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Highly dispersed networks generated by enhanced redirection

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Abstract. We analyze growing networks that are built by enhanced redirection. Nodes are sequentially added and each incoming node attaches to a randomly chosen 'target' node with probability 1-r, or to the parent of the target node with probability r. When the redirection probability r is an increasing function of the degree of the parent node, with $r \to 1$ as the parent degree diverges, networks grown via this enhanced redirection mechanism exhibit unusual properties, including (i) multiple macrohubs, i.e., nodes with degrees proportional to the number of network nodes N; (ii) non-extensivity of the degree distribution in which the number of nodes of degree k, N_k , scales as $N^{\nu-1}/k^{\nu}$, with $1 < \nu < 2$; (iii) lack of self-averaging, with large fluctuations between individual network realizations. These features are robust and continue to hold when the incoming node has out-degree greater than 1 so that networks contain closed loops. The latter networks are strongly clustered; for the specific case of double attachment, the average local clustering coefficient is $\langle C_i \rangle = 4 \ln 2 - 2 = 0.77258...$

Keywords: growth processes, network dynamics, random graphs, networks, clustering techniques

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1. Introduction

Models for the growth of complex networks often involve mechanisms that are based on global knowledge of the network. For example, in preferential attachment [1]–[7], nodes are added sequentially and each links to existing target nodes in the network according to an attachment rate A_k that is an increasing function of the degree k of the target node. According to this rule, incoming nodes must 'know' the degree distribution of the entire network to correctly choose a target node. In real networks, however, it is not feasible that any new node has such detailed global knowledge.

The impracticality of implementing a growth rule based on global knowledge has motivated alternatives to preferential attachment that rely on the incoming nodes exploiting only local knowledge of a small portion of the network. Examples include attachment via spatial locality [8]–[10] and node similarity [11]. In this work, we focus on the local growth rule that exploits redirection [12]–[18]. Here, each incoming node selects a target node at random and links either to this target node (probability 1-r) or to the parent of the target node (probability r). This redirection rule is based on the network being directed so that the parent(s) of any node is well defined. If each new

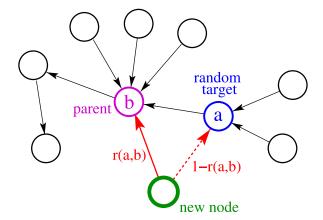


Figure 1. Illustration of enhanced redirection. The redirection probability r(a,b) depends on the degree a of a randomly chosen target node and the degree b of its parent, with r(a,b) an increasing function of b and $r \to 1$ as $b \to \infty$. In this example, the degree of the parent node is large, so the new node is likely to attach to the parent.

node has only one outgoing link, redirection produces networks with a tree topology; it is straightforward to extend redirection to allow each incoming node to attach to more than one node in the network [15].

The surprising feature of redirection with a fixed redirection probability r is that it is mathematically equivalent to the global growth rule of shifted linear preferential attachment, where the rate of attaching a new node to a pre-existing node of degree k is $A_k = k + \lambda$ with $\lambda = r^{-1} - 2$, see [13]. Redirection is also highly efficient because one only needs to select a random node and identify its parent to add a node to the network. The time to create a network of N nodes via redirection therefore scales linearly with N.

The utility of redirection as a simple and efficient procedure that is equivalent to linear preferential attachment motivates us to exploit models that use slightly more comprehensive (but still local) degree information around the target node. Specifically, we allow the redirection probability r(a,b) to depend on the degrees of the target and parent nodes, a and b respectively (figure 1). In hindered redirection, r(a,b) is a decreasing function of the parent degree b, a rule that leads to sub-linear preferential attachment growth [19]. In this work, we investigate the complementary situation of enhanced redirection, for which the redirection probability r is an increasing function of the parent degree b, with $r \to 1$ as $b \to \infty$. This seemingly-innocuous redirection rule gives rise to networks with several intriguing and practically relevant properties.

(i) Appearance of multiple macrohubs. Macrohubs are nodes whose degrees are a finite fraction of N. While macrohubs arise in other models [4, 13], [20]– [22], the resulting networks are singular, with nearly all nodes attached to a single macrohub. In the cases of superlinear preferential attachment, where $A_k \sim k^{\gamma}$ with $\gamma > 1$ [4, 13, 20], and in the fitness model, where the attachment rate is proportional to both the degree k and the fitness of the target [22, 23], a single macrohub arises that is connected to almost all other nodes of the network. In contrast, enhanced redirection networks

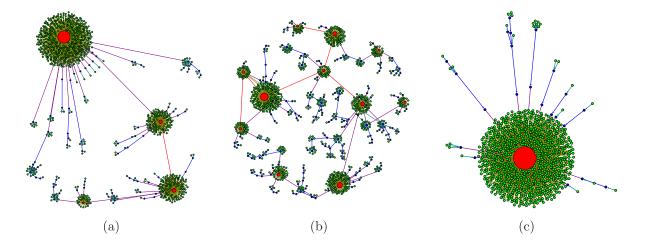


Figure 2. Enhanced redirection networks of $N=10^3$ nodes for $\lambda=\frac{3}{4}$ (see equation (3)) starting from the same initial state. (a) Maximum degree $k_{\rm max}=548$, $\mathcal{C}=66$ core $(k\geq 2)$ nodes, and maximum depth $D_{\rm max}=10$. (b) $k_{\rm max}=\mathcal{C}=154$, $D_{\rm max}=12$ (smallest $k_{\rm max}$ out of 10^3 realizations). (c) $k_{\rm max}=963$, with $\mathcal{C}=23$ and $D_{\rm max}=6$ (largest $k_{\rm max}$ out of 10^3 realizations). Green, nodes of degree 1; blue, degrees 2–20; red, degree > 20. The link color is the average of the endpoint node colors.

are highly disperse (figure 2), with interconnected hub-and-spoke structures that are reminiscent of airline route networks [7, 21], [24]–[26].

(ii) Non-extensivity. In many sparse networks, the degree distribution is extensive, with the number of nodes of degree k, N_k , proportional to N. This happens, for example, in linear preferential attachment where additionally the degree distribution has an algebraic tail, $N_k \sim N/k^{\nu}$ for $k \gg 1$, with $\nu > 2$. In contrast, enhanced redirection leads to the non-extensive scaling

$$N_k \sim \frac{N^{\nu-1}}{k^{\nu}} \qquad \text{with } \nu < 2.$$
 (1)

The allowed range of the exponent ν is key. While past empirical studies have observed networks with degree exponent in the range $1 < \nu < 2$ [27] (see table 1 for some examples), the range $1 < \nu < 2$ is mathematically inconsistent for sparse networks because it leads to a divergent average degree as $N \to \infty$ whenever the degree distribution obeys the standard scaling $N_k \sim N k^{-\nu}$. A simple resolution of this dilemma is to relax the hypothesis of extensivity. We shall see that in enhanced redirection almost all nodes have degree 1 (leaves). More precisely, the number $\mathcal{C} \equiv N - N_1$ of core nodes (nodes with degree > 1) grows sub-linearly with N, namely as $N^{\nu-1}$ with the exponent ν in the range $1 < \nu < 2$. All N_k with $k \ge 2$ also grow as $N^{\nu-1}$. This anomalous scaling can therefore be summarized as follows:

$$C = N - N_1 \simeq c_1 N^{\nu - 1}, \qquad N_k \simeq c_k N^{\nu - 1} \quad \text{for } k \ge 2,$$
 (2)

Table 1. Networks with degree exponent $\nu < 2$. All examples are simple graphs (at most one link between any node pair); the first 10 are undirected and the remainder are directed. Data at http://konect.uni-koblenz.de/networks.

Network	Exponent ν	Average degree	Network size
Orkut	1.27	76.281	3072 441
Catster friendships	1.36	72.803	149700
Dogster friendships	1.40	40.048	426820
arXiv hep-ph	1.47	224.14	28093
arXiv hep-th	1.47	213.44	22908
Wikipedia conflict	1.50	34.644	118100
Hamsterster full	1.52	13.711	16630
Hamsterster friendships	1.54	13.491	1858
Flickr	1.73	43.742	105938
Internet topology	1.86	9.8618	34761
Wikipedia, Italian	1.48	28.457	1204009
Wikipedia, German	1.50	28.811	2166669
LiveJournal	1.56	28.465	4847571
Wikipedia, French	1.62	22.165	2212682
OpenFlights	1.79	20.756	2939

where the c_k are constants. This scaling satisfies the sum rule $\sum_{1 \leq k \leq N} N_k = N$ and leads to a finite average degree $\langle k \rangle$ without imposing an artificial cutoff in the degree distribution.

(iii) Lack of self-averaging. Different realizations of enhanced redirection are visually diverse when starting from the same initial condition (figure 2). Basic network measures, such as the number of nodes of fixed degree, N_k with any $k \geq 2$, or the number of core nodes \mathcal{C} , vary significantly between realizations and do not converge as $N \to \infty$. For instance, the ratio of the mean deviation to the average, $\sqrt{\langle \mathcal{C}^2 \rangle - \langle \mathcal{C} \rangle^2}/\langle \mathcal{C} \rangle$, converges to a positive constant when $N \to \infty$, thereby manifesting the lack of self-averaging. In contrast, preferential attachment networks do self-average, as the relative deviations in N_k or \mathcal{C} systematically decrease as N increases [28].

In the next section we formally define our enhanced redirection models. In section 3 we provide analytical and numerical arguments that justify the properties (i), (ii), (iii) given above. Some of these arguments substantially extend our findings that were reported in [29]. Most results in section 4, and all results in section 5 about enhanced redirection with multiple attachments, are new.

2. Enhanced redirection model

We define the initial network to be a single node that is linked to itself, so that the root node is its own parent and its own child. The initial conditions have a weak and mostly

quantitative influence on asymptotic network properties. Thus we shall generally use the above simple initial condition; we will explicitly define other initial conditions in the few cases where such a modification is more amenable to analysis.

Links are directed so that the parent and children of any node are well defined. In sections 3 and 4 we investigate models in which each node has out-degree equal to 1, and thus a unique parent. This growth rule produces tree networks if the starting network is a tree. Our networks are trees with the exception of the initial self-loop.

Nodes are introduced one by one. Each incoming node first picks a random target node. If the degrees of the target and parent nodes are a and b, respectively, then the new node (see figure 1)

- (i) attaches to the target with probability 1 r(a, b); or
- (ii) attaches to the parent of the target with probability r(a, b).

Two natural (but by no means unique) choices for the redirection probability are

$$r(a,b) = 1 - b^{-\lambda}, \qquad r(a,b) = \frac{a^{\lambda}}{a^{\lambda} + b^{\lambda}}, \qquad \lambda > 0.$$
 (3)

Our results are robust with respect to the form of the redirection probability, as long as $r(a,b) \to 1$ as $b \to \infty$. For concreteness, we focus on the redirection probability $r(a,b) = 1 - b^{-\lambda}$. In appendix A we compare some results for this case with corresponding results for the redirection probability $r(a,b) = a^{\lambda}/(a^{\lambda} + b^{\lambda})$. This comparison indicates that the two models are qualitatively the same.

3. Degree distribution and lack of self-averaging

We first study the degree distribution $\langle N_k \rangle$, the average number of nodes of degree k; to avoid notational clutter we drop the angle brackets henceforth. Simulation results clearly show that the degree distribution has the anomalous scaling behaviors given by equation (2) (figure 3). The exponent ν depends on the redirection parameter λ , but is always less than 2 (figure 4) so that the degree distribution decays very slowly in k. Because $\nu < 2$, equation (2) implies that the number of nodes of degree 1 grows more rapidly with N than the number of core nodes. Thus, visually, a typical network is dominated by its leaves.

We employ the master equation approach to understand the anomalous scaling in equation (2). The degree distribution evolves according to

$$\frac{\mathrm{d}N_k}{\mathrm{d}N} = \frac{(1 - f_{k-1})N_{k-1} - (1 - f_k)N_k}{N} + \frac{(k-2)t_{k-1}N_{k-1} - (k-1)t_kN_k}{N} + \delta_{k,1}.$$
 (4)

Here, f_k and t_k are defined as the respective probabilities that an incoming link is redirected from a node of degree k, and redirected to a node of degree k. The first ratio in equation (4) accounts for instances in which the incoming node attaches directly to the target node. Thus the term $(1 - f_k)N_k/N$ gives the probability that the randomly selected target node has degree k and the incoming node is not redirected. The term is

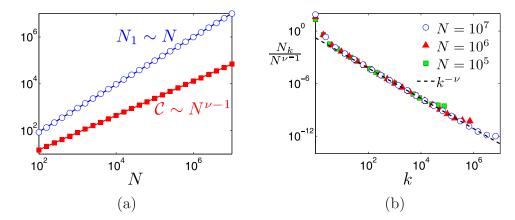


Figure 3. (a) N_k versus N and (b) $N_k/N^{\nu-1}$ versus k for enhanced redirection with $\lambda = \frac{3}{4}$ and $\nu = 1.73$ (determined numerically; see figure 4). Data are based on 10^4 realizations, with equally-spaced bins on a logarithmic scale in (b). The lines in (a) show the prediction of equation (2), while the line in (b) shows the k dependence from the numerical solution of (8).

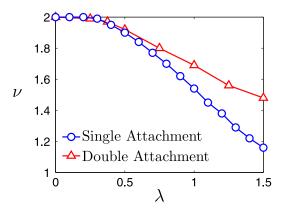


Figure 4. Degree distribution exponent ν versus λ for enhanced redirection with single attachment (\circ) and double attachment (Δ). Each data point is determined from fits of N_k versus N, as in figures 3(a) and 11.

negative because the target node degree increases from k to k+1, which causes N_k to decrease. Similarly, the second ratio corresponds to instances in which the incoming node is redirected to the parent. Thus the term $(k-1)t_kN_k/N$ gives the probability that one of the $(k-1)N_k$ children of a degree k node is targeted and that the incoming node is redirected to the parent. The term $\delta_{k,1}$ arises because each newly added node has degree 1.

The probabilities f_k and t_k are defined by

$$f_k = \sum_{b \ge 1} \frac{r(k, b)N(k, b)}{N_k}, \qquad t_k = \sum_{a \ge 1} \frac{r(a, k)N(a, k)}{(k - 1)N_k}.$$
 (5)

Here, the correlation function N(a, b) is defined as the number of nodes of degree a that have a parent of degree b. Thus f_k is the probability of redirecting from a node of degree

k, averaged over all such target nodes, and t_k is the probability of redirecting to a node of degree k, averaged over all the $(k-1)N_k$ children of nodes of degree k. Defining $\alpha_k = (k-1)t_k + 1 - f_k$, equation (4) can be written in the canonical form

$$\frac{\mathrm{d}N_k}{\mathrm{d}N} = \frac{\alpha_{k-1}N_{k-1} - \alpha_k N_k}{N} + \delta_{k,1}.\tag{6}$$

Substituting equation (2) into the master equations (4) gives the recursions

$$(\nu - 1)c_1 N^{\nu - 2} = \alpha_1 (1 - c_1 N^{\nu - 2}) \qquad k = 1$$

$$(\nu - 1)c_2 N^{\nu - 2} = \alpha_1 (1 - c_1 N^{\nu - 2}) - \alpha_2 c_2 N^{\nu - 2} \qquad k = 2$$

$$c_k = \frac{\alpha_{k-1}}{\alpha_k + \nu - 1} c_{k-1} \qquad k \ge 3.$$
(7)

We eliminate the common factor in the first two lines to obtain $c_2 = \nu/(\alpha_2 + \nu)$, which, combining with the recursion for $k \geq 3$, gives the product solution

$$c_k = c_1 \frac{\nu - 1}{\alpha_k} \prod_{j=2}^k \left(\frac{\alpha_j}{\alpha_j + \nu - 1} \right). \tag{8}$$

For an explicit solution, we need the analytic form for α_k , which requires the probabilities f_k and t_k . For redirection probability $r(a,b) = 1 - b^{-\lambda}$, the quantities f_k and t_k reduce to

$$f_k = \sum_{b>1} \frac{(1-b^{-\lambda})N(k,b)}{N_k} \equiv 1 - \langle b^{-\lambda} \rangle, \qquad t_k = \sum_{a>1} \frac{(1-k^{-\lambda})N(a,k)}{(k-1)N_k} = 1 - k^{-\lambda}, \tag{9}$$

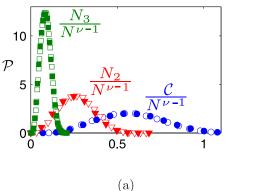
where we use the sum rule $\sum_{a\geq 1} N(a,k) = (k-1)N_k$. We now combine equation (9) with $\alpha_k = (k-1)t_k + 1 - f_k$, to give $\alpha_k = k - k^{1-\lambda} + k^{-\lambda} - f_k \to k$ in the large-k limit. Finally, using $\alpha_k \sim k$ in the product solution (8) gives the asymptotic behavior

$$c_k \sim c_1 \frac{\nu - 1}{k} \prod_{j=2}^k \left(\frac{j}{j + \nu - 1} \right) \sim k^{-\nu}.$$
 (10)

Thus the degree distribution exhibits anomalous scaling, $N_k \sim N^{\nu-1}/k^{\nu}$, with $\nu < 2$, as given in (1).

Numerical simulations show that the exponent ν is a decreasing function of λ (figure 4). For $\lambda \to 0$, enhanced redirection becomes equivalent to random attachment, for which the degree distribution is extensive, $N_k \sim N$ and $\nu \to 2$. As λ increases, attachment to a single node becomes progressively more likely and $\nu \to 0$. We also checked that the exponent ν is not affected by different initial conditions such as an initial loop of different sizes. However, finer details of the degree distribution, such as the probability distribution for the maximal degree and the number of core nodes, do depend on the initial condition.

One of the visually striking features of enhanced redirection networks is that they display large fluctuations from realization to realization, as is apparent from the examples in figure 2. To quantify these fluctuations, let us study $P(N_k)$, the distributions of the number of nodes of fixed degree k. For networks that are grown by preferential attachment,



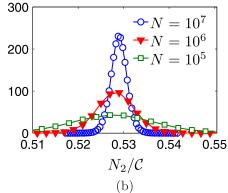


Figure 5. Probability densities in enhanced redirection for (a) $\mathcal{C}/N^{\nu-1}$, $N_2/N^{\nu-1}$, and $N_3/N^{\nu-1}$ for $N=10^6$ (open symbols) and $N=10^7$ nodes (closed symbols) and (b) N_2/\mathcal{C} . Data are based on 10^5 realizations with $\lambda=\frac{3}{4}$ and $\nu=1.73$.

this distribution becomes progressively sharper as N increases [23], as long as the degree is not close to its maximal value. Thus the *average* number of nodes of a given degree can be regarded as the set of variables that fully characterizes the degree distribution. It is only the nodes of the highest degree that fail to self-average [30].

In contrast, for enhanced redirection networks, essentially all geometrical features are non self-averaging, as illustrated by the distributions of $C/N^{\nu-1}$, $N_2/N^{\nu-1}$, $N_3/N^{\nu-1}$, etc, which do not sharpen as N increases (figure 5). Since the number of core nodes C and the number of nodes of fixed degree N_k for $k \geq 2$ both scale as $N^{\nu-1}$ (equation (39)), scaled distributions of $N_k/N^{\nu-1}$ and $C/N^{\nu-1}$ would progressively sharpen as N increases if self-averaging holds. The lack of self-averaging implies a sensitive dependence on initial conditions where events early in the evolution have lasting effects on the network structure. Surprisingly, the ratios N_k/C are self-averaging for $k \geq 2$, as the distributions N_k/C do sharpen as N increases (figure 5). The self-averaging of these ratios suggests that the degree distributions given a value of C are statistically similar, even though the overall number of core nodes C varies widely between realizations.

4. Singular structures

Because of the tendency to connect to high-degree nodes, enhanced redirection networks tend to be dominated by one or a few high-degree nodes. In this section, we explore some of the consequences of this attraction to high-degree nodes.

4.1. Macrohubs

Macrohubs always arise when $\lambda > 0$, but they are easily detectable only when the exponent ν is notably smaller than 2. Figure 4 shows that this happens when $\lambda \gtrsim 0.4$, and in this range macrohubs are clearly observed in all network realizations. When the redirection parameter is small, $0 < \lambda < 0.4$, it may be necessary to grow the network to an astronomically large value of N to detect macrohubs with certainty.

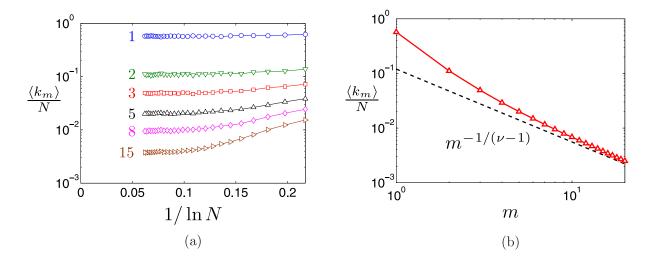


Figure 6. (a) The averages $\langle k_m \rangle / N$ versus $1/\ln N$ for selected m, illustrating that these values approach a finite fraction as $N \to \infty$. Each data point corresponds to 10^3 realizations. (b) $\langle k_m \rangle / N$ versus m for $N = 10^6$.

There are usually many macrohubs, whose degree is proportional to N, as shown in figure 6(a). To estimate k_m , the degree of the mth largest macrohub, we use the extremal criterion [31],

$$\sum_{k \ge k_m} N_k \sim m. \tag{11}$$

This equation merely states that there should be of the order of m nodes of degree k_m or larger. Thus k_m indeed gives an estimate for the value of the mth largest degree. Combining this criterion with the asymptotic $N_k \sim N^{\nu-1}/k^{\nu}$ from equation (1) gives

$$k_m \sim N/m^{1/(\nu-1)}. (12)$$

The basic feature is that the degrees of macrohubs scale linearly with N. In contrast, for networks with an extensive degree distribution of the form $N_k \sim N/k^{\nu}$ and $\nu > 2$, the above extremal criterion gives sub-linear growth: $k_m \sim (N/m)^{1/(\nu-1)}$.

The degrees of the macrohubs substantially depend on the early stages of network growth, but once a set of macrohubs emerges (with degrees k_1, k_2, k_3, \ldots), the probability of attaching to a macrohub of degree k_m asymptotically approaches k_m/N . This preferential attachment to macrohubs is similar to a Pólya urn process for filling an urn with balls of several colors [32, 33]. In the urn process, a ball is drawn at random from the urn and then replaced along with an additional matching-color ball. The different colors in the urn process correspond to different macrohubs in enhanced redirection. If there are k_m balls of the mth color in an urn of N total balls, then the probability of choosing color m, and thus increasing k_m , is given by k_m/N . For the Pólya urn process, the ultimate fractions of balls of different colors do not self-average; the same is expected for the scaled degrees of macrohubs in enhanced redirection.

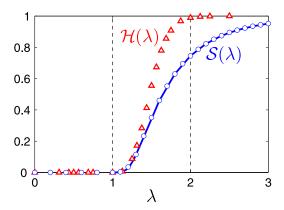


Figure 7. Probabilities for a star (\circ) and a hairball graph (Δ), \mathcal{S} and \mathcal{H} , respectively. For each λ , the data are based on 10^5 realizations. The solid curve represents the numerical evaluation of the product in (13).

4.2. Star graphs

Because of the tendency to link to high-degree nodes, it is possible that a star graph arises in which a single node is connected to every other node of the network. As we now show, the probability for such a star to occur is non-zero for $\lambda > 1$. For the initial condition of a single node with a self-loop, the star contains N-1 leaves, while the root node has degree N+1. The probability S_N to build such a star graph is

$$S_N(\lambda) = \prod_{n=1}^{N-1} \left\{ \frac{1}{n} + \frac{n-1}{n} \left[1 - (n+1)^{-\lambda} \right] \right\}.$$
 (13)

The factor 1/n accounts for the new node attaching to the root in a network of n nodes, while the second term accounts for first choosing a leaf and then redirecting to the root. As shown in appendix A, the asymptotic behavior of (13) is

$$S_N(\lambda) \to \begin{cases} S(\lambda) & \lambda > 1\\ A/N & \lambda = 1 \end{cases}$$

$$\exp\left(-\frac{N^{1-\lambda}}{1-\lambda}\right) \qquad 0 < \lambda < 1$$

$$\frac{1}{(N-1)!} \qquad \lambda = 0,$$
(14)

where $S(\lambda) = \lim_{N \to \infty} S_N(\lambda)$ is a monotonically increasing function of λ when $\lambda > 1$ and $A = \pi^{-1} \sinh \pi \approx 3.676$.

For $0 < \lambda \le 1$, the probability of a star graph asymptotically approaches zero as $N \to \infty$. In this range, the network typically has many macrohubs with average sizes distributed according to equation (12). For $\lambda > 1$, a star graph occurs with a positive probability (figure 7), with a continuous phase transition at $\lambda = 1$. As shown in appendix A, this phase transition has an infinite order because all derivatives of $\mathcal{S}(\lambda)$ vanish at $\lambda = 1$.

4.3. Size distributions of macrohubs

We now study the size distribution of the largest macrohub, the second largest macrohub, etc. By 'size' we mean the degree of a macrohub, so that there is no confusion between the size of the largest macrohub (a quantity characterizing one node) and the degree distribution which specifies the number of nodes of a fixed degree.

The degree k_m of the *m*th largest macrohub scales linearly with the total number of nodes N, and therefore the corresponding size distribution $M_m(k_m, N)$ approaches the scaling form

$$M_m(k_m, N) \to \frac{1}{N} \mathcal{M}_m(x), \qquad x = \frac{k_m}{N},$$
 (15)

in the $N \to \infty$ limit. We do not know how to compute the scaling functions $\mathcal{M}_m(x)$, but some generic properties of these functions can be established without calculations. For instance, the scaling function $\mathcal{M}_m(x)$ vanishes when x > 1/m. Indeed, the mth largest macrohub has maximal degree $k_m = N/m$, which corresponds to the situation when the first m largest macrohubs all have equal maximal possible size N/m. Thus $\mathcal{M}_m(x)$ is singular at x = 1/m: $\mathcal{M}_m(x) \equiv 0$ when x > 1/m, and $\mathcal{M}_m(x) > 0$ when x < 1/m. Consider now the most interesting function $\mathcal{M}_1(x)$, which describes the scaled degree of the largest macrohub. It has a singularity at x = 1, and also a singularity at x = 1/2, as at this point the second largest macrohub can emerge. Continuing this line of reasoning, we conclude that $\mathcal{M}_1(x)$ has infinitely many singularities that are located at x = 1, $1/2, 1/3, 1/4, \ldots$. The emergence of these progressively weaker singularities is a generic feature and they arise in numerous examples including random walks, random maps, spin glasses, fragmentation, etc, which are characterized by the lack of self-averaging, see, e.g., [34]–[38]. Similarly, the scaling function $\mathcal{M}_m(x)$ has singularities at x = 1/m, $1/(m+1), 1/(m+2), \ldots$

The presence of infinitely many singularities (partly) explains why it is difficult to compute the scaling functions $\mathcal{M}_m(x)$. Fortunately, it is possible to probe the asymptotic behavior of $\mathcal{M}_m(x)$ near the maximal possible size x = 1/m. Consider the most interesting case of the largest macrohub. To determine the asymptotic behavior of $\mathcal{M}_1(x)$ in the $x \to 1$ limit, we notice that it can be extracted from the probability to build a star graph. Comparing (14) with (15) we find that in the marginal case of $\lambda = 1$

$$\mathcal{M}_1(1) = A = \frac{\sinh(\pi)}{\pi},\tag{16}$$

while in the $0 < \lambda < 1$ range, the scaling function $\mathcal{M}_1(x)$ very rapidly vanishes near the upper limit,

$$\ln \mathcal{M}_1(x) \sim -(1-x)^{-(1-\lambda)} \quad \text{when } x \to 1.$$
 (17)

The asymptotic behavior of $\mathcal{M}_2(x)$ in the $x \to \frac{1}{2}$ limit can be similarly extracted from the probability to build a 2-star graph. We outline some of these calculations in the following subsection.

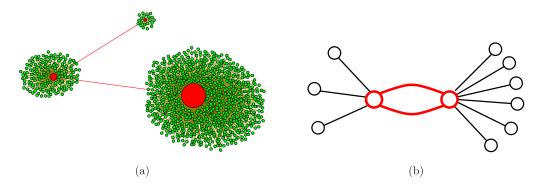


Figure 8. (a) Example hairball network of $N = 10^3$ nodes generated by enhanced redirection with $\lambda = 2.5$. There are three core nodes (degree ≥ 2) with degrees k = 819, 163, and 21. (b) Idealized hairball graph $H_{m,n}$, with (m,n) = (3,6). The initial nodes and links are highlighted.

4.4. Hairballs

A slightly less singular variant of the star graph is what we term the 'hairball' graph (figure 8). A hairball consists of multiple linked stars in which there are no nodes of degree 2. A star is thus a special case of a hairball that consists of a single ball. The star probability is therefore always less than the hairball probability $\mathcal{H}(\lambda)$ (see also figure 7), and the latter appears to reach 1 for $\lambda > 2$. Thus enhanced redirection networks undergo two distinct phase transitions: (i) emergence of a star graph when $\lambda > 1$ and (ii) the vanishing of nodes of degree 2 when $\lambda > 2$.

It is difficult to determine the hairball probability $\mathcal{H}(\lambda)$ analytically, as the number of high-degree nodes and their degrees are unspecified. To gain insight, we consider the more tractable probabilities for concrete types of hairballs. Let us start with the simplest hairball that contains two macrohubs. We also modify the initial condition to make the calculations cleaner; namely, two nodes in a cycle of size two (figure 8(b)). A hairball with two macrohubs is thus a network where all nodes, apart from the two initial nodes, are leaves. Suppose that the network has reached the stage when one initial node is connected to m leaves and the other initial node is connected to n leaves. Let $H_{m,n}$ be the probability to reach such an (m, n) hairball. This hairball can arise from an (m-1, n) or an (m, n-1) hairball, which (by definition) occur with probabilities $H_{m-1,n}$ and $H_{m,n-1}$. Generally,

$$H_{m,n} = g_{m-1,n} H_{m-1,n} + g_{n-1,m} H_{m,n-1}, (18)$$

with coefficients g that depend on the redirection rule. For redirection probability $r(a, b) = 1 - b^{-\lambda}$,

$$g_{m-1,n-1} = \frac{1}{m+n} \left[m - \frac{m}{(m+1)^{\lambda}} + \frac{1}{(n+1)^{\lambda}} \right].$$
 (19)

Instead of simulating enhanced redirection networks and looking for hairballs, we can use the recurrence (18) to calculate the exact values $H_{m,n}$ for any (m,n) starting from the obvious initial condition $H_{0,0} = 1$.

The recurrences (18) are readily solvable in the limit when either m or n vanishes. In this case

$$H_{N-2,0} = H_{0,N-2} = \frac{1}{2} \prod_{n=1}^{N-2} \left[1 - \frac{n-1}{n^{\lambda+1}} \right]$$
 (20)

for $N \geq 3$. Hence the probability to generate the star graph, $S_N = H_{N-2,0} + H_{0,N-2}$, is

$$S_N(\lambda) = \prod_{n=1}^{N-2} \left[1 - \frac{n-1}{n^{\lambda+1}} \right].$$

Using the same analysis as that used to derive equation (14), we find the asymptotic behaviors

$$S_N(\lambda) \to \begin{cases} S(\lambda) & \lambda > 1, \\ B/N & \lambda = 1, \\ \exp\left(-\frac{N^{1-\lambda}}{1-\lambda}\right) & 0 < \lambda < 1, \end{cases}$$
 (21)

with $S(\lambda) = \prod_{n \ge 1} [1 - (n-1)/n^{\lambda+1}]$ and

$$B = \prod_{n=1}^{\infty} \left[1 + \frac{1}{n(n+1)} \right] = \frac{1}{\pi} \cosh\left(\frac{\pi\sqrt{3}}{2}\right) = 2.428189792\dots$$

The differences between (14) and the above formulae stem from the different initial conditions.

The most interesting behavior arises when both m and n are large and comparable: $m \sim n \sim N$. In this regime we employ a continuum approach. We treat m and n as continuous variables and expand $H_{m-1,n}$ and $H_{m,n-1}$ in Taylor series to lowest order,

$$H_{m-1,n} = H - \frac{\partial H}{\partial m}, \qquad H_{m,n-1} = H - \frac{\partial H}{\partial n},$$
 (22)

where $H \equiv H_{m,n}$. Substituting the expansions (22) into (18) and using (19), we recast the original recurrence into a partial differential equation that depends on λ . When $0 < \lambda < 1$, we obtain

$$m\frac{\partial H}{\partial m} + n\frac{\partial H}{\partial n} = -(m^{1-\lambda} + n^{1-\lambda})H. \tag{23}$$

The controlling factor of the solution is given by

$$H_{m,n} \sim \exp\left[-\frac{m^{1-\lambda} + n^{1-\lambda}}{1-\lambda}\right], \qquad 0 < \lambda < 1.$$
 (24)

To find the sub-leading factors, it would be necessary to refine (23) by keeping lower-order terms.

In the marginal case $\lambda = 1$, the partial differential equation becomes

$$m\frac{\partial H}{\partial m} + n\frac{\partial H}{\partial n} = -3H,\tag{25}$$

whose solution, which satisfies the necessary symmetry requirement $H_{m,n} = H_{n,m}$, is

$$H_{m,n} = \frac{C}{(mn)^{3/2}} \tag{26}$$

with some amplitude C that cannot be computed in the framework of the continuum approximation.

Setting m = n = N/2 (so that the corresponding total number of nodes in the network is N+2) we see that the probability of such an (N/2, N/2) graph scales as N^{-3} . This is precisely the probability that the second largest macrohub has the maximal possible size N/2. The scaling behavior (15) of the size distribution of the second largest hub is compatible with the N^{-3} extremal behavior if

$$\mathcal{M}_2(x) \sim \left(\frac{1}{2} - x\right)^2 \quad \text{when } x \to \frac{1}{2}.$$
 (27)

For $\lambda > 1$, the partial differential equation becomes

$$m\frac{\partial H}{\partial m} + n\frac{\partial H}{\partial n} = -H. \tag{28}$$

The remarkable feature of this equation is its universality (independence of λ) for large m and n. Solving (28) we get

$$H_{m,n} = \frac{C_2(\lambda)}{\sqrt{mn}}, \qquad \lambda > 1. \tag{29}$$

The multiplicative constant factor $C_2(\lambda)$ cannot be determined within the continuum framework. Setting again m = n = N/2 we find that the second largest macrohub has the maximal possible size N/2 with probability $2C_2(\lambda)/N$, which in conjunction with the scaling behavior (15) tells us that

$$\mathcal{M}_2\left(\frac{1}{2}\right) = 2C_2(\lambda) \quad \text{when } \lambda > 1.$$
 (30)

Let us briefly discuss the general case of a hairball with p macrohubs. To simplify the analysis we again modify the initial condition by taking the initial network to be a cycle of p nodes. Generalizing the above analysis, we find that in the marginal case of $\lambda = 1$, the governing equation is

$$m_1 \frac{\partial H}{\partial m_1} + \dots + m_p \frac{\partial H}{\partial m_p} = -(2p-1)H,$$
 (31)

whose solution is

$$H_{m_1,\dots,m_p} = C_p \left(\prod_{j=1}^p m_j \right)^{-2+1/p}.$$
 (32)

When $\lambda > 1$, the governing equation for H is

$$m_1 \frac{\partial H}{\partial m_1} + \dots + m_p \frac{\partial H}{\partial m_p} = -(p-1)H,$$
 (33)

whose solution is

$$H_{m_1,\dots,m_p} = C_p(\lambda) \left(\prod_{j=1}^p m_j\right)^{-1+1/p}.$$
 (34)

Using these results we extract the asymptotic behavior of the scaled size distribution $\mathcal{M}_p(x)$ of the pth largest macrohub in the $x \to 1/p$ limit,

$$\mathcal{M}_{p}(x) \sim \begin{cases} \left(\frac{1}{p} - x\right)^{2p-2} & \lambda = 1, \\ \left(\frac{1}{p} - x\right)^{p-2} & \lambda > 1. \end{cases}$$
 (35)

Finally, when $\lambda < 1$, the asymptotic behavior of the scaled size distribution $\mathcal{M}_p(x)$ is extracted from (24), and its generalization to an arbitrary p, to give

$$\ln \mathcal{M}_p(x) \sim -\left(\frac{1}{p} - x\right)^{-(1-\lambda)}.$$
(36)

4.5. Root node

To further appreciate the role of macrohubs let us now consider the evolution of the degree of the root node. We return to our default initial condition of a single root node that is linked to itself. Let R(k, N) be the probability that the root has degree k in a network of N nodes. This probability was previously determined analytically for random attachment and linear preferential attachment networks [30], where it was shown that the root degree is broadly distributed for preferential attachment. For enhanced redirection, this probability obeys the difference equation

$$R(k, N+1) = \left\{ \frac{1}{N} + \frac{k-3}{N} \left[1 - \frac{1}{(k-1)^{\lambda}} \right] \right\} R(k-1, N) + \left\{ 1 - \frac{1}{N} - \frac{k-2}{N} \left[1 - \frac{1}{k^{\lambda}} \right] \right\} R(k, N)$$
(37)

with initial condition R(2,1) = 1. The first term gives the probability that the root has degree k-1 when the (N+1)th node connects to it, either directly or by redirection from the k-3 children of the root. The second term is the probability that the root has degree k and the (N+1)th node does *not* connect to the root. Iterating (37) numerically we obtain the numerical results for R(k,N) shown in figure 9.

The behavior of the moments $\langle k^d \rangle_N = \sum_k k^d R(k,N)$ helps to shed light on the behavior of the root degree distribution R(k,N). From equation (37), the mean degree evolves according to

$$N\langle k\rangle_{N+1} = (N+1)\langle k\rangle_N - \langle k^{1-\lambda}\rangle_N + 2\langle k^{-\lambda}\rangle_N - 1.$$
(38)

While this recurrence is not closed, we can drop the second term on the right-hand side of (38) as $N \to \infty$ because $\langle k^{1-\lambda} \rangle_N \ll \langle k \rangle_N$ (since $\lambda > 0$). The following terms are even

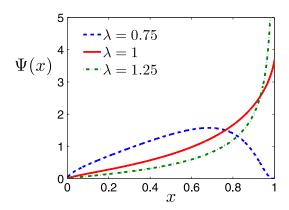


Figure 9. The scaling function from equation (39) for the root degree probability distribution.

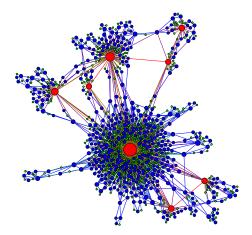


Figure 10. A network of $N=10^3$ nodes which has been built by enhanced redirection with double attachment, equation (42) with $\lambda=0.75$. The maximum degree in this example is $k_{\text{max}}=623$.

smaller. Hence (38) simplifies to $\langle k \rangle_{N+1} \simeq (1+1/N) \langle k \rangle_N$, from which $\langle k \rangle \sim N$. Similarly, the recurrence for the variance $\sigma_N = \langle k^2 \rangle_N - \langle k \rangle_N^2$ shows that $\langle \sigma \rangle_{N+1} \simeq (1+2/N) \langle \sigma \rangle_N$. This implies that $\langle \sigma \rangle \sim N^2$. Thus we conclude that the degree of the root is *not* self-averaging. The lack of self-averaging arises because the early evolution steps play a huge role in determining the root degree for large N.

Since the root degree is non self-averaging with $\langle k \rangle \sim N$, we anticipate that when $k, N \to \infty$, the probability distribution R(k, N) admits the scaling form (figure 9)

$$R(k,N) = N^{-1}\Psi(x), \qquad x = \frac{k}{N}.$$
 (39)

When $\lambda \leq 1$, the scaled distribution $\Psi(x)$ is a smooth function on [0,1], but when $\lambda > 1$, $\Psi(x)$ additionally contains a singular component $\mathcal{R}(\lambda)\delta(x-1)$ that accounts for the probability to create a star or a hairball about the root node. Therefore,

$$\mathcal{R}(\lambda) \ge \mathcal{S}(\lambda) = \prod_{n=1}^{\infty} \left[1 - \frac{n-1}{n(n+1)^{\lambda}} \right]. \tag{40}$$

Consider the extreme behaviors at $x \to 1$. This limit essentially corresponds to the probability of forming the star. Thus for $\lambda = 1$ we expect $\Psi(1) = \pi^{-1} \sinh \pi \approx 3.676$. This agrees with simulation results. Further, $\Psi(1) = 0$ for $\lambda < 1$; more precisely, the scaling function $\Psi(x)$ near x = 1 is essentially the same as the scaling function $\mathcal{M}_1(x)$ describing the largest macrohub, equation (17). Hence $\ln \Psi(x) \sim -(1-x)^{-(1-\lambda)}$ as $x \to 1$. This is also compatible with our numerical results.

To estimate the asymptotic behavior of $\Psi(x)$ for $x \to 0$, we have computed the probabilities that the root has the smallest possible degrees k = 3 and k = 4 (appendix B). From these results we infer the asymptotic behavior

$$\Psi(x) \sim x^{1-3^{-\lambda}} \qquad \text{as } x \to 0, \tag{41}$$

which is compatible with our numerical results.

5. Enhanced redirection with multiple attachments

In our discussion thus far, each new node added to the network has a single outgoing link. The resulting network is therefore a tree, except for closed loops that were part of the initial condition. It is therefore worthwhile to check whether the many anomalous features of enhanced redirection still exist if we allow the out-degree of each node to be larger than 1 so that closed loops can be created.

For simplicity, we consider the attachment rule in which a new node makes exactly two connections to existing nodes of the network—double attachment. We choose the initial condition of a single node with two self-loops, so that the root node is its own parents. Nodes are added sequentially according to the following rules.

- (i) The new node links to a randomly selected target node.
- (ii) The new node links to one of the two parents of the target with probabilities

$$r_1(a,b) = \frac{a^{\lambda}}{a^{\lambda} + b^{\lambda}}, \qquad r_2(a,b) = \frac{b^{\lambda}}{a^{\lambda} + b^{\lambda}},$$
 (42)

where where a and b are the degrees of parent 1 and parent 2.

We choose these redirection probabilities so that attachment to a given parent becomes increasingly likely as its degree increases: $r_1(a,b) \to 1$ as $a \to \infty$ and $r_2(a,b) \to 1$ as $b \to \infty$. Figure 10 shows a typical network that has been generated by rules (i) and (ii) with $\lambda = 0.75$.

5.1. Degree distribution

The two links created by each new node arise from two qualitatively different mechanisms. One link arises by random attachment, a mechanism that leads to the random recursive tree. The second link is to one of two parents, and it is selected via enhanced redirection which, as we have shown earlier, leads to a broad and non-extensive degree distribution, $N_k \sim N^{\nu-1}/k^{\nu}$ with $\nu < 2$. Because of the competition between these two mechanisms, one may anticipate that N_k scales linearly with N for sufficiently small k, while for large k the degree distribution N_k scales as $N^{\nu-1}$ with $\nu < 2$. This is indeed the case, as illustrated in figure 11.

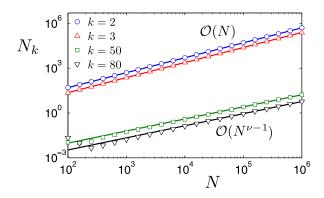


Figure 11. N_k versus N for degrees k=2, 3, and degrees k=50, 80 for $\lambda=0.75$. The lines indicate the scalings $N_k \sim N$ and $N_k \sim N^{\nu-1}$ with $\nu=1.8$.

To understand how two different scaling regimes emerge, we study the master equation that governs N_k (compare with equation (4)),

$$\frac{dN_k}{dN} = \left[\frac{N_{k-1}}{N} - \frac{N_k}{N}\right] + \left[\frac{\tau_{k-1}(k-3)N_{k-1}}{N} - \frac{\tau_k(k-2)N_k}{N}\right] + \delta_{k,2}.$$
 (43)

The terms in the first set of brackets account for attachment to a randomly selected target. Similarly, the terms in the second set of brackets account for redirection. Here,

$$\tau_k = \sum_b \frac{r_1(k,b)N(k,b)}{(k-2)N_k} \tag{44}$$

is the probability that an incoming node attaches to the degree-k parent of a random target, where N(a,b) is the number of nodes with parents of degree a and b. Thus τ_k is the probability of redirection to a degree-k parent averaged over all $(k-2)N_k$ children of this parent node. In (43), the expression $\tau_k(k-2)N_k/N$ gives the probability that the incoming node initially targets one of the $(k-2)N_k$ children of a node of degree k and then redirects to this parent. The term $\delta_{k,2}$ accounts for each newly-created node having degree 2.

To determine N_k , we separately analyze equation (43) for small and for large k. When k is small, we make the ansatz $N_k = c_k N$ and substitute into equation (43). Rearranging, we find the recursion relation for c_k for k > 2,

$$c_k = \frac{1 + (k - 3)\tau_{k-1}}{2 + (k - 2)\tau_k} c_{k-1},\tag{45}$$

while $c_2 = \frac{1}{2}$. This recursion has the product solution

$$c_k = \frac{1}{2^{k-1}} \prod_{j=3}^k \frac{1 + (j-3)\tau_{j-1}}{1 + (j-2)\tau_j/2}.$$
(46)

Since redirection to a low-degree parent is unlikely, we approximate the redirection probability as $\tau_k = 0$ for small k. With this approximation, the product in equation (46) equals 1, and the degree distribution reduces to $N_k \approx N/2^{(k-1)}$.

For large k, we substitute the non-extensive scaling ansatz $N_k = c_k N^{\nu-1}$ into equation (43), and rearrange to obtain the recursion

$$c_k = \frac{1 + (k - 3)\tau_{k-1}}{\nu + (k - 2)\tau_k} c_{k-1}.$$
(47)

This gives the product solution

$$c_k = c_\ell \prod_{j=\ell+1}^k \frac{1 + (j-3)\tau_{j-1}}{\nu + (j-2)\tau_j},\tag{48}$$

where ℓ is the degree above which the non-extensive scaling ansatz is valid. In the limit of large k, we approximate $\tau_k = 1$ in equation (48) because the probability of redirection to a high-degree parent node approaches 1. This gives the asymptotic behavior

$$c_k = c_\ell \prod_{j=\ell+1}^k \frac{j-2}{j-2+\nu} \sim k^{-\nu}.$$
 (49)

Combined with the non-extensive ansatz, the degree distribution for large k is $N_k \sim N^{\nu-1}/k^{\nu}$. In the above derivation, the precise value of ℓ only affects c_k up to a multiplicative factor but not the scaling behavior at large k.

We now define the crossover degree k^* as the value that separates the small-k extensive scaling regime from the large-k non-extensive regime. To estimate k^* , we find the value at which the N_k in the small- and large-k approximations coincide. This leads to the transcendental equation $N/2^{(k^*-1)} = N^{\nu-1}/(k^*)^{\nu}$, whose solution gives $k^* \sim \ln N$ to lowest order in N.

To summarize, the limiting degree distributions are

$$N_k = \begin{cases} N/2^{(k-1)} & k \ll k^*, \\ cN^{\nu-1}/k^{\nu} & k \gg k^*, \end{cases}$$
 (50)

with c a constant. As shown in figure 12, the agreement between this prediction and the degree distribution from Monte Carlo simulations is excellent. The double attachment rule also produces many macrohubs whose degrees grow linearly with N and also obey the same scaling behavior (12) as in single-attachment enhanced redirection.

5.2. Clustering coefficient

The new feature of double attachment is that the resulting network contains closed loops. A basic question about this type of network is whether it is homogeneous or highly clustered. We measure the level of clustering by the local clustering coefficient $C_i(k)$ [7] for a given node i of degree k. This quantity is defined as the ratio of the actual number of links between the neighbors of node i to the k(k-1)/2 possible links between these neighbors if they were all connected. For the complete graph, the local clustering coefficient equals 1 for every node, while in a tree network the local clustering coefficient is everywhere zero.

To compute the local clustering coefficient for double-attachment networks, consider an arbitrary node i with degree k. There are two ways that attachment can occur to this node: (i) a new node can attach directly to node i and to one of its parents, or

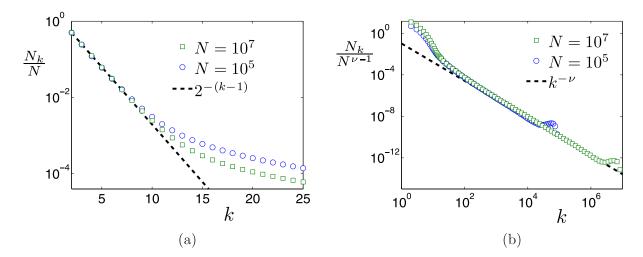


Figure 12. The degree distribution for double attachment with $\lambda = 0.75$ and $\nu = 1.8$ for $N = 10^5$, 10^6 , and 10^7 . In (a), the scaled degree distribution N_k/N is plotted to show the data collapse for small k. Similarly, (b) plots $N_k/N^{\nu-1}$ to show the data collapse at large k.

(ii) a new node can attach to one of the k-2 children of node i and to node i itself. In either case, the degree of node i increases by one and the number of links between the neighbors of node i also increases by one. When node i is first created it necessarily has degree 2 and a single link between its neighbors. Thus, when node i reaches degree k, it will have k-1 links between its neighbors. Therefore its clustering coefficient is $C_i(k) = (k-1)/[k(k-1)/2] = 2/k$. This 1/k scaling of the local clustering coefficient is seen in many real-world networks [39]–[42].

The average local clustering coefficient is therefore

$$\langle C_i \rangle = \frac{1}{N} \sum_{k=2}^{2N} C(k) N_k = \frac{2}{N} \sum_{k=2}^{2N} \frac{N_k}{k}.$$
 (51a)

We partition the sum according to whether k is smaller or greater than k^* using equation (50). This gives

$$\langle C_i \rangle = \sum_{k=2}^{k^*} \frac{1}{2^{(k-2)} k} + N^{\nu-2} \sum_{k=k^*+1}^{2N} \frac{2}{k^{\nu+1}}, \tag{51b}$$

with $k^* \sim \ln N$ the cutoff between the two scaling regimes. As $N \to \infty$, the second term vanishes, because $\nu < 2$, and the first term asymptotically approaches

$$\langle C_i \rangle \to \sum_{k=2}^{\infty} 2^{-(k-2)} k^{-1} = 4 \ln 2 - 2 = 0.77258...$$
 (51c)

This large value for this coefficient indicates a highly clustered network. For comparison, empirical studies found $\langle C_i \rangle = 0.79$ for actor collaboration networks [39], 0.68 for co-authorship networks [40], and 0.14 for blogging networks [41]. By contrast, the Erdős–Rényi random graphs have a vanishing mean local clustering coefficient; more precisely it decreases with N according to $\langle C_i \rangle \sim N^{-1}$ [42].

6. Conclusion

Enhanced redirection is an appealing mechanism that produces networks with a variety of unusual features, including the existence of multiple macroscopic hubs, anomalous scaling of the degree distribution, and lack of self-averaging. Unlike other models that produce macrohubs, enhanced redirection is based solely on local growth rules and does not assume intrinsic differences between nodes. Networks grown by enhanced redirection are highly disperse and typically consist of a set of loosely connected macrohubs that is reminiscent of airline route networks [7, 21], [24]–[26].

Intriguingly, the degree distribution decays more slowly than k^{-2} . Such an anomalously slow decay is mathematically consistent with a finite average degree only if the number of nodes of fixed degree scales sub-linearly with the number of nodes N. Enhanced redirection may thus provide the mechanism that underlies the wide range of networks [27] whose degree distributions apparently decay more slowly than k^{-2} .

We also combined the enhanced redirection mechanism with the simplest random attachment to produce networks that contain closed loops. The resulting degree distribution exhibits an unusual combination of extensive and non-extensive scaling. The clustering coefficient in these networks is large, as is observed in many real networks, and thus this rule produces highly clustered networks with numerous macrohubs.

Acknowledgments

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Appendix A. Star probability

To derive the asymptotic behaviors given in (14) we first rewrite (13) as

$$S_N(\lambda) = \prod_{n=1}^{N-1} \left[1 - \frac{n-1}{n(n+1)^{\lambda}} \right]. \tag{A.1}$$

When $\lambda > 1$, the product on the right-hand side of (A.1) converges to

$$S(\lambda) = \prod_{n=1}^{\infty} \left[1 - \frac{n-1}{n(n+1)^{\lambda}} \right]. \tag{A.2}$$

The probability $S(\lambda)$ is clearly positive and an increasing function of $\lambda > 1$ (see also figure 7). Numerical evaluation of the product gives, for example, $S(2) \approx 0.74562$, $S(4) \approx 0.984$, $S(6) \approx 0.999$.

When $\lambda \leq 1$, the product on the right-hand side of (A.1) converges to zero as $N \to \infty$. Consider first the case where $0 < \lambda < 1$. We take the logarithm of (A.1) and expand the logarithm. Since the dominant contribution arises for large n, we can replace the sum by an integral to yield

$$\ln S_N(\lambda) = \sum_{n=1}^{N-1} \ln \left[1 - \frac{n-1}{n(n+1)^{\lambda}} \right] \sim -\int^N \frac{\mathrm{d}n}{n^{\lambda}} = -\frac{N^{1-\lambda}}{1-\lambda}.$$

In the marginal case of $\lambda=1$ we use $\prod_{1\leq n\leq N-1}[1-(n+1)^{-1}]=N^{-1}$ to rewrite equation (A.1) at $\lambda=1$ as

$$NS_N(1) = \prod_{n=1}^{N-1} \frac{1 - \frac{n-1}{n(n+1)}}{1 - \frac{1}{n+1}} = \prod_{n=1}^{N-1} \left[1 + n^{-2} \right].$$
(A.3)

The last product converges to $A = \pi^{-1} \sinh \pi$ as $N \to \infty$ [43]. In the extreme case of $\lambda = 0$, we simplify equation (13) to give $S_N(0) = 1/(N-1)!$. These give the results summarized in equation (14). It is worth emphasizing that the third line in (14), $\exp[-N^{1-\lambda}/(1-\lambda)]$, represents only the controlling factor in the asymptotic behavior. Subdominant, and possibly less singular, contributions as $N \to \infty$ have been neglected.

The precise behavior of $S(\lambda)$ in the $\lambda - 1 \to 0^+$ limit can be extracted from (A.2). Taking the logarithm of the infinite product for $S(\lambda)$, expanding the logarithm, and separating the terms that converge and diverge as $\lambda \downarrow 1$, one gets

$$\ln \mathcal{S}(\lambda) = -\zeta(\lambda) + 2 + \sum_{n \ge 1} \left(\frac{n}{(n+1)(n+2)} + \ln \left[1 - \frac{n}{(n+1)(n+2)} \right] \right) + \mathcal{O}(\lambda - 1).$$

Recalling the asymptotic behavior of the zeta function, $\zeta(\lambda) = (\lambda - 1)^{-1} + \gamma + \mathcal{O}(\lambda - 1)$, where $\gamma \approx 0.5772$ is Euler's constant, we conclude that

$$S(\lambda) \simeq \exp\left(-\frac{1}{\lambda - 1} + C\right),$$
 (A.4)

with

$$C = 2 - \gamma + \sum_{n \ge 1} \left[\frac{n}{(n+1)(n+2)} + \ln\left(1 - \frac{n}{(n+1)(n+2)}\right) \right] \approx 1.3018.$$

Let us now compute the probability to form the star for the model with redirection probability $r(a,b) = a^{\lambda}/(a^{\lambda} + b^{\lambda})$. One gets

$$S_N(\lambda) = \prod_{n=1}^{N-1} \left[1 - \frac{n-1}{n} \frac{1}{1 + (n+1)^{\lambda}} \right],$$

from which

$$S_N(\lambda) \to \begin{cases} \widehat{S}(\lambda) & \lambda > 1\\ \widehat{A}/N & \lambda = 1 \end{cases}$$

$$\exp\left(-\frac{N^{1-\lambda}}{1-\lambda}\right) & 0 < \lambda < 1$$

$$\frac{N}{2^{N-1}} & \lambda = 0,$$
(A.5)

where $\widehat{\mathcal{S}}(\lambda) = \prod_{n\geq 1} \left[1 - ((n-1)/n)\left(1/(1+(n+1)^{\lambda})\right)\right]$ and the amplitude \widehat{A} is found by employing the same construction as that used for equation (A.3) to yield

$$\widehat{A} = \prod_{n=1}^{\infty} \left[1 + \frac{2n+1}{n^2(n+2)} \right] = \frac{2}{\pi} \cosh\left(\frac{\pi\sqrt{3}}{2}\right) = 4.856379592\dots$$

Again, the third line in equation (A.5) represents only the controlling factor in the asymptotic behavior in which subdominant contributions such as $N \to \infty$ have been neglected.

The predictions (14) and (A.5) for the two models (3) are qualitatively the same for the same values of λ . More precisely, qualitatively different networks emerge only in the extreme case $\lambda = 0$, which is not interesting to us as our goal is to study enhanced redirection. (When $\lambda = 0$, the first model in (3) leads to uniform attachment without redirection, the process that generates recursive random trees, while the second model in (3) leads to constant redirection probability r = 1/2, which is equivalent to a strictly linear preferential attachment process.)

Appendix B. Smallest root degree

The minimal possible degree of the root is k=3 when $N\geq 2$. Using (37) we get

$$R(3, N+1) = \left[1 - \frac{2 - 3^{-\lambda}}{N}\right] R(3, N).$$
(B.1)

Starting from R(3,2) = 1 and iterating (B.1) we find

$$R(3,N) = \frac{\Gamma(N-2+3^{-\lambda})}{\Gamma(N)\Gamma(3^{-\lambda})} \simeq \frac{1}{\Gamma(3^{-\lambda})} N^{-2+3^{-\lambda}}.$$
 (B.2)

Let us now probe the behavior of the second smallest degree k = 4. Using (37) we obtain the recurrence

$$R(4, N+1) = \frac{2-3^{-\lambda}}{N} R(3, N) + \left[1 - \frac{3-2^{1-2\lambda}}{N}\right] R(4, N).$$
 (B.3)

Making the substitution

$$R(4, N) = \frac{\Gamma(N - 3 + 2^{1 - 2\lambda})}{\Gamma(N)} U(N)$$
(B.4)

and using (B.2) we recast (B.3) into a simple recursion

$$U(N+1) = U(N) + \frac{2 - 3^{-\lambda}}{\Gamma(3^{-\lambda})} \frac{\Gamma(N-2 + 3^{-\lambda})}{\Gamma(N-2 + 2^{1-2\lambda})},$$

from which

$$U(N) = \frac{2 - 3^{-\lambda}}{\Gamma(2^{1 - 2\lambda})} + \frac{2 - 3^{-\lambda}}{\Gamma(3^{-\lambda})} \sum_{j=1}^{N-3} \frac{\Gamma(j + 3^{-\lambda})}{\Gamma(j + 2^{1 - 2\lambda})}.$$
 (B.5)

Using equations (B.4) and (B.5) one deduces the asymptotic behavior

$$R(4, N) \simeq \frac{1}{\Gamma(3^{-\lambda})} \frac{2 - 3^{-\lambda}}{1 + 3^{-\lambda} - 2^{1 - 2\lambda}} N^{-2 + 3^{-\lambda}}.$$
 (B.6)

The asymptotic behaviors (B.2) and (B.6) exhibit the same dependence on the number of nodes N. Generally, $R(k,N) \sim N^{-2+3^{-\lambda}}$ when $N \gg 1$ and k is kept finite. This, in conjunction with the scaling form (39), leads to the small-x behavior given in (41) for the scaled root degree distribution.

References

- [1] Yule G U, 1925 Phil. Trans. R. Soc. B **213** 21
- [2] Simon H A, 1955 Biometrika **42** 425
- [3] Barabási A-L and Albert R, 1999 Science 286 509
- [4] Krapivsky P L, Redner S and Leyvraz F, 2000 Phys. Rev. Lett. 85 4629
- [5] Dorogovtsev S N, Mendes J F F and Samukhin A N, 2000 Phys. Rev. Lett. 85 4633
- [6] Dorogovtsev S N and Mendes J F F, 2003 Evolution of Networks: From Biological Nets to the Internet and WWW (Oxford: Oxford University Press)
- [7] Newman M E J, 2010 Networks: An Introduction (Oxford: Oxford University Press)
- [8] Fabrikant A, Koutsoupias E and Papadimitriou C H, 2002 Automata, Languages and Programming (Lecture Notes in Computer Science vol 2380) (Berlin: Springer) pp 110–22
- [9] Colizza V, Banavar J R, Maritan A and Rinaldo A, 2004 Phys. Rev. Lett. 92 198701
- [10] Barthelemy M, 2011 Phys. Rep. 499 1
- [11] Papadopoulos F, Kitsak M, Ángeles Serrano M, Boguñá M and Krioukov D, 2012 Nature 489 537
- [12] Kleinberg J, Kumar R, Raghavan P, Rajagopalan S and Tomkins A, 1999 Proc. Int. Conf. Combinatorics and Computing (Lecture Notes in Computer Science) vol 1627 (Berlin: Springer) pp 1–18
- [13] Krapivsky P L and Redner S, 2001 Phys. Rev. E 63 066123
- [14] Vázquez A, 2003 Phys. Rev. E 67 056104
- [15] Rozenfeld H and ben-Avraham D, 2004 Phys. Rev. E 70 056107
- [16] Krapivsky P L and Redner S, 2005 Phys. Rev. E 71 036118
- [17] Lambiotte R and Ausloos M, 2007 Europhys. Lett. 77 58002
- [18] Ben-Naim E and Krapivsky P L, 2010 J. Stat. Mech. P06004
- [19] Gabel A and Redner S, 2013 J. Stat. Mech. P02043
- [20] Krapivsky P L and Krioukov D, 2008 Phys. Rev. E **78** 026114
- [21] i Cancho R F and Solé R V, 2003 Statistical Mechanics of Complex Networks (Lecture Notes in Physics no 625) (Berlin: Springer) p 114
- Bianconi G and Barábasi A-L, 2001 Phys. Rev. Lett. 86 5632
 Bianconi G and Barábasi A-L, 2001 Europhys. Lett. 54 436
- [23] Krapivsky P L and Redner S, 2002 Comput. Netw. 39 261
- [24] Bryan D L and O'Kelly M E, 1999 J. Reg. Sci. 39 275
- [25] Han J J et al. 2004 Nature **430** 88
- [26] Guimera R, Mossa S, Turtschi A and Amaral L A N, 2005 Proc. Nat. Acad. Sci. USA 102 7794
- [27] Kunegis J, Blattner M and Moser C, 2013 arXiv:1303.6271
- [28] Krapivsky P L and Redner S, 2002 J. Phys. A: Math. Gen. **35** 9517
- [29] Gabel A, Krapivsky P L and Redner S, 2013 Phys. Rev. E 88 050802
- [30] Krapivsky P L and Redner S, 2002 Phys. Rev. Lett. 89 258703
- [31] Gumbel E J, 1958 Statistics of Extremes (New York: Columbia University Press)
- [32] Eggenberger F and Pólya G, 1923 Z. Angew. Math. Mech. 3 279
- [33] Johnson N and Kotz S, 1977 Urn Models and Their Applications: An Approach to Modern Discrete Probability Theory (New York: Wiley)
 - Mahmoud H M, 2008 Pólya Urn Models (London: Chapman and Hall)
- [34] Derrida B and Flyvbjerg H, 1987 J. Phys. A: Math. Gen. 20 5273 Derrida B and Flyvbjerg H, 1987 J. Physique 48 971
 - Derrida B and Bessis D, 1988 J. Phys. A: Math. Gen. 21 L509
- [35] Higgs P G, 1995 Phys. Rev. E **51** 95
- [36] Frachebourg L, Ispolatov I and Krapivsky P L, 1995 Phys. Rev. E 52 R5727
- [37] Derrida B and Jung-Muller B, 1999 J. Stat. Phys. **94** 277
- [38] Krapivsky P L, Grosse I and Ben-Naim E, 2000 Phys. Rev. E 61 R993 Krapivsky P L, Ben-Naim E and Grosse I, 2004 J. Phys. A: Math. Gen. 37 2863
- [39] Watts D and Strogatz S H, 1998 Nature **393** 440
- [40] Acedo F J, Barroso C, Casanueva C and Galán J L, 2006 J. Manage. Stud. 43 970
- [41] Fu F, Liu L and Wang L, 2008 Physica A **387** 675
- [42] Ravasz E and Barabási A-L, 2003 Phys. Rev. E 67 026112
- [43] Abramowitz M and Stegun A, 1965 Handbook of Mathematical Functions (New York: Dover)