

The Generalized Lotka Volterra Equation

R.S. breaking + Cavity

Last lecture we introduced the G.L.V. equations

ρ is non-reciprocal

$$\frac{dN_c}{dt} = g_c N_c (r_c - N_c - \sum A_{ij} N_j) \quad l=1 \dots S$$

$$r_c \sim N(\bar{r}, \sigma_r^2)$$

$$A_{ij} = \frac{\mu}{S} + \sigma a_{ij}$$

$$\langle a_{ij} \rangle = 0 \quad \langle a_{ij} a_{kl} \rangle = \frac{1}{S} (\delta_{ij} \delta_{kl} + \rho \delta_{ik} \delta_{jl})$$

The key to these equations is that species could go extinct. i.e. $N_c = 0$ was fixed point.

Hence, D.O.F's can "appear" and "disappear".

We have S species initially in region species pool and some $S^* \leq S$ can survive...

Analyzing this is trying to understand the relationship between "diversity" (how big S^*) is and stability (is the ecosystem stable to small perturbations...)

There are conceptually two different kinds of perturbations...

(1) Say the system reaches a fixed point \vec{N}^* . We can $\vec{N}^* \rightarrow \vec{N}^* + \delta \vec{N}$ and ask if this grows or shrinks (usual stability for dynamical system / ODE)

(2) A second kind of stability is perturbation to changes in parameters. We will be concerned primarily with

$$\vec{r} \rightarrow \vec{r} + \delta \vec{r}$$

These are related and we will come back to this..

Since May's argument, lots of questions about how mathematically we see large diverse systems since for "generic" linearized dynamics"

$$\frac{d\vec{N}}{dt} \approx (I - L) \delta \vec{N} \quad L=1, \dots, S^*$$

Linearized matrix assumed random

we expect there to be transition for S^* sufficiently big (positive eigenvalue).. $L_{ij} \sim N(0, \sigma_L^2)$

To answer this question, we can ask about steady states of C.L.V. $r_i \sim N(\bar{r}, \sigma_r^2)$

$$0 = N_i (r_i - N_i - \sum_j A_{ij} N_j)$$

We can define an effective carrying capacity N_i^{eff}

$$r_c^{eff} = r_c - \sum_j A_{cj} N_j = r_c - \mu \langle N \rangle - \sum_j a_{cj} N_j$$

Notice that this clearly a "Mean-field" since it is sum over many variables.

Also notice that this is "almost linear" problem we solved...

We will need one more criteria... We require the steady-state to be non-invasible... Namely if there exists $N_c^* > 0$, we take that steady state... $0 = N_c (r_c^{eff} - N_c)$

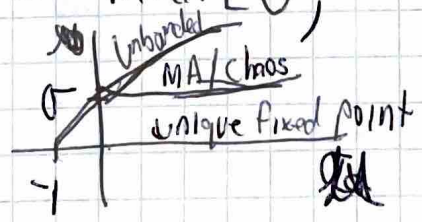
So that

$$N_c^* = \max [0, r_c - N_c - \sum A_{cj} N_j]$$

$$N_c^* = \max [0, r_c^{eff} \langle N \rangle]$$

Solution of linear equations

Phase Diagram



just $r_c^{eff} < 0$ means species goes extinct..

So now let us as before assume R.S. this is just the statement that

r_c^{eff} should be modeled as Gaussian with mean and variance...

As before to calculate this mean and variance we imagine "invading" the ecosystem with new species $N_0 \dots$

$$0 = N_c (r_c - N_c - \sum_{j \neq c} A_{cj} N_j - A_{c0} N_0) = N_c (r_c - N_c - \mu \langle N \rangle - \sum_{j \neq c} A_{cj} N_j - a_{c0} N_0)$$

$$0 = N_0 (r_0 - N_0 - \sum_{j=1}^N A_{0j} N_j) = N_0 (r_0 - N_0 - \mu \langle N \rangle - a_{00} N_0 - \sum_{j=1}^N A_{0j} N_j)$$

We proceed almost identical to as before...

Notice $r_c \rightarrow r_c + \delta r_c$ with $\delta r_c = -A_{c0} N_0$

So we can define susceptibility matrix..

$$\chi_{ij} = \frac{\partial N_i}{\partial r_j}$$

and write

$$N_c = N_{c10} - \sum_{j=1}^S \chi_{cj} a_{j0} N_0$$

Substituting we get to leading order in $1/S$

$$0 = N_0 [r_0 - N_0 - \mu \langle N \rangle - \sigma \sum_j a_{0j} (N_{j10} - \sum_k \chi_{jk} a_{k0} N_0)]$$

There are two solutions for N_0

$$N_0 = 0$$

$$N_0 = \frac{\bar{r} + \delta r_0 - \sigma \sum_j a_{0j} N_{j10}}{1 - \sigma^2 \sum_j \chi_{jj} a_{0j} a_{j0}}$$

Since we are taking large S limit we can substitute

$$\sum_{j=1}^S a_{0j} a_{j\omega} v_j = \sum_{j=1}^S \langle a_{0j} a_{j\omega} \rangle v_j$$

$$= \frac{\rho}{S} \sum_{j=1}^S v_j = \rho \tilde{v}$$

$$\tilde{v} = \frac{1}{S} \sum_j v_j \dots$$

$$N_0 = \max \left(0, \frac{r + \delta r_0 - \sigma_R^2 \sum_{j=1}^S a_{0j} N_{j10} - \mu \langle N \rangle}{1 - \sigma^2 \rho \tilde{v}} \right)$$

Gaussian Field R.S.
 r_0^{eff}

$$\langle r_0^{\text{eff}} \rangle = r - \mu \langle N \rangle$$

$$\langle N \rangle = \frac{1}{S} \sum_j N_j$$

$$\langle \delta r_0^{\text{eff}} \rangle = \sigma_R^2 + \sigma_A^2 \langle N \rangle^2$$

$$\langle N^2 \rangle = \frac{1}{S} \sum_j N_j^2$$

$$N_0 = \max \left(0, \frac{r - \mu \langle N \rangle + \sqrt{\sigma_R^2 + \sigma_A^2 \langle N \rangle^2} Z_N}{1 - \rho \sigma^2 \tilde{v}} \right)$$

Unit normal variable

As before we will invoke self-averaging and see species distribution is self-consistent truncated Gaussian

Let us call fraction of surviving species $\frac{S^*}{S} = \phi_N$
then we can

$$\phi_{N_0} = \phi_N = \frac{1}{S} \sum_{i=1}^S \theta(N_{i0})$$

$$\langle N_0 \rangle = \langle N \rangle = \frac{1}{S} \sum_{i=1}^S N_{i0}$$

$$\langle N_0^2 \rangle = \langle N^2 \rangle = \frac{1}{S} \sum_{i=1}^S N_{i0}^2$$

$$\left\langle \frac{dN_0}{dt_0} \right\rangle = \frac{1}{S} \sum_{i=1}^S v_{i0}$$

So notice

$$\bar{v} = \left\langle \frac{\partial N_0}{\partial t_0} \right\rangle = \frac{\phi_N}{1 - p_0^2 v}$$

Non-zero fraction since

$$\frac{\partial N_0}{\partial t_0} = 0 \text{ if } N_0 = 0$$

Useful to define $y = ma(a, \frac{cx+q}{b})$

$$\frac{\partial N_0}{\partial t_0} = \frac{1}{1 - p_0^2 v} \text{ if } N_0 \neq 0$$

$$\langle y^j \rangle = \frac{1}{\sqrt{2\pi}} \int_{-c/q}^{\infty} e^{-\frac{x^2}{2}} \left(\frac{c}{b}x + \frac{q}{b} \right)^j dx$$

$$= \left(\frac{c}{b} \right)^j \frac{1}{\sqrt{2\pi}} \int_{-q/c}^{\infty} e^{-\frac{x^2}{2}} \left(x + \frac{q}{c} \right)^j dx$$

$$= \left(\frac{c}{b}\right)^d W_1\left(\frac{a}{c}\right)$$

⑦

where

$$W_1(\Delta) = \frac{1}{2\pi} \int_{-\Delta}^{\infty} e^{-z^2} (z + \Delta) dz$$

(Truncated Gaussian)
distribution integral

Need identity

In terms of $W_1(\Delta)$ get $W_2(\Delta) = W_0(\Delta) + \Delta$

$$v = \frac{\phi_N}{1 - p\sigma^2}$$

$$\phi_N = \frac{W_0\left(\frac{r - \mu\langle N \rangle}{\sqrt{\sigma_r^2 + \sigma^2\langle N^2 \rangle}}\right)}{\sqrt{\sigma_r^2 + \sigma^2\langle N^2 \rangle}}$$

$$\langle N \rangle = \frac{\sqrt{\sigma_r^2 + \sigma^2\langle N^2 \rangle}}{1 - p\sigma^2 v} W_1\left(\frac{r - \mu\langle N \rangle}{\sqrt{\sigma_r^2 + \sigma^2\langle N^2 \rangle}}\right)$$

$$\langle N^2 \rangle = \frac{\sigma_r^2 + \sigma^2\langle N^2 \rangle}{(1 - p\sigma^2 v)^2} W_2\left(\frac{r - \mu\langle N \rangle}{\sqrt{\sigma_r^2 + \sigma^2\langle N^2 \rangle}}\right)$$

This allows us to calculate M.F. R.S. distribution exactly,

Truncated Gaussian with means, moments, and non-zero fraction given by above self-consistent equations..

So far we have not said anything about the stability of this Replica Symmetric ansatz..

We would like the M.F. solution to exist and be stable to perturbations..

Happens numerically

If it is not stable, then we say that ^{as σ_c increased} there is Replica-Symmetry breaking...

\Rightarrow check when infinitely susceptible to small perturb

To do this we imagine taking the surviving species and perturbing them

$$\vec{N}^* \rightarrow N^* + \epsilon \vec{\eta}$$

↙ strength
↘ random vector

$\langle \vec{\eta} \cdot \vec{\eta} \rangle \sim \text{finite } \bar{\eta}^2$

We can calculate the effect this will have by looking at our cavity species

$$N_0 = \max \left[0, \frac{r_0 - \mu \langle N \rangle - \sigma \sum a_{0j} N_{j0} + \epsilon \eta_0}{1 - \rho \sigma^2 \nu} \right]$$

We would like to measure sensitivity to perturbation

$$\left\langle \left(\frac{dN_0}{d\epsilon} \right)^2 \right\rangle$$

(square because want to know when variance diverges)

Note that

$$\frac{dN_0}{d\varepsilon} = 0 \quad \text{if } N_0 = 0$$

$$\frac{dN_0}{d\varepsilon} = - \frac{\sigma \sum_{j=1}^{S^*} a_{0j} \frac{\partial N_{j0}}{\partial \varepsilon} + \eta_0 a_{00}}{1 - \rho \sigma^2 v}$$

$$\left\langle \left(\frac{dN_0}{d\varepsilon} \right)_+^2 \right\rangle =$$

(Over non zero) value

$$\frac{\sigma^2 \frac{S^*}{S} \sum_{j=1}^{S^*} \left(\frac{\partial N_{j0}}{\partial \varepsilon} \right)^2 + \eta_0^2 \sigma^2}{(1 - \rho \sigma^2 v)^2}$$

$\bar{\eta}^2 = \sum_{i=1}^S \eta_i^2$
 which we choose to scale to be finite
 Variance of random variable

We now write

$$\frac{1}{S^*} \sum_{j=1}^S \left(\frac{\partial N_{j0}}{\partial \varepsilon} \right)^2 = \left\langle \left(\frac{\partial N_{j0}}{\partial \varepsilon} \right)_+^2 \right\rangle$$

non-zero average

But self averaging say

$$\left\langle \left(\frac{\partial N_{j0}}{\partial \varepsilon} \right)_+^2 \right\rangle = \left\langle \left(\frac{\partial N_0}{\partial \varepsilon} \right)_+^2 \right\rangle$$

So that

$$\left[(1 - \rho \sigma^2 v)^2 - \sigma^2 \phi_s \right] \left\langle \left(\frac{\partial N_0}{\partial \varepsilon} \right)_+^2 \right\rangle = \bar{\eta}_0^2 \sigma^2 \frac{1}{(1 - \rho \sigma^2 v)}$$

$$= 0 \quad \left\langle \left(\frac{\partial N_0}{\partial \varepsilon} \right)_+^2 \right\rangle = \frac{\bar{\eta}_0^2 \sigma^2}{(1 - \rho \sigma^2 v) \left[(1 - \rho \sigma^2 v)^2 - \sigma^2 \phi_s \right]}$$

This susceptibility diverges when

$$(1 - \rho \sigma_c^2 v_c)^2 = \sigma_c^2 \phi_s$$

(c - critical subscript)
 v_c, σ_c

Use

$$\phi_s = v_c (1 - \rho \sigma_c^2 v_c)$$

To get

$$(1 - \rho \sigma_c^2 v_c)(1 - \rho \sigma_c^2 v_c - \sigma_c^2 v_c) = 0$$

Non-zero since
by assumption
R.S. solution
exists

$$v_c = \frac{1}{\sigma_c^2 (1 + \rho)}$$

zero

We still need another equation to solve for v_c
to do this we note

$$(1 - \rho \sigma_c^2 v_c) = \left(1 - \frac{\rho \sigma_c^2}{\rho \sigma_c^2 + \sigma_c^2}\right) = \frac{\sigma_c^2}{\rho \sigma_c^2 + \sigma_c^2} = \frac{1}{1 + \rho}$$

~~From this we find that~~
 ~~$\phi_s = v_c (1 - \rho \sigma_c^2 v_c)$~~
 ~~$\phi_s = \frac{1}{1 + \rho}$~~
~~(HW)~~
~~to show that~~
 ~~$\frac{1}{\sigma_c^2} = \frac{1 + \rho}{\sqrt{2}}$~~
 ~~$\frac{1}{\sigma_c^2} = \frac{1 + \rho}{\sqrt{2}}$~~

and use (HW)
identity $w_2(\Delta) = w_0(\Delta) + \Delta$
to find
$$\frac{1}{\sigma_c^2} = \frac{1 + \rho}{\sqrt{2}}$$