



Time	Saturday, 26.08	Sunday, 26.08	Monday, 27.08	Tuesday, 28.08	Wednesday, 29.08	Thursday, 30.08	Thu
8:30	Registration						8:30
9:00	Breakfast						8:30
9:10	(Plenary Session) K. Lehtinen	(Plenary Session) K. Lehtinen	(Plenary Session) B. Barlow	(Invited Speech) C. Johnston		Plenary 8:30	8:30
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## XXIII INTERNATIONAL SYMPOSIUM ON THE JAHN-TELLER EFFECT

VIBRONIC COUPLING AND ELECTRON-PHONON INTERACTION  
IN MOLECULES AND CRYSTALS

BASIC JAHN-TELLER THEORY  
JAHN-TELLER AND PSEUDO JAHN-TELLER  
COUPLING IN MOLECULES AND CRYSTALS  
CONICAL INTERSECTIONS  
SUPERCONDUCTIVITY  
PEROVSKITES, FULLERENES, GRAPHENE,  
AND RELATED SYSTEMS

AUG 27 - SEPT 1, 2016 TARTU, ESTONIA

Website: [www.jahn-teller.ee](http://www.jahn-teller.ee)  
(1500-2100)

### Jahn-Teller Crystals – New Class of Smart Materials

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Smart materials are materials with properties that undergo significant and quick changes under environmental disturbances such as changes of temperature, pressure (hydrostatic and uniaxial), external magnetic and electric fields, light intensity, pH, and so on. Most of the material properties are defined by their electronic spectra. The electronic spectra undergo the most dramatic changes due to environmental forces when there is a degeneracy or pseudo degeneracy (small energy gaps) of the electronic states of the materials. In these materials, as a rule, the vibronic (electron-phonon) interactions are significant. All that leads to the conclusion that the Jahn-Teller systems are good candidates in the search for smart materials.

Different electronic and structural phase transitions are typical in the multicenter Jahn-Teller systems – Jahn-Teller crystals. If one accepts the fact that phase transitions are characterized by anomalously big material property changes, it becomes clear that Jahn-Teller multiferroics are especially good candidates for new smart materials and their applications.

While among the smart Jahn-Teller crystals there are representatives of all types of materials – metals, dielectrics, semiconductors, and superconductors, this presentation will concentrate on dielectric materials. Various old and new unusual smart properties will be analyzed.

Examples of practical applications of smart Jahn-Teller materials and corresponding patents will be discussed.

1. M. Kaplan, B. Vekhter, "Cooperative Phenomena in Jahn-Teller Crystals", Plenum Press, New York/London (1995).
2. M. Kaplan, G. Zimmerman, "Elastic Anomalies in Jahn-Teller Crystals with Competing Orderings", J. of Phys.: Conf. Ser., **428**, 012033 (2013).
3. M. Kaplan, "Jahn-Teller Crystals – New Class of Smart Materials", 4<sup>th</sup> Annual World Congress of Advanced Materials, p.81, Chongqing, China (2015).

### Two-level behavior of Ti<sup>3+</sup> *d*-orbitals in NaTiSi<sub>2</sub>O<sub>6</sub>

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NaTiSi<sub>2</sub>O<sub>6</sub> is a spin 1/2 member of quasi one-dimensional magnets of the pyroxene family. It is a Mott insulator, which crystallizes in the monoclinic space group *C2/c* at high temperatures [1, 2]. Its crystal lattice consists of slightly distorted TiO<sub>6</sub> octahedra, which are skew edge connected into zigzag chains, bridged by SiO<sub>4</sub> tetrahedra. In coordinates used in [3], the threefold *t<sub>2g</sub>* level of Ti<sup>3+</sup> splits due to distortion and crystal field into low-lying doublet with *d<sub>xy</sub>* and *d<sub>yz</sub>* orbitals and a single *d<sub>xz</sub>* orbital removed to 0.3 eV higher energy. This orbital degeneracy causes cooperative Jahn-Teller (JT) effect through strong electron-lattice coupling [3]. At *T<sub>c</sub>* = 210 K the lattice undergoes structural phase transition into triclinic *P1* space group [1, 2] with dimerization of Ti zigzag chains. The phase transition is actually orbital ordering and the creation of *d<sub>xy</sub>* dimers with spin singlet ground state, which is the orbital analogue of spin-Peierls transition. As a result a spin-gap of order 500-700 K [1, 4] and 53 meV [5] opens and the magnetic susceptibility starts rapidly to decrease. Theoretical calculation for the pyroxene family magnets in [6] gives for Ti-Ti dimer in NaTiSi<sub>2</sub>O<sub>6</sub> antiferromagnetic exchange coupling *J<sub>intra</sub>* = 396 K and weak exchange between the dimers *J<sub>inter</sub>* = 5K, close to the experimental data.

In order to gain more insight on local spin dynamics of the compound, we have performed <sup>29</sup>Si NMR spin-lattice relaxation study on NaTiSi<sub>2</sub>O<sub>6</sub>, [7]. We have found that at low temperature the magnetization recovery curve is nonexponential, with initial increase of the magnetization as square root of the time, which is characteristic to relaxation on paramagnetic impurities with no spin diffusion [8, 9]. Detailed analysis of the magnetization recovery gives a concentration of such impurities [9], which are in the present case mostly undimerized Ti<sup>3+</sup> ions. We have found that temperature dependence of this concentration *c(T)* follows activation type behaviour *c(T) = c<sub>0</sub> exp(-E<sub>a</sub>/T)*, with the energy *E<sub>a</sub>* = 300 ± 20K. We ascribe this activation energy to the splitting of the lower *d*-orbital doublet. We note that *E<sub>a</sub>* is very close to phonon Raman modes 221 and 209 cm<sup>-1</sup>, which anti-cross at *T<sub>c</sub>*, when the dynamic JT phase changes into the static one [3]. We show that the concentration of the unpaired Ti<sup>3+</sup> ions and the magnetic susceptibility data can be well described within a two level model, with 300 K splitting between the levels and the additional singlet-triplet energy gap of order 650 K.

1. M. Isobe *et al.*, Novel phase transition in spin-1/2 linear chain systems: NaTiSi<sub>2</sub>O<sub>6</sub> and LiTiSi<sub>2</sub>O<sub>6</sub>, J. Phys. Soc. Japan **71**, 1423 (2002).
2. G. J. Redhammer, *et al.*, Single-crystal structure refinement of NaTiSi<sub>2</sub>O<sub>6</sub> clinopyroxene at low temperatures (298 < T < 100K), Acta Crystallogr., sect. B: Struct. Sci **B59**, 730 (2003).
3. M. J. Konstantinovic *et al.*, Orbital dimerization in NaTiSi<sub>2</sub>O<sub>6</sub>: An orbital analogue of the spin-Peierls phase transition, Phys. Rev. **B69**, 020409(R) (2004).
4. P. J. Baker *et al.*, Muon-spin relaxation on the dimerized spin-1/2 chains NaTiSi<sub>2</sub>O<sub>6</sub> and TiOCl, Phys. Rev. **B75**, 094404 (2007).
5. H. J. Silverstein *et al.*, Direct measurement of the spin gap in a quasi-one-dimensional clinopyroxene: NaTiSi<sub>2</sub>O<sub>6</sub>, Phys. Rev. **B90**, 140402(R) (2014).
6. S. V. Streltsov and D. I. Khomskii, Electronic structure and magnetic properties of pyroxenes (Li,Na)TM(Si,Ge)<sub>2</sub>O<sub>6</sub>: Low-dimensional magnets with 90° bonds, Phys. Rev. **B77**, 064405 (2008).
7. R. Rästa *et al.*, <sup>29</sup>Si NMR study of NaTiSi<sub>2</sub>O<sub>6</sub>, this conference.
8. W. E. Blumberg, Nuclear spin-lattice relaxation caused by paramagnetic impurities, Phys. Rev. **119**, 79 (1959).
9. J. R. Bodart *et al.*, Recovery of nuclear magnetization under extreme inhomogeneous broadening, Phys. Rev. **B**, **54**, 15291 (1996).