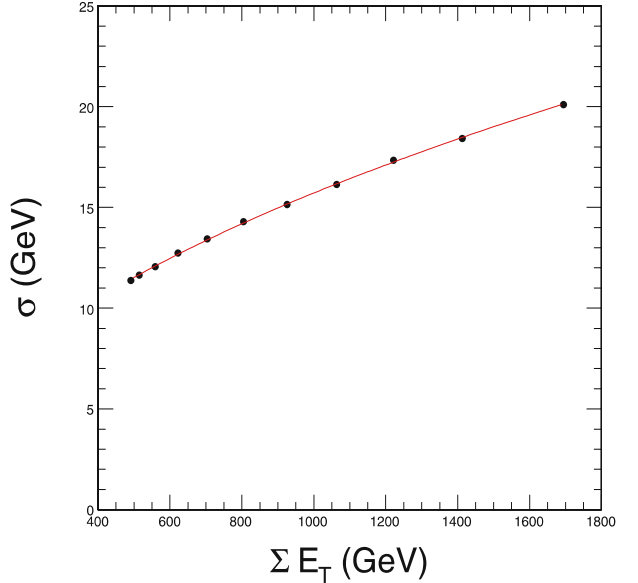


Fig. 15. Missing transverse energy resolution in a direction orthogonal to the jet axis vs. $\sum E_T$ of the entire event. The resolution is comparable to that observed in soft collisions (Fig. 4). A fit (curve) gives $\sigma = [(9.9 \text{ GeV})^2 + (0.48 \text{ GeV}^{1/2} \sqrt{\sum E_T - 350 \text{ GeV}})^2]^{1/2}$

to that of the most energetic jet. Figure 13 shows the φ difference minus π between the highest E_T jet and $\mathbf{E}_T^{\text{miss}}$, indicating that the reconstructed $\mathbf{E}_T^{\text{miss}}$ vector tends to be opposite the leading jet in QCD events.

The direction (φ angle) of reconstructed missing transverse energy in the jet region is strongly correlated to that in the unclustered region indicating that the underlying event contributes to the overall balance of $\mathbf{E}_T^{\text{miss}}$. Figure 14 shows the distribution of $\Delta\varphi - \pi = \varphi_j - \varphi_u - \pi$, where φ_j and φ_u are the azimuthal angles of the reconstructed missing transverse energy vectors in the jet and unclustered regions, for two different QCD samples, (a) $30 \text{ GeV}/c < \hat{p}_T < 50 \text{ GeV}/c$ and (b) $50 \text{ GeV}/c < \hat{p}_T < 80 \text{ GeV}/c$, and four different cone sizes ($R = 0.2, 0.4, 0.6$, and 0.8). The correlation is stronger for lower p_T jets where the importance of the underlying event for overall $\mathbf{E}_T^{\text{miss}}$ balance is greater.

As an additional check of detector $\mathbf{E}_T^{\text{miss}}$ performance, one may look at the resolution in a direction orthogonal to the jet axis. By choosing this direction, the observed resolution is independent of the reconstructed jet resolution, but rather is dominated by the underlying event and pileup activity. This distribution is shown in Fig. 15 vs. reconstructed $\sum E_T$. The result may be fit to the form (curve in Fig. 15)

$$\sigma = \left[(9.9 \text{ GeV})^2 + \left(0.48 \text{ GeV}^{1/2} \sqrt{\sum E_T - 350 \text{ GeV}} \right)^2 \right]^{1/2}.$$

As may be expected, the resolution is comparable (to about 20%) to that observed in soft collisions (Fig. 4).

8 Corrections to E_T^{miss}

Earlier studies [29] show that using a calibrated jet energy scale can restore the average missing transverse energy scale in processes containing genuine E_T^{miss} , but similar methods do not significantly reduce the fake reconstructed E_T^{miss} in QCD events, which mainly comes from the stochastic effect of the calorimeter jet energy response.

Samples of $t\bar{t}$ and W +jet with leptonic decays were used to investigate the reconstruction of events which contain a true E_T^{miss} . These samples were chosen because of their large background contributions to many new physics searches as well as their differing event topologies. Events having isolated muons or electrons with transverse momentum ($p_T > 20$ GeV) and pseudorapidity ($|\eta| < 3.0$ for electrons and $|\eta| < 2.4$ for muons) were selected after full detector reconstruction. Jets are reconstructed using the iterative cone algorithm with a cone size $R = 0.5$, seed threshold $E_T = 1.0$ GeV, minimum jet threshold $E_T = 15$ GeV, and calorimeter tower threshold $E_T = 0.5$ GeV. Muons were reconstructed using the level 3 muon code from ORCA [15],[22]. Figure 16 shows the resulting inclusive jet p_T spectrum from QCD (solid line), W +jet with $W p_T$ between 40 and 300 GeV (dashed line), and $t\bar{t}$ (dotted line). Figure 17 shows the reconstructed E_T^{miss} and $\sum E_T$ spectrum in $t\bar{t}$ and W +jet events.

Corrections were made to the jet energy scale using the technique described in [28]. The main purpose of the correction, which restores the average reconstructed jet E_T in a cone to the average generated E_T in the cone, is to account for the nonlinearity of the calorimetry. The

size of the correction, which depends on η , is 10%–20% at jet $p_T = 350$ GeV/c. In basing the jet corrections on QCD events, which is done in order to suppress fake E_T^{miss} from badly reconstructed jets, we introduce a p_T -dependent bias into the $t\bar{t}$ and W +jet samples due to their different par-

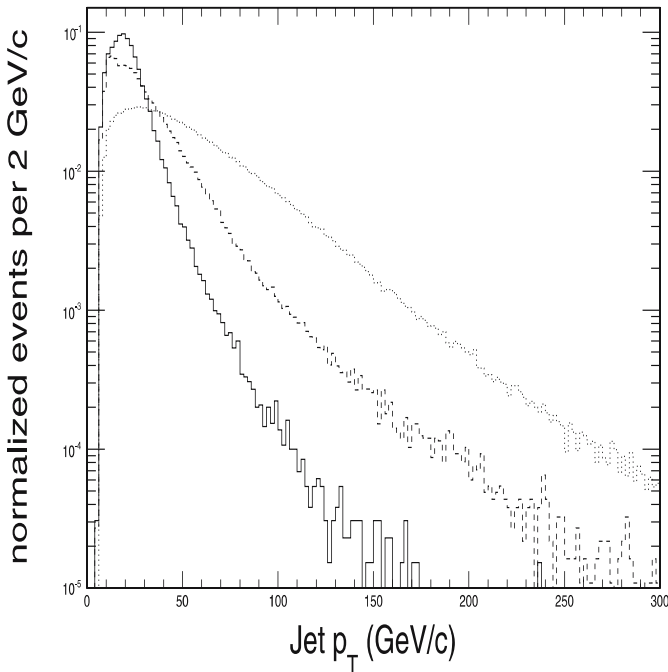


Fig. 16. Inclusive jet p_T spectrum for events having an isolated lepton with $p_T > 20$ GeV from QCD (solid line), W +jet with $W p_T$ between 40 and 300 GeV (dashed line), and $t\bar{t}$ (dotted line)

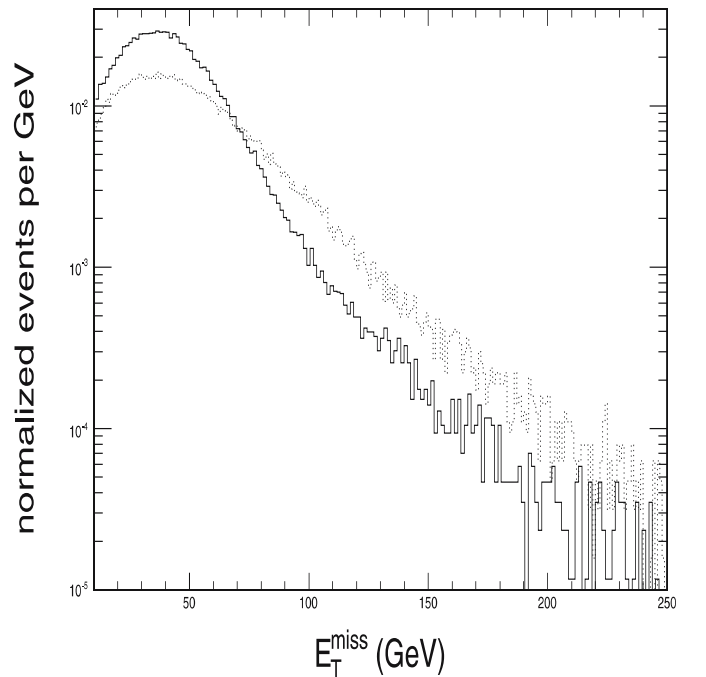
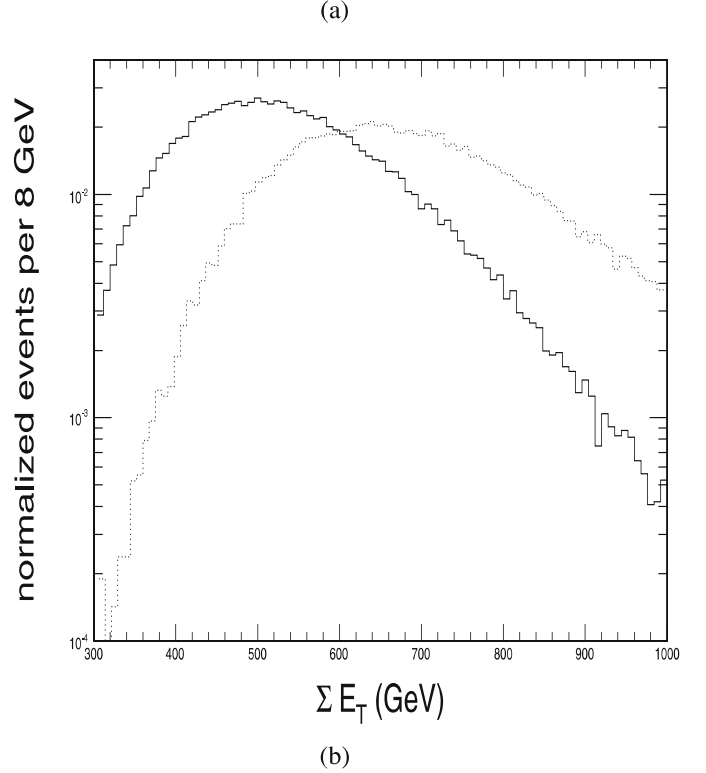


Fig. 17. Reconstructed **a** $\sum E_T$ and **b** E_T^{miss} in $t\bar{t}$ events (dashed line) and W +jet events with $W p_T$ between 40 and 300 GeV (solid line)

ton composition. The correction factors are calculated as a function of jet p_T and η with QCD events. The value of E_T^{miss} is calculated as before with the towers corresponding to the jets having been corrected. Figure 18 shows the missing transverse energy resolution vs. reconstructed E_T^{miss} before (open circles) and after (closed circles) jet correc-

tions for (a) $t\bar{t}$ events and (b) W +jet events. It is evident that the missing transverse energy resolution is better for the W +jet events, relative to $t\bar{t}$ events, at low values of E_T^{miss} where there is a significant contribution to the value of $\sum E_T$ from a well-measured electron from W decay in the electromagnetic calorimeter, an effect also observed in

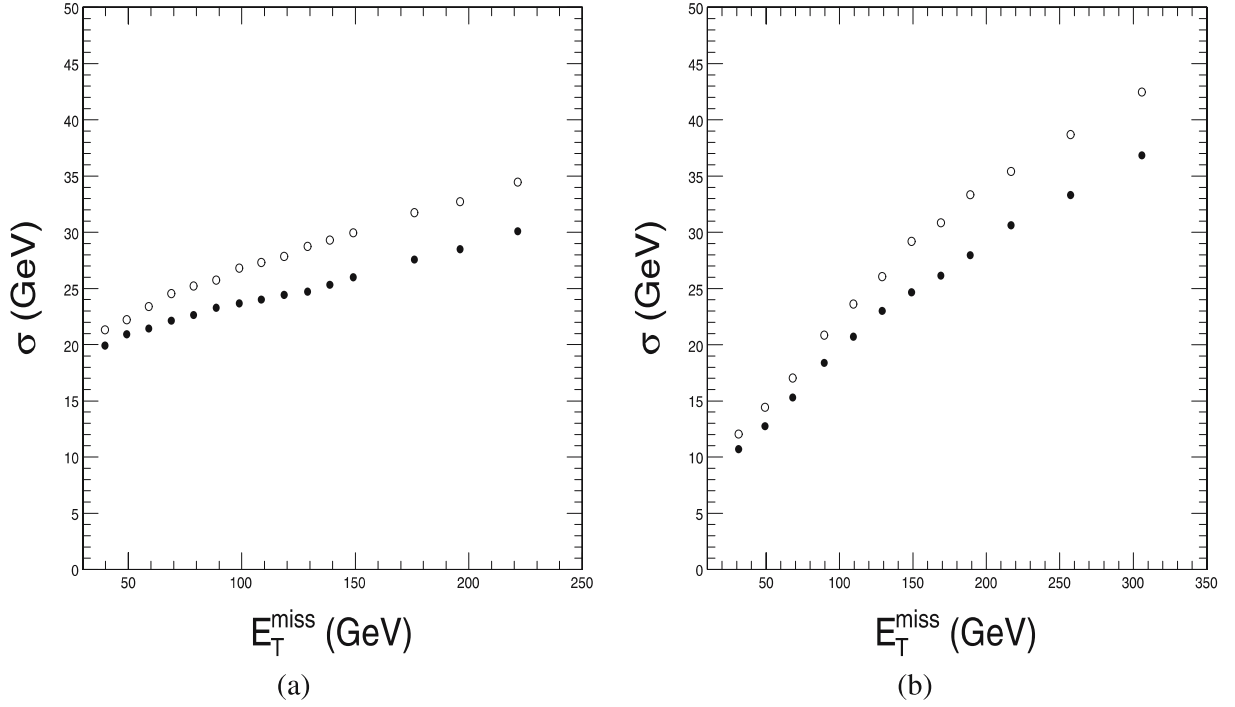


Fig. 18. Missing transverse energy resolution before (*open circles*) and after (*closed circles*) jet corrections vs. reconstructed E_T^{miss} for **a** $t\bar{t}$ events and **b** W +jet events

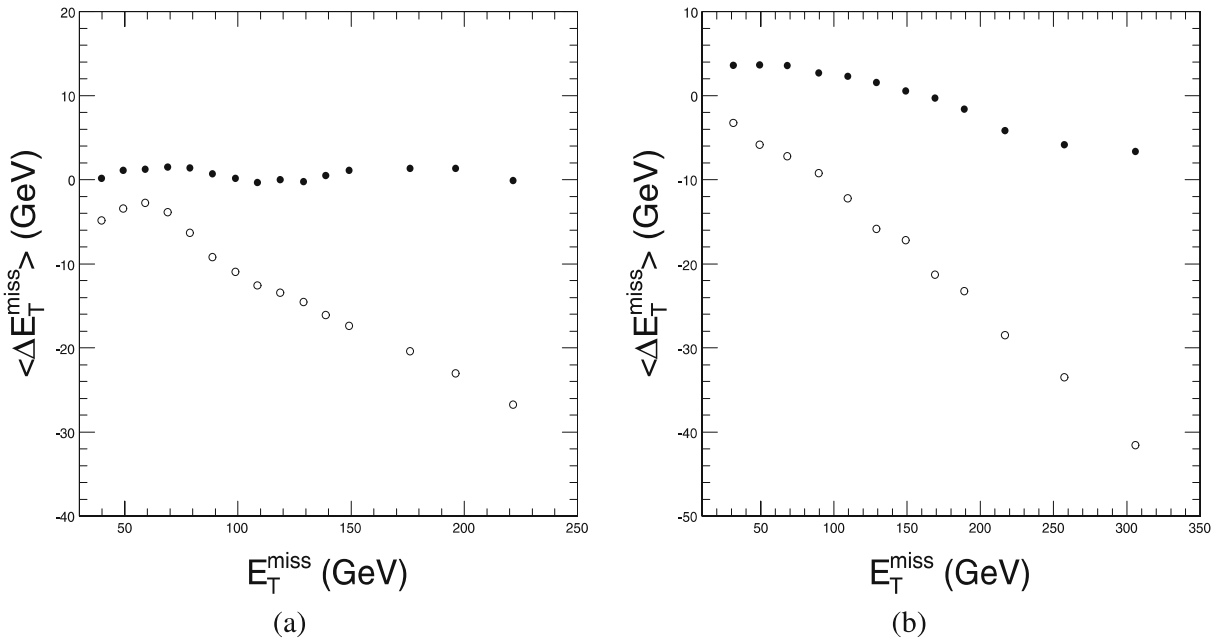


Fig. 19. Error in the reconstructed missing transverse energy scale vs. reconstructed E_T^{miss} before (*open circles*) and after (*closed circles*) jet corrections for **a** $t\bar{t}$ events and **b** W +jet events

UA1 [5]. At large values of E_T^{miss} , the W +jet events have worse resolution due to the W recoiling against a high E_T jet. Figure 19 shows the error in the reconstructed missing transverse energy scale (the average value of reconstructed E_T^{miss} minus true E_T^{miss}) vs. reconstructed E_T^{miss} before (open circles) and after (closed circles) jet corrections for (a) $t\bar{t}$ events and (b) W +jet events.

9 Angular resolution

To establish a signature for missing transverse energy in an event, it is likely necessary to show that the vector E_T^{miss} does not have the same φ as an energetic jet whose energy fluctuation in the calorimeter may well have been the cause of the reconstructed E_T^{miss} . Figure 20a shows the E_T^{miss} angular resolution as a function of reconstructed E_T^{miss} in $t\bar{t}$ events with leptonic decays having $p_T^\ell > 20$ GeV before (open circles) and after (closed circles) Monte Carlo corrections to the jet energies. The resolution is calculated by fitting the distribution $\varphi_r - \varphi_g$, where φ_r (φ_g) is the azimuthal angle of the reconstructed (generated) missing transverse energy vector. The angular resolution for measured $E_T^{\text{miss}} = 45$ GeV is comparable to the jet cone size (0.5 radian) but improves at $E_T^{\text{miss}} = 100$ GeV to the jet core size (0.2 radian) [13]. At larger values of E_T^{miss} , the resolution approaches that of the hadronic tower size (0.1 radian). The Monte Carlo corrections of the missing transverse energy vector make only a small improvement, approximately 0.02 radians for reconstructed $E_T^{\text{miss}} = 100$ GeV, to the angular resolution.

The missing transverse energy angular resolution has also been studied for W +jet events with leptonic decays

having $p_T^\ell > 20$ GeV, which have a different topology than $t\bar{t}$ events. In these events, the Wp_T is approximately twice the average p_T of an energetic neutrino which causes a non-zero reconstructed missing transverse energy. In this case, the missing transverse energy tends to be opposite in φ to an energetic jet whose mismeasurement contributes directly to the reconstructed E_T^{miss} . Figure 20b shows the missing transverse energy φ resolution as a function of reconstructed E_T^{miss} , in agreement with the angular resolution determined from $t\bar{t}$ events (Fig. 20a).

10 Summary

The missing transverse energy performance of the CMS detector has been evaluated using fully simulated minimum bias, QCD dijet, $t\bar{t}$, and W +jet events with high statistics. For minimum bias events without pileup, a resolution of 6.1 GeV is computed, corresponding to a stochastic contribution of $0.63\sqrt{\sum E_T}$ GeV $^{1/2}$, where $\sum E_T$ is the summed transverse energy in all calorimeter towers. This result is in agreement with expectations from measurements at UA1 and CDF, taking into account differences in calorimeter resolution. When the contribution of pileup is included, the resolution degrades according to the overall deposited $\sum E_T$ with the same stochastic coefficient. For QCD events, we expect $\sigma = [(3.8 \text{ GeV})^2 + (0.97 \text{ GeV}^{1/2}\sqrt{\sum E_T})^2 + (0.012 \sum E_T)^2]^{1/2}$ resulting in a resolution of 45 GeV for jet events with reconstructed transverse momentum of 800 GeV/c, again in accord with measurements made at UA1 and CDF. A sample of $t\bar{t}$ events with lepton decays leading to true missing transverse energy was used to determine the azimuthal angle

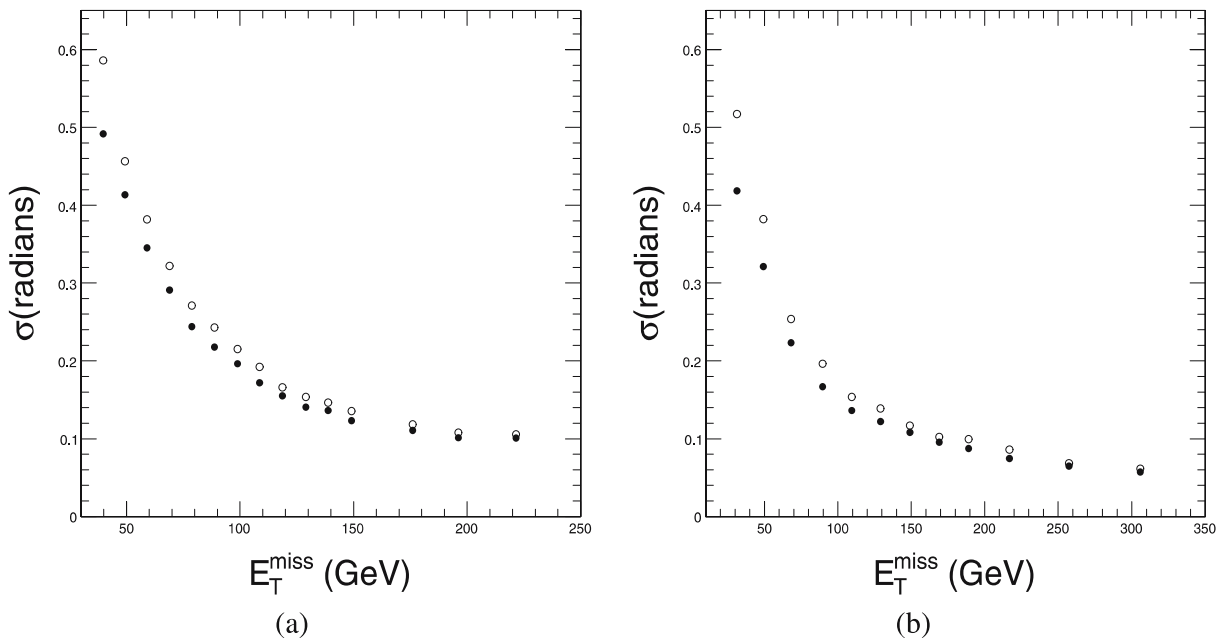


Fig. 20. Phi resolution of the vector E_T^{miss} as a function of reconstructed E_T^{miss} in events with leptonic decays having $p_T^\ell > 20$ GeV for **a** $t\bar{t}$ and **b** W +jet events before (open circles) and after (closed circles) Monte Carlo corrections to the jet energies

resolution to be 0.1 radians (0.2 radians) for a reconstructed missing transverse energy of 200 GeV (100 GeV).

Acknowledgements. It is a pleasure to thank Salavat Abdullin, Daniel Denegri, Gigi Rolandi, Maria Spiropulu, and Roberto Tenchini for many useful suggestions. We also wish to acknowledge the support of the CMS HCAL/JetMET working group.

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