

“Measuring” physical observables

Order parameter of ferromagnetic transition: Magnetization

$$M = \sum_{i=1}^N \sigma_i, \quad m = \frac{M}{N}$$

Expectation vanishes for finite system; calculate $\langle |m| \rangle$, $\langle m^2 \rangle$

Susceptibility: Linear response of $\langle m \rangle$ to external field

$$E = E_0 - hM, \quad E_0 = J \sum_{i,j} \sigma_i \sigma_j$$
$$\chi = \left. \frac{d\langle m \rangle}{dh} \right|_{h=0} \quad (\text{we can also consider } h > 0 \text{ here})$$

Deriving Monte Carlo **estimator**

$$\langle m \rangle = \frac{1}{Z} \sum_S m e^{-(E_0 - hM)/T}, \quad Z = \sum_S e^{-(E_0 - hM)/T}$$

$$\chi = -\frac{dZ/dh}{Z^2} \sum_S m e^{-(E_0 - hM)/T} + \frac{1}{Z} \frac{1}{T} \sum_S m M e^{-(E_0 - hM)/T}$$

$$\frac{dZ}{dh} = \frac{1}{T} \sum_S M e^{-(E_0 - hM)/T}$$

$$\chi = \frac{1}{N} \frac{1}{T} (\langle M^2 \rangle - \langle M \rangle^2) = \frac{1}{N} \frac{1}{T} \langle M^2 \rangle, \quad (h = 0)$$

Extrapolating to infinite size, this gives the correct result only in the disordered phase (gives infinite susceptibility for $T < T_c$)

We can also define the susceptibility estimator as

$$\chi = \frac{1}{N} \frac{1}{T} (\langle M^2 \rangle - \langle |M| \rangle^2)$$

Gives correct infinite-size extrapolation for any T

Specific heat

$$C = \frac{1}{N} \frac{dE}{dT} = \frac{1}{N} \frac{d}{dT} \sum_C E(C) e^{-E(C)/T} = \frac{1}{N} \frac{1}{T^2} (\langle E^2 \rangle - \langle E \rangle^2)$$

Correlation function

$$C(\vec{r}) = \langle \sigma_i \sigma_{j(\vec{r}, i)} \rangle$$

Average over all spins i

$$C(\vec{r}) = \frac{1}{N} \sum_{i=1}^N \langle \sigma_i \sigma_{j(\vec{r}, i)} \rangle$$

Statistical errors (“error bars”)

Calculation based on M “bins”. What is the statistical error?

Consider M independent calculations (each based on n configs)

Statistically independent averages $\bar{A}_i, i = 1, \dots, M$

Full average

$$\bar{A} = \frac{1}{M} \sum_{i=1}^M A_i$$

Standard deviation

$$\sigma' = \sqrt{\frac{1}{M} \sum_{i=1}^M (\bar{A}_i^2 - \bar{A}^2)}$$

But, we want the standard deviation of the average

$$\sigma = \sqrt{\frac{1}{M(M-1)} \sum_{i=1}^M (\bar{A}_i^2 - \bar{A}^2)}$$

The bins have to be long enough (# of MC steps, n , large enough) to be essentially statistically independent (can be quantified by “autocorrelations” - later)

Simulations with an external magnetic field

$$E = -J \sum_{\langle i,j \rangle} \sigma_i \sigma_j - h \sum_i \sigma_i$$

here $J > 0$ for ferromagnet
-we added minus sign in front
-a matter of taste...

For $h > 0$, the average magnetization $\langle M \rangle > 0$

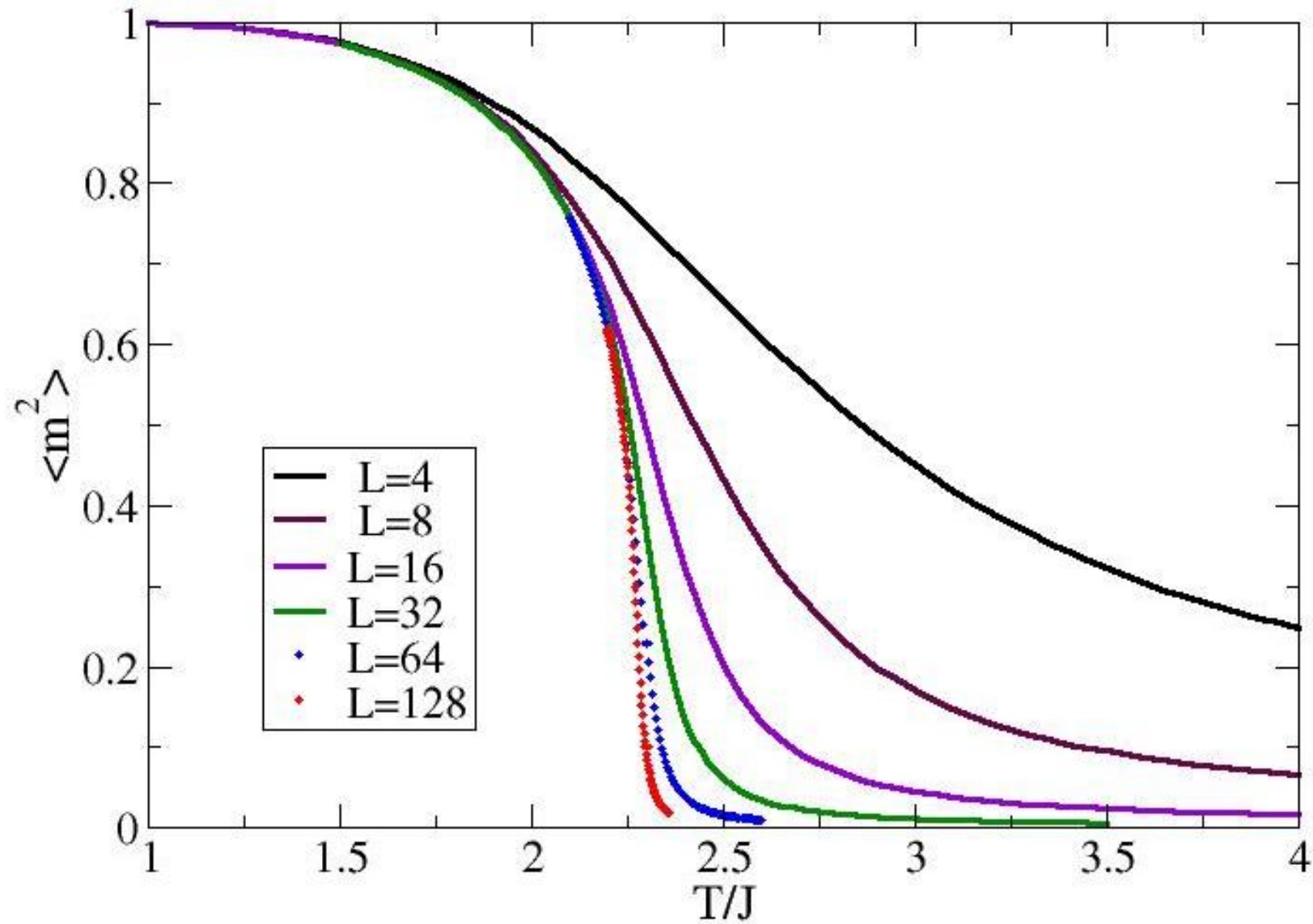
Simple change in the acceptance probability

$$P(S \rightarrow \tilde{S}_j) = \min \left[\frac{W(\tilde{S}_j)}{W(S)}, 1 \right]$$

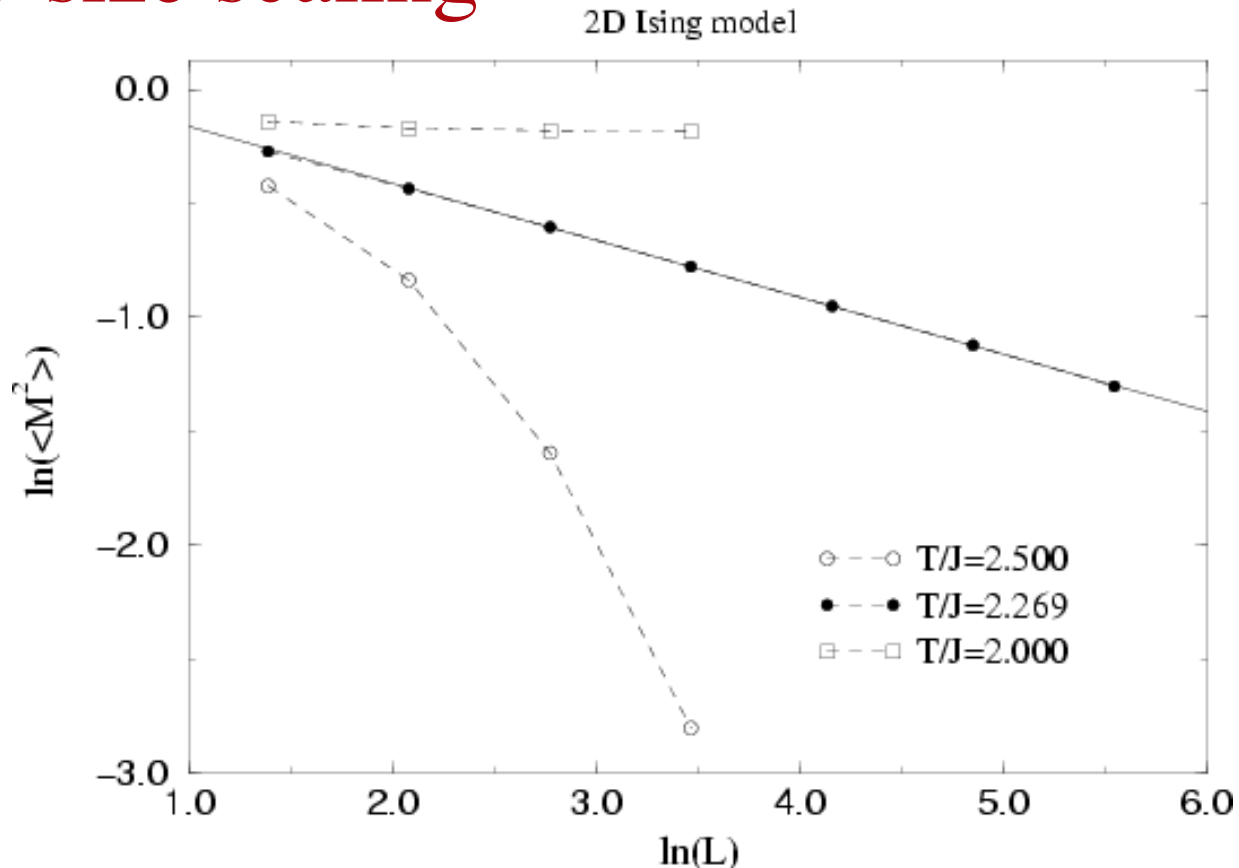
$$\frac{W(\tilde{S}_j)}{W(S)} = \exp \left[-\frac{2J}{T} \sigma_j \left(\sum_{\delta(j)} \sigma_{\delta(j)} - h \right) \right]$$

Look at the program `ising2d.jl`

Squared magnetization for different system sizes (no external field): development of phase transition



Finite-size scaling



$T > T_c$: $\langle M^2 \rangle \rightarrow 0$ (as $1/L^2$, trivial from short-range correlations)

$T = T_c$: $\langle M^2 \rangle \rightarrow 0$ (non-trivial power law)

$T < T_c$: $\langle M^2 \rangle \rightarrow \text{constant} > 0$

Extracting an exponent: $A = aL^\alpha \longrightarrow \ln(A) = \ln(a) + \alpha \ln(L)$

-Power-law: straight line when plotted on log-log scale

-The $1/L^2$ form for $T/J=2.5$ not yet seen because of cross-over behavior; close to the critical point, larger L required

Critical behavior and scaling

Correlation length ξ defined in terms of correlation function

$$C(\vec{r}_{ij}) = \langle \sigma_i \sigma_j \rangle - \langle \sigma_i \rangle^2 \sim e^{-r_{ij}/\xi}, \quad \vec{r}_{ij} \equiv |\vec{r}_i - \vec{r}_j|$$

The correlation length diverges at the critical point

$$\xi \sim t^{-\nu}, \quad t = \frac{|T - T_c|}{T_c} \quad (\text{reduced temperature})$$

ν is an example of a **critical exponent**

Universality

Critical exponents do not depend on microscopic details of the interactions; only on the dimensionality of the system and the order parameter:

- Ising, gas/liquid (scalar Z_2 -symmetric order parameter)
- XY spins, phase of superconductor (2D, $O(2)$ order parameter)
- Heisenberg spins (3D, $O(3)$ order parameter)

Phase transitions fall into **universality classes** characterized by different sets of critical exponents

Other critical exponents

Order parameter for $T < T_c$ (e.g., magnetization)

$$\langle m \rangle \sim (T_c - T)^\beta$$

In practice, calculate $\langle |m| \rangle$, $\langle m^2 \rangle$

Susceptibility corresponding to order

$$\chi = \frac{1}{N} \frac{1}{T} (\langle M^2 \rangle - \langle |M| \rangle^2)$$

Diverges at the critical point

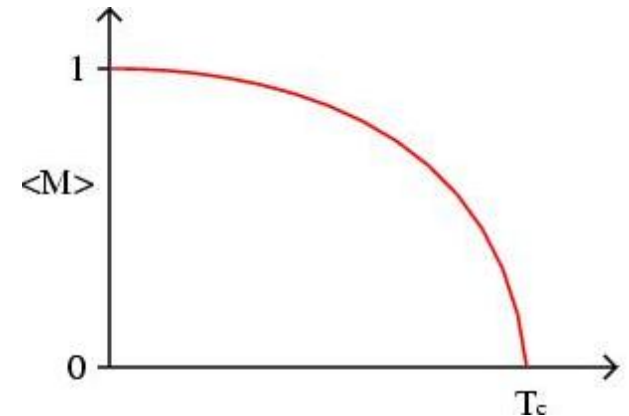
$$\chi \sim t^{-\gamma}$$

Specific heat $C = \frac{1}{N} \frac{1}{T^2} (\langle E^2 \rangle - \langle E \rangle^2)$

Singular at T_c

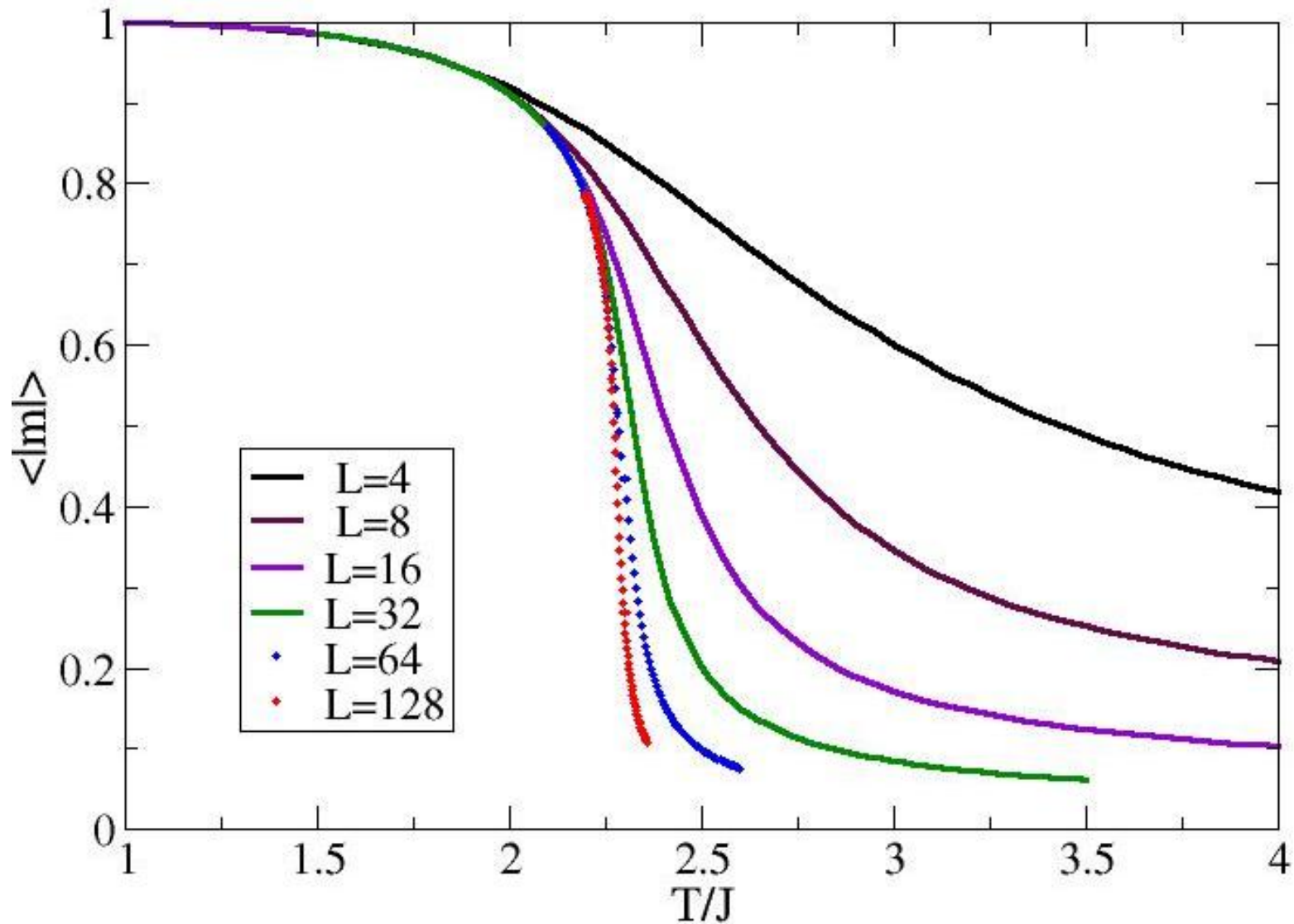
$$C \sim t^{-\alpha}$$

The exponent α can be positive or negative (no divergence if negative; 0 can correspond to log divergence)

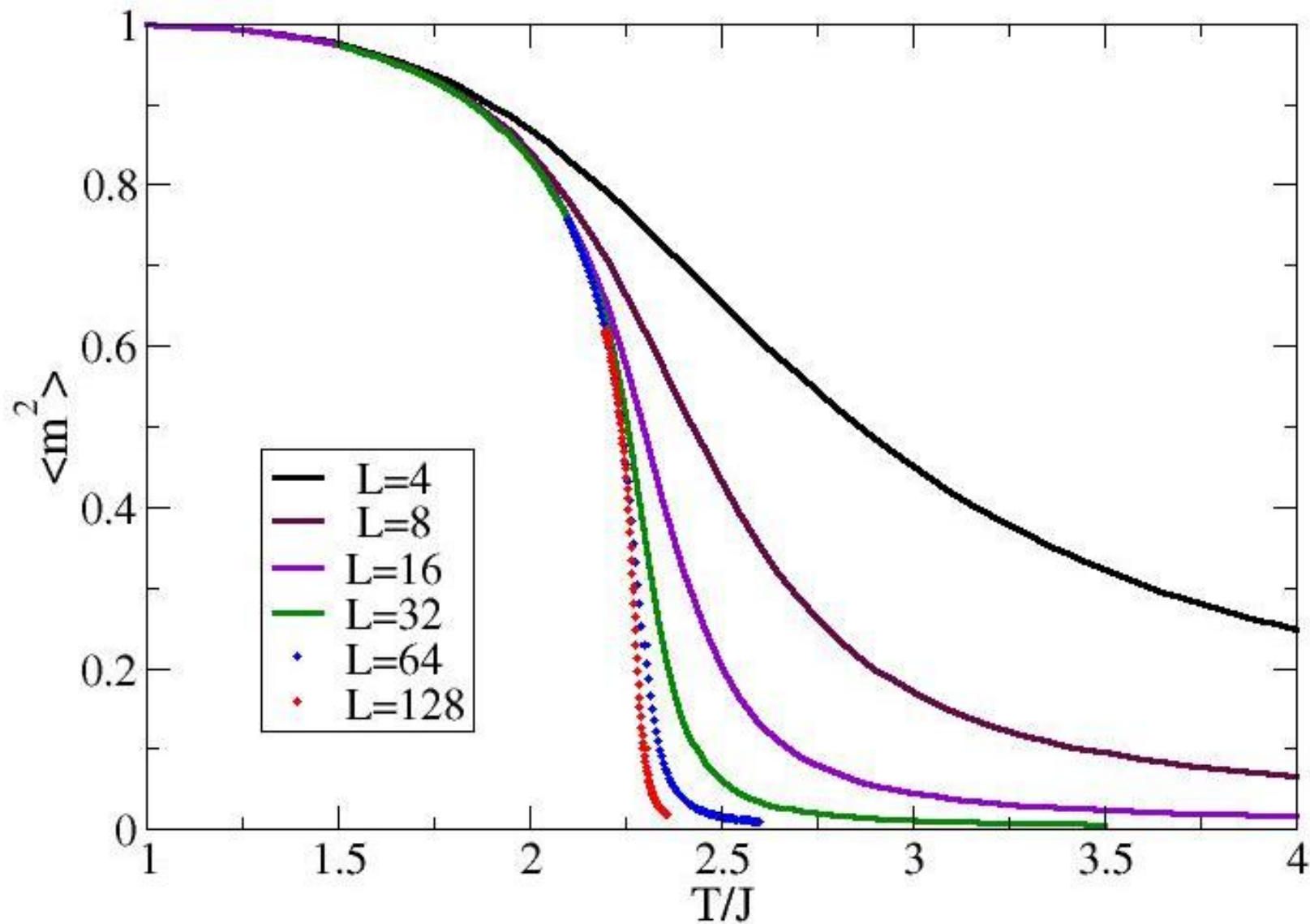


Magnetization of 2D Ising ferromagnet

$\langle |m| \rangle \sim (T_c - T)^\beta$, $(T < T_c)$ for infinite system



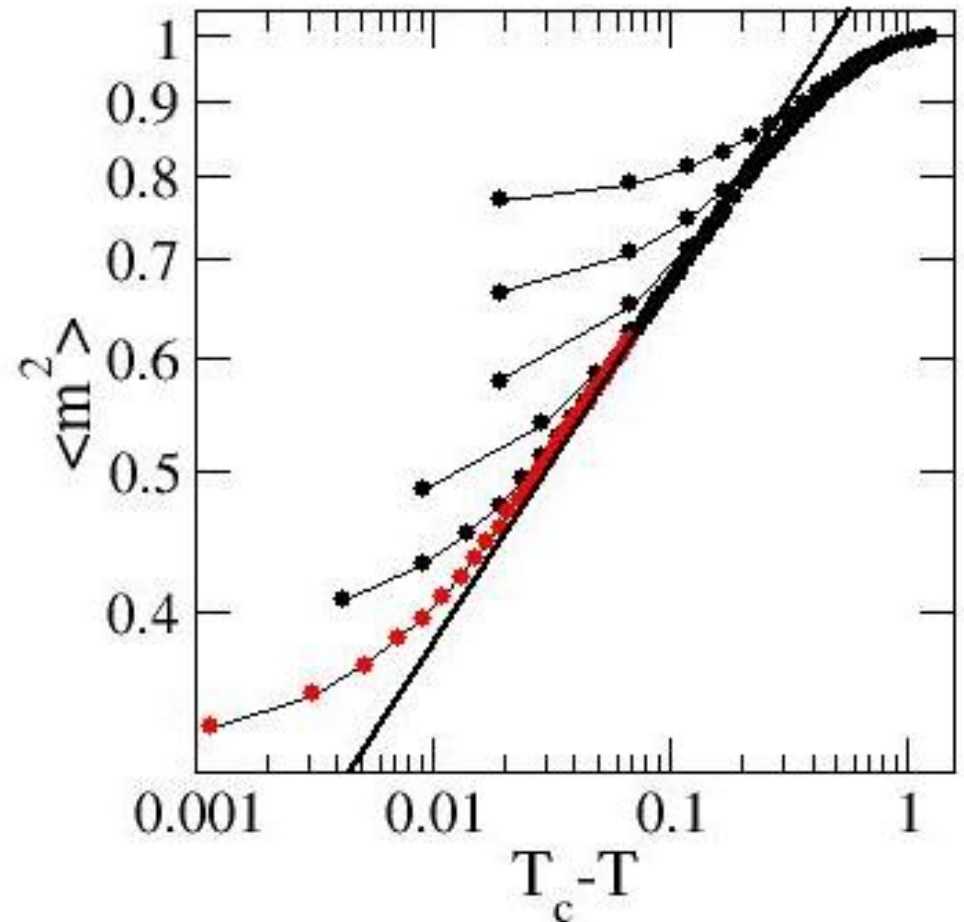
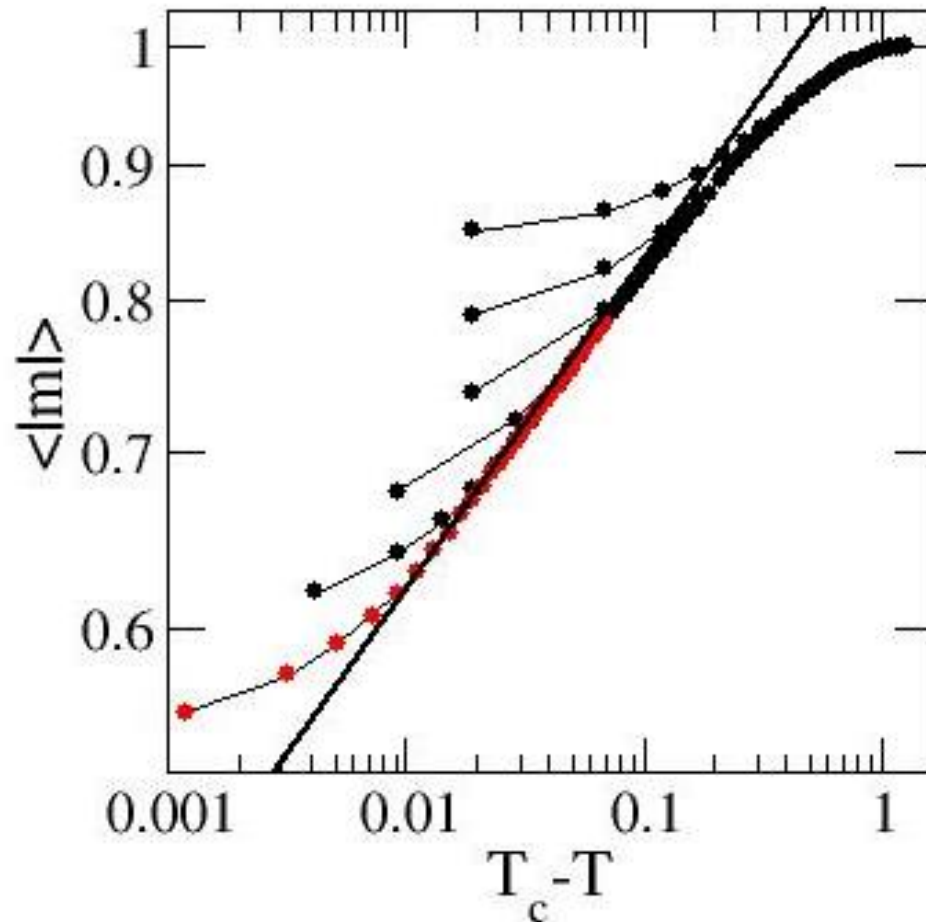
Magnetization squared $\langle m^2 \rangle \sim (T_c - T)^{2\beta}$, $(T < T_c)$



The exponent β can be extracted for large L

Comparison with known 2D Ising model exponent

$$\beta = 1/8$$



If T_c is not known, use it as an adjustable parameter and look for power-law behavior

Finite-size scaling

For a system of length L , the correlation length $\xi \leq L$

Express divergent quantities in terms of correlation length, e.g.,

$$\xi \sim t^{-\nu}, \quad \chi \sim t^{-\gamma} \sim \xi^{\gamma/\nu}$$

The largest value is obtained by substituting $\xi \rightarrow L$

$$\chi_{\max} \sim L^{\gamma/\nu}$$

At what T does the maximum occur?

$$\xi = at^{-\nu} = L \Rightarrow t \sim L^{-1/\nu}$$

The peak position of a divergent quantity can be taken as T_c for finite L (different quantities will give different T_c)

γ, ν can be extracted by studying peaks in $\xi(T)$

Similarly for specific heat;

$$C_{\max} \sim L^{\alpha/\nu}$$