

## Second version of Second Exam, Solutions PY541

1) The figure is drawn below. Note that there are 3 independent rotation modes and 3 vibration modes. (In general the number of vibration modes is  $3n-6$  for a non-collinear molecule of  $n$  atoms; we subtract off the 3 rotation modes and 3 center of mass translation modes from the  $3n$  degrees of freedom.)

2a) As one student observed, this problem is not physical. Orbital angular momentum must be an integer, so  $L=1/2$  is impossible. (I should have said  $L=1$ ,  $S=1/2$ , for example.) Nonetheless, it is straightforward to answer the question. The figure is drawn on the next page. At low  $T$  the entropy per molecule is  $k$  times the logarithm of the groundstate degeneracy which is zero. At high  $T$  it is  $k$  times the total number of state,  $k \ln 4$ . In between the entropy must increase since its derivative gives the heat capacity which is generally positive.

b)

$$Z_1 = 1 + 3e^{-\Delta/kT}. \quad (1)$$

(I ignore ground state energy which doesn't have any effect on the entropy.) Using the Helmholtz free energy,  $A = -kT \ln Z$ , I obtain the internal entropy per atom from these degrees of freedom:

$$S/N = -\frac{\partial A}{\partial T} = k \ln \left( 1 + 3e^{-\Delta/kT} \right) + \frac{3\Delta}{T(e^{\Delta/kT} + 3)} \quad (2)$$

Noting that the argument of the logarithm becomes 1 at low  $T$  and 4 at high  $T$  and that the second term is exponentially small at low  $T$  and drops off as  $\Delta/T$  at high  $T$  we see that this has the low and high  $T$  limits of the figure drawn below.

3) The maximum density which is possible without macroscopic occupation of the groundstate, achieved when  $\mu = 0$ , is:

$$n_e(T) = \int \frac{d^2p}{(2\pi\hbar)^2} \frac{1}{e^{pc/kT} - 1}. \quad (3)$$

This integral is finite since it behaves as:

$$\int \frac{dp}{p} \quad (4)$$

at small  $p$ , a finite integral. Therefore, when the actual density exceeds  $n_e(T)$ , Bose condensation must occur. Thus  $T_c$  is determined by the condition:  $n = n_e(T_c)$  where the function  $n_e(T)$  is determined by Eq. (3). We may reduce this to a dimensionless integral by a rescaling. Defining  $x = pc/kT$ , Eq. (3) becomes:

$$n_e(T) = (kT/2\pi\hbar c)^2 2\pi \int_0^\infty \frac{dx x}{e^x - 1}. \quad (5)$$

Letting  $I$  stand for this dimensionless integral:

$$I \equiv \int_0^\infty \frac{dx x}{e^x - 1}, \quad (6)$$

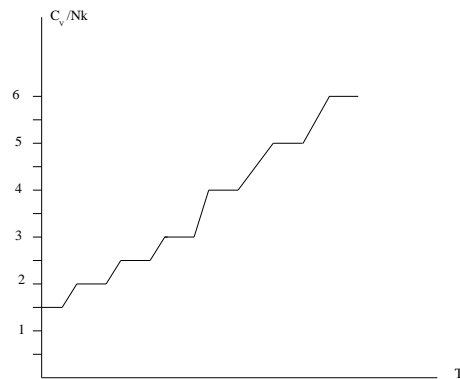


FIG. 1:

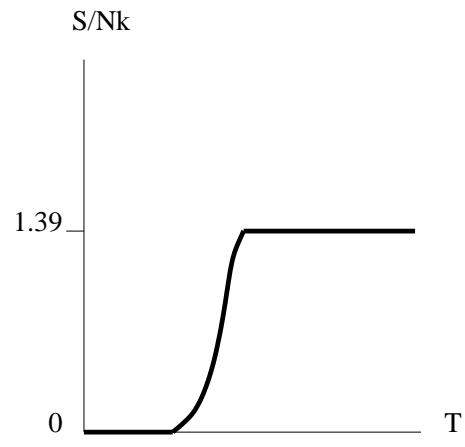


FIG. 2:

we obtain:

$$T_c = \sqrt{2\pi n/I}(\hbar c/k) \quad (7)$$