

Problem Set 5: Calculus on Manifolds I

1 The Gradient (5 points)

Consider the function on $f : \mathbb{R}^3 \rightarrow \mathbb{R}$ defined in Cartesian coordinates by $f(x, y, z) = x^2 + y^2 + z^2$.

- (a) Compute the gradient $df = a_i dx^i$ in Cartesian coordinates, i.e. find expressions for a_i .

SOLUTION:

The gradient is

$$df = \frac{\partial f}{\partial x} dx + \frac{\partial f}{\partial y} dy + \frac{\partial f}{\partial z} dz = 2x dx + 2y dy + 2z dz. \quad (1)$$

- (b) Evaluate $df(V)$ for $V = x\partial_x - y\partial_y + \partial_z$.

SOLUTION:

Given that $dx^\mu(\partial_\nu) = \delta_\nu^\mu$, we have

$$df(V) = 2x^2 - 2y^2 + 2z. \quad (2)$$

- (c) Express df in spherical coordinates (r, θ, ϕ) .

SOLUTION:

We can see immediately that $f(x, y, z) = r^2$, and so without any further computation, we have $df = 2r dr$.

- (d) Evaluate $df(V)$ for $V = \partial_r - \partial_\theta + \partial_\phi$.

SOLUTION: As before, $dx^\mu(\partial_\nu) = \delta_\nu^\mu$, and therefore $df(V) = 2r$.

2 Properties of the Lie Bracket (10 points)

Prove the following properties of the Lie bracket between two vector fields X and Y :

- (a) The Lie bracket is linear, i.e. for functions f, g on M and a, b real numbers, $[X, Y](af + bg) = a[X, Y]f + b[X, Y]g$.

SOLUTION: Starting from the definition of the Lie bracket, we have

$$[X, Y](f) = X(Y(f)) - Y(X(f)). \quad (3)$$

From the linearity of X and Y (since they are simply directional derivatives),

$$[X, Y](af + bg) = X(Y(af + bg)) - Y(X(af + bg)) = X(aY(f) + bY(g)) - Y(aX(f) + bX(g)). \quad (4)$$

Now applying linearity again,

$$\begin{aligned} [X, Y](af + bg) &= aX(Y(f)) + bX(Y(g)) - aY(X(f)) - bY(X(g)) \\ &= a[X, Y](f) + b[X, Y](g), \end{aligned} \quad (5)$$

as required.

(b) The Lie bracket obeys Leibniz's rule, i.e. $[X, Y](fg) = f[X, Y]g + g[X, Y]f$.

SOLUTION: Again, remembering that X and Y are simply directional derivatives, and therefore individually satisfy Leibniz's rule, we have

$$[X, Y](fg) = X(Y(fg)) - Y(X(fg)) = X(fY(g) + gY(f)) - Y(fX(g) + gX(f)). \quad (6)$$

The object $X(g)$ is simply taking the directional derivative at every point on the manifold of the function g , and so is yet another scalar function. We can therefore apply Leibniz's rule for each of X and Y to get

$$\begin{aligned} [X, Y](fg) &= X(f)Y(g) + fX(Y(g)) + X(g)Y(f) + gX(Y(f)) \\ &\quad - Y(f)X(g) - fY(X(g)) - Y(g)X(f) - gY(X(f)) \\ &= f[X(Y(g)) - Y(X(g))] + g[X(Y(f)) - Y(X(f))] \\ &= f[X, Y]g + g[X, Y]f, \end{aligned} \quad (7)$$

as required.

(c) In terms of the components of $X \equiv X^\mu \partial_\mu$ and $Y \equiv Y^\nu \partial_\nu$, the Lie bracket is given by $[X, Y]^\mu = X^\nu \partial_\nu Y^\mu - Y^\nu \partial_\nu X^\mu$.

SOLUTION: In terms of components, we see that

$$[X, Y] = X^\mu \partial_\mu Y^\nu \partial_\nu - Y^\mu \partial_\mu X^\nu \partial_\nu = (X^\mu \partial_\mu Y^\nu - Y^\mu \partial_\mu X^\nu) \partial_\nu. \quad (8)$$

and therefore

$$[X, Y]^\mu = X^\nu \partial_\nu Y^\mu - Y^\nu \partial_\nu X^\mu, \quad (9)$$

as required.

(d) $[X, Y]^\mu$ transforms under a change of coordinates as a vector.

SOLUTION: Under a change of coordinates, we see that

$$\begin{aligned} X^\nu \partial_\nu Y^\mu &\mapsto X'^\alpha \partial'_\alpha Y'^\beta \\ &= \frac{\partial x'^\alpha}{\partial x^\nu} X^\nu \frac{\partial x^\nu}{\partial x'^\alpha} \partial'_\alpha \left(\frac{\partial x'^\beta}{\partial x^\mu} Y^\mu \right) \\ &= X^\nu \left[\frac{\partial x'^\beta}{\partial x^\mu} \partial_\nu Y^\mu + Y^\mu \frac{\partial x'^\beta}{\partial x^\mu \partial x^\nu} \right]. \end{aligned} \quad (10)$$

Similarly, we have

$$Y^\nu \partial_\nu X^\mu \mapsto Y^\nu \left[\frac{\partial x'^\beta}{\partial x^\mu} \partial_\nu X^\mu + X^\mu \frac{\partial x'^\beta}{\partial x^\mu \partial x^\nu} \right], \quad (11)$$

which we can obtain easily by just swapping $X \leftrightarrow Y$. Therefore,

$$\begin{aligned} X^\nu \partial_\nu Y^\mu - Y^\nu \partial_\nu X^\mu &\mapsto X^\nu \left[\frac{\partial x'^\beta}{\partial x^\mu} \partial_\nu Y^\mu + Y^\mu \frac{\partial x'^\beta}{\partial x^\mu \partial x^\nu} \right] - Y^\nu \left[\frac{\partial x'^\beta}{\partial x^\mu} \partial_\nu X^\mu + X^\mu \frac{\partial x'^\beta}{\partial x^\mu \partial x^\nu} \right] \\ &= X^\nu \left[\frac{\partial x'^\beta}{\partial x^\mu} \partial_\nu Y^\mu + Y^\mu \frac{\partial x'^\beta}{\partial x^\mu \partial x^\nu} \right] - Y^\nu \left[\frac{\partial x'^\beta}{\partial x^\mu} \partial_\mu X^\nu + X^\nu \frac{\partial x'^\beta}{\partial x^\nu \partial x^\mu} \right] \\ &= \frac{\partial x'^\beta}{\partial x^\mu} (X^\nu \partial_\nu Y^\mu - Y^\mu \partial_\mu X^\nu), \end{aligned} \quad (12)$$

which therefore transforms as a vector. The pesky part which also was the reason why the partial derivative of a vector isn't a tensor drops out. Note that in the second to last line, I am simply relabeling $\mu \leftrightarrow \nu$, which I can do because the indices are summed over and are dummy indices. I also used the fact that partial derivatives commute (at least in physics).

3 The Lie Derivative of the Metric (10 points)

Prove that the Lie derivative of the metric is

$$(\mathcal{L}_X g)_{\mu\nu} = X^\alpha \partial_\alpha g_{\mu\nu} + g_{\mu\alpha} \partial_\nu X^\alpha + g_{\alpha\nu} \partial_\mu X^\alpha. \quad (13)$$

SOLUTION: The Lie derivative is defined as follows: for vector fields X and Y , as well as a scalar field f , we have $\mathcal{L}_X f = X(f)$, and $\mathcal{L}_X Y = [X, Y]$. Furthermore, the Lie derivative obeys linearity and Leibniz's rule.

With this in mind, let W and Y be vector fields on the manifold. Consider the object $g_{\mu\nu} W^\mu Y^\nu$, which is a scalar. The Lie derivative acting on a scalar should give

$$\begin{aligned} \mathcal{L}_X (g_{\mu\nu} W^\mu Y^\nu) &= X^\alpha \partial_\alpha (g_{\mu\nu} W^\mu Y^\nu) \\ &= W^\mu Y^\nu X^\alpha \partial_\alpha g_{\mu\nu} + X^\alpha g_{\mu\nu} \partial_\alpha (W^\mu Y^\nu). \end{aligned} \quad (14)$$

On the other hand, applying Leibniz's rule,

$$\begin{aligned} \mathcal{L}_X (g_{\mu\nu} W^\mu Y^\nu) &= \mathcal{L}_X g_{\mu\nu} \cdot W^\mu Y^\nu + g_{\mu\nu} \mathcal{L}_X (W^\mu Y^\nu) \\ &= \mathcal{L}_X g_{\mu\nu} \cdot W^\mu Y^\nu + g_{\mu\nu} (Y^\nu \mathcal{L}_X W^\mu + W^\mu \mathcal{L}_X Y^\nu). \end{aligned} \quad (15)$$

But

$$\begin{aligned} Y^\nu \mathcal{L}_X W^\mu + W^\mu \mathcal{L}_X Y^\nu &= Y^\nu (X^\alpha \partial_\alpha W^\mu - W^\alpha \partial_\alpha X^\mu) + W^\mu (X^\alpha \partial_\alpha Y^\nu - Y^\alpha \partial_\alpha X^\nu) \\ &= X^\alpha \partial_\alpha (W^\mu Y^\nu) - (W^\alpha \partial_\alpha X^\mu) Y^\nu - (Y^\alpha \partial_\alpha X^\nu) W^\mu \end{aligned} \quad (16)$$

Therefore, comparing Eqs. (14) and (15), we have

$$\begin{aligned} \mathcal{L}_X g_{\mu\nu} \cdot W^\mu Y^\nu &= W^\mu Y^\nu X^\alpha \partial_\alpha g_{\mu\nu} + g_{\mu\nu} [(W^\alpha \partial_\alpha X^\mu) Y^\nu + (Y^\alpha \partial_\alpha X^\nu) W^\mu] \\ &= W^\mu Y^\nu X^\alpha \partial_\alpha g_{\mu\nu} + W^\mu Y^\nu g_{\alpha\nu} \partial_\mu X^\alpha + W^\mu Y^\nu g_{\mu\alpha} \partial_\nu X^\alpha, \end{aligned} \quad (17)$$

where we have done some relabeling of the dummy indices in the last line. We can therefore conclude that

$$\mathcal{L}_X g_{\mu\nu} = X^\alpha \partial_\alpha g_{\mu\nu} + g_{\mu\alpha} \partial_\nu X^\alpha + g_{\alpha\nu} \partial_\mu X^\alpha, \quad (18)$$

as required.

4 (SG 11.2) Rigid Rotations on the Sphere (10 points)

The metric on the unit sphere in polar coordinates is

$$g = d\theta \otimes d\theta + \sin^2 \theta d\phi \otimes d\phi. \tag{19}$$

Consider

$$V_x = -\sin \phi \partial_\theta - \cot \theta \cos \phi \partial_\phi, \tag{20}$$

which is a vector field of a rigid rotation about the x -axis. Show that $\mathcal{L}_{V_x} g = 0$, i.e. that rigid rotations do not deform the metric.

SOLUTION: For convenience, let me just rename V_x as V . Then the Lie derivative of the metric is

$$(\mathcal{L}_V g)_{\mu\nu} = V^\alpha \partial_\alpha g_{\mu\nu} + g_{\mu\alpha} \partial_\nu V^\alpha + g_{\alpha\nu} \partial_\mu V^\alpha. \tag{21}$$

Component by component, noting that the derivative of $g_{\theta\theta}$ in any coordinate is zero, since it is a constant, we have

$$\begin{aligned} (\mathcal{L}_V g)_{\theta\theta} &= 2g_{\theta\theta} \partial_\theta V^\theta \\ &= 0, \end{aligned} \tag{22}$$

since V^θ only depends on ϕ . Next,

$$\begin{aligned} (\mathcal{L}_V g)_{\theta\phi} &= g_{\theta\theta} \partial_\phi V^\theta + g_{\phi\phi} \partial_\theta V^\phi \\ &= -\cos \phi + \sin^2 \theta \cdot (\csc^2 \theta \cos \phi) \\ &= 0. \end{aligned} \tag{23}$$

$(\mathcal{L}_V g)_{\phi\theta} = (\mathcal{L}_V g)_{\theta\phi} = 0$ by symmetry (you can check by permuting the indices yourself). Finally,

$$\begin{aligned} (\mathcal{L}_V g)_{\phi\phi} &= V^\theta \partial_\theta g_{\phi\phi} + 2g_{\phi\phi} \partial_\phi V^\phi \\ &= -\sin \phi (2 \sin \theta \cos \theta) + 2 \sin^2 \theta (\cot \theta \sin \phi) \\ &= 0. \end{aligned} \tag{24}$$

Therefore $\mathcal{L}_{V_x} g = 0$, as required.

5 Wedge Product (10 points)

Consider two one-forms λ, η on \mathbb{R}^2 , defined in Cartesian coordinates by

$$\lambda = 2xy dx + (x^2 + y^2)dy, \quad \eta = e^{xy} dx - dy. \tag{25}$$

- (a) Compute the wedge product $\lambda \wedge \eta$.

SOLUTION: The wedge product is

$$\begin{aligned}\lambda \wedge \eta &= (2xy \, dx + (x^2 + y^2) \, dy) \wedge (e^{xy} \, dx - dy) \\ &= -2xy \, dx \wedge dy + (x^2 + y^2) e^{xy} \, dy \wedge dx \\ &= -(2xy + (x^2 + y^2) e^{xy}) \, dx \wedge dy.\end{aligned}\tag{26}$$

(b) Compute $\lambda \wedge \eta(V, W)$ where $V = x\partial_x + y\partial_y$ and $W = \partial_x - \partial_y$.

SOLUTION: We have

$$\begin{aligned}\lambda \wedge \eta(V, W) &= -(2xy + (x^2 + y^2) e^{xy}) \, dx \wedge dy(V, W) \\ &= -(2xy + (x^2 + y^2) e^{xy})(-x - y) \\ &= (x + y)(2xy + (x^2 + y^2) e^{xy}).\end{aligned}\tag{27}$$

(c) Express $\lambda \wedge \eta$ in polar coordinates (ρ, θ) .

SOLUTION: Switching to polar coordinates, we have $x = \rho \cos \theta$, $y = \rho \sin \theta$, and

$$\begin{aligned}xy &= \rho^2 \cos \theta \sin \theta, \\ (x^2 + y^2) e^{xy} &= \rho^2 e^{\rho^2 \cos \theta \sin \theta}.\end{aligned}\tag{28}$$

Furthermore,

$$\begin{aligned}dx \wedge dy &= \left(\frac{\partial x}{\partial \rho} d\rho + \frac{\partial x}{\partial \theta} d\theta \right) \wedge \left(\frac{\partial y}{\partial \rho} d\rho + \frac{\partial y}{\partial \theta} d\theta \right) \\ &= \rho \cos^2 \theta \, d\rho \wedge d\theta + (-\rho \sin \theta)(\sin \theta) d\theta \wedge d\rho \\ &= \rho(\sin^2 \theta + \cos^2 \theta) d\rho \wedge d\theta \\ &= \rho \, d\rho \wedge d\theta.\end{aligned}\tag{29}$$

Thus,

$$\lambda \wedge \eta = -(2\rho^2 \cos \theta \sin \theta + \rho^2 e^{\rho^2 \cos \theta \sin \theta}) \rho \, d\rho \wedge d\theta.\tag{30}$$

6 Exterior Derivatives (10 points)

Compute the exterior derivatives of the following forms:

(a) $f = xy$.

SOLUTION:

$$df = \frac{\partial f}{\partial x} dx + \frac{\partial f}{\partial y} dy = y \, dx + x \, dy.\tag{31}$$

(b) $A = y \, dx + xy \, dy$.

SOLUTION:

$$\begin{aligned}dA &= dy \wedge dx + (y dx + x dy) \wedge dy \\ &= (y - 1)dx \wedge dy.\end{aligned}\tag{32}$$

(c) $F = yx dx \wedge dy + xz dy \wedge dz.$ **SOLUTION:**

$$\begin{aligned}dF &= (y dx + x dy) \wedge dx \wedge dy + (x dz + z dx) \wedge dy \wedge dz \\ &= z dx \wedge dy \wedge dz.\end{aligned}\tag{33}$$

(d) $\omega = xyz(dx \wedge dy + dy \wedge dz + dz \wedge dx).$ **SOLUTION:**

$$\begin{aligned}d\omega &= (yz dx + xz dy + xy dz) \wedge (dx \wedge dy + dy \wedge dz + dz \wedge dx) \\ &= (xy + yz + xz)dx \wedge dy \wedge dz.\end{aligned}\tag{34}$$