

Problem Set 6: Solutions

1) For the diagrams see graphs in Fig.1. The expansion involves:

a) 12 diagrams like graph1 where

$$(a,b,c,d) = (1,2,3,4); (1,3,2,4); (1,3,4,2); (1,4,3,2); (2,4,1,3); (2,1,4,3); (3,1,2,4); (3,2,1,4); (1,2,4,3); (1,4,2,3); (2,1,3,4); (2,3,1,4)$$

b) 4 diagrams like graph2 where

$$(a,b,c,d) = (2,1,4,3); (3,2,1,4); (2,3,1,4); (3,4,2,1)$$

c) 12 diagrams like graph3 where

$$(a,b,c,d) = (1,2,3,4); (1,3,2,4); (1,4,3,2); (2,3,1,4); (2,4,1,2); (2,1,4,3); (3,1,2,4); (3,2,4,1); (3,4,1,2); (4,1,2,3); (4,2,3,1); (4,3,1,2)$$

d) 3 diagrams like graph4 where

$$(a,b,c,d) = (1,2,4,3); (1,3,2,4); (1,2,3,4)$$

e) 6 diagrams like graph5 where

$$(a,b,c,d) = (1,3,2,4); (1,2,3,4); (1,2,4,3); (2,4,3,1); (2,3,4,1); (3,2,4,1)$$

f) 1 diagram like graph6 where

$$(a,b,c,d) = (1,2,4,3)$$

The sum S_4 of the clusters is

$$S_4 = 12V \left(\int d^3x f(\vec{x}) \right)^3 + 4V \left(\int d^3x f(\vec{x}) \right)^3 + \quad (1)$$

$$+ 12V \int d\vec{x}_1 f(\vec{x}_1) \int d\vec{x}_2 d\vec{x}_3 f(\vec{x}_2 - \vec{x}_3) f(\vec{x}_2) f(\vec{x}_3) + \quad (2)$$

$$+ 3V \int d\vec{x}_1 d\vec{x}_2 d\vec{x}_3 f(\vec{x}_2 - \vec{x}_1) f(\vec{x}_3 - \vec{x}_2) f(\vec{x}_1) f(\vec{x}_3) + \quad (3)$$

$$+ 6V \int d\vec{x}_1 d\vec{x}_2 d\vec{x}_3 f(\vec{x}_2 - \vec{x}_3) f(\vec{x}_1 - \vec{x}_3) f(\vec{x}_1) f(\vec{x}_2) f(\vec{x}_3) + \quad (4)$$

$$+ V \int d\vec{x}_1 d\vec{x}_2 d\vec{x}_3 f(\vec{x}_2 - \vec{x}_3) f(\vec{x}_1 - \vec{x}_3) f(\vec{x}_1 - \vec{x}_2) f(\vec{x}_1) f(\vec{x}_2) f(\vec{x}_3) \quad (5)$$

and

$$b_4 = \frac{S_4}{V \lambda^{12} 4!}$$

2) As shown in class,

$$\left(\frac{\partial T}{\partial P} \right)_H = \frac{1}{C_P} \left[T \left(\frac{\partial V}{\partial T} \right)_P - V \right]$$

where,

$$\left(\frac{\partial V}{\partial T} \right)_P = \frac{\partial(V, P)}{\partial(T, P)} = \frac{\partial(V, P)}{\partial(V, T)} \frac{\partial(V, T)}{\partial(T, P)} = - \frac{(\partial P / \partial T)_V}{(\partial P / \partial V)_T}$$

Now, $P = nkT [1 - b_2 \lambda^3 n]$ then

$$\left(\frac{\partial P}{\partial T} \right)_V = nk \left[1 - b_2 \lambda^3 n - \frac{d(b_2 \lambda^3)}{dT} T n \right]$$

and

$$\left(\frac{\partial P}{\partial V}\right)_T = -\frac{nkT}{V} + \frac{2n^2 b_2 \lambda^3 kT}{V}$$

so,

$$\left(\frac{\partial V}{\partial T}\right)_P = \frac{nk \left[1 - b_2 \lambda^3 n - \frac{d(b_2 \lambda^3)}{dT} T n \right]}{\frac{nkT}{V} [1 - 2nb_2 \lambda^3]} = \frac{V}{T} \left[1 + b_2 \lambda^3 n - \frac{d(b_2 \lambda^3)}{dT} T n \right]$$

Then,

$$T \left(\frac{\partial V}{\partial T} \right)_P - V = V \left[b_2 \lambda^3 n - \frac{d(b_2 \lambda^3)}{dT} T n \right]$$

and

$$\left(\frac{\partial T}{\partial P}\right)_H = \frac{n}{C_P} \left[b_2 \lambda^3 - \frac{d(b_2 \lambda^3)}{dT} T \right]$$

3) a) It is convenient to assume that particles feel the same potentials at walls, so that so particle can get closer to either wall than r_0 (this assumption does not matter at $L \rightarrow \infty$).

$$W_N = N! \int_{-L/2+r_0}^{L/2-Nr_0} dx_N \dots \int_{x_4+r_0}^{L/2-3r_0} dx_3 \int_{x_3+r_0}^{L/2-2r_0} dx_2 \int_{x_2+r_0}^{L/2-r_0} dx_1$$

where we assume $x_1 > x_2 > \dots > x_N$. The $N!$ factor takes care of all other orders. Then,

$$\int_{x_2+r_0}^{L/2-r_0} dx_1 = \frac{L}{2} - 2r_0 - x_2$$

$$\int_{x_3+r_0}^{L/2-2r_0} \left(\frac{L}{2} - 2r_0 - x_2 \right) dx_2 = \frac{1}{2} \left(\frac{L}{2} - 2r_0 - x_3 \right)^2$$

$$\int_{x_4+r_0}^{L/2-3r_0} \frac{1}{2} \left(\frac{L}{2} - 2r_0 - x_3 \right)^2 dx_3 = \frac{1}{3!} \left(\frac{L}{2} - 4r_0 - x_4 \right)^3$$

and so on. Finally we get,

$$W_N = (L - (N+1)r_0)^N \approx (L - Nr_0)^N$$

Therefore,

$$Z_N = \frac{(L - Nr_0)^N}{N! \lambda^N}$$

Then,

$$A = -kT \ln(Z_N) \approx -kTN \ln \left[\frac{e(L - Nr_0)}{N\lambda} \right]$$

where I used Stirling's formula.

$$P = - \left(\frac{\partial A}{\partial L} \right)_{N,T} = \frac{kTN}{L - Nr_0} = \frac{kTN}{L} \left[1 + \frac{N}{L} r_0 + \left(\frac{N}{L} r_0 \right)^2 + \dots \right]$$

b) $f(x) = e^{-\beta u(x)} - 1$, then $f(x) = -1$ for $|x| < r_0$ and $f(x) = 0$ for $|x| > r_0$.

$$b_2 = \frac{1}{2\lambda} \int_{-\infty}^{\infty} f(x) dx = -\frac{r_0}{\lambda}$$

$$b_3 - 2b_2^2 = \frac{1}{3!\lambda^2} \int f(x_1) f(x_2 - x_1) f(x_2) dx_1 dx_2 = -\frac{1}{3!\lambda^2} \int_{-r_0}^{r_0} dx_1 \int_{-r_0}^{r_0} dx_2 \theta(r_0 - |x_1 - x_2|) = -\frac{1}{2} \left(\frac{r_0}{\lambda}\right)^2.$$

So,

$$[1 - b_2 n \lambda + (4b_2^2 - 2b_3) n^2 \lambda^2] = [1 + r_0 n + r_0^2 n^2]$$

4)

$$b_2 = \frac{4\pi}{2\lambda^3} \int r^2 [e^{-\beta u(r)} - 1] dr = \quad (6)$$

$$= \frac{2\pi}{\lambda^3} \left[- \int_0^{r_0} r^2 dr + \int_{r_0}^{r_1} r^2 [e^{\beta u_0} - 1] dr \right] = \quad (7)$$

$$= \frac{2\pi}{\lambda^3} \left[-\frac{r_0^3}{3} + (e^{\beta u_0} - 1) \frac{1}{3} (r_1^3 - r_0^3) \right] = \quad (8)$$

$$= \frac{2\pi}{3\lambda^3} [-r_0^3 e^{\beta u_0} + r_1^3 (e^{\beta u_0} - 1)] \quad (9)$$

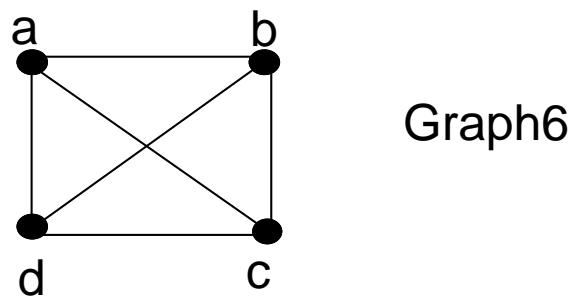
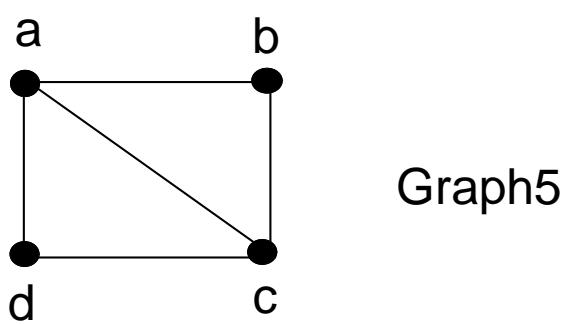
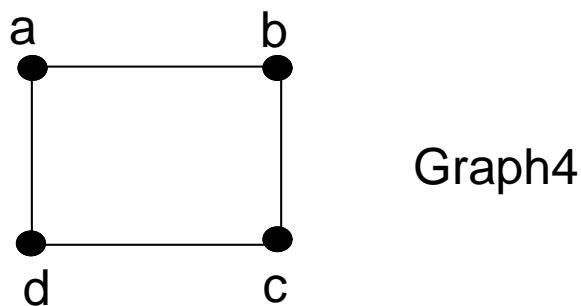
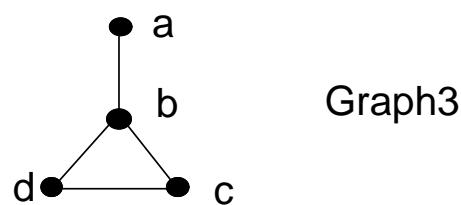
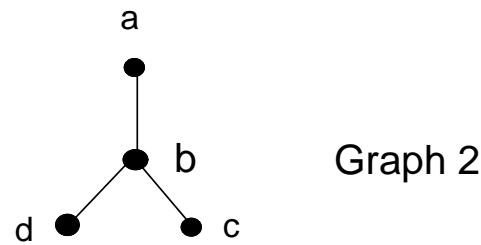


FIG. 1: (a) Graphs for expansion in Problem 1.