

Selected topics in High-energy QCD physics

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QCD - Features

 $\mathcal{L}_{QCD} = \bar{\psi} \left[i \gamma^{\mu} \partial_{\mu} - m \right] \psi - g_s \bar{\psi} \gamma^{\mu} G^a_{\mu} \frac{\lambda_a}{2} \psi - \frac{1}{4} G^a_{\mu\nu} G^{\mu\nu}_a$

- $G^a_{\mu\nu} = \partial_\mu G^a_\nu \partial_\nu G^a_\mu g_s f^a_{bc} G^b_\mu G^c_\nu$
- Interactions arise from fundamental symmetry principles: SU(3)_c
- Visible phenomena (e.g. proton) emerge through complex structure of the vacuum (e.g formation of Hadrons from quarks/gluons)
- Fundamental differences to QED:
 - O Self-interacting: highly nonlinear
 - O Interaction increases at large distances: Confinement
 - O Interaction decreases at small distances: Asymptotic freedom
 - O Strong coupling: as >>aem
 - O Topological excitations

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- QCD Profound differences between hadrons and other many-body systems
 - Atoms, molecules, nuclei,...:
 - O Constituents can be removed
 - Exchanged boson generating interaction may be subsumed into static potential (e.g. photon into Coulomb potential)
 - Most of mass from fermion constituents

• Nucleons:

- O Quarks are confined
- O Gluons are essential degrees of freedom
- Most of mass generated by interactions (~99%)
- Exploration of QCD: Analytical (e.g. perturbative) / non-perturbative (e.g. Lattice QCD -Numerical solution on space-time lattice) methods in comparison to experimental results

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matters



QCD - Perturbative Side



Photo PRAS



Discovery of asymptotic freedom in the theory of strong interaction (Quantum Chromo Dynamics): Nobel prize in physics 2004



4 tracks 4.1 GeV 4.3 GeV 4.3 GeV 4.3 GeV 4.3 GeV 4.3 GeV 4.3 GeV

TASSO Collaboration, R. Brandelik et al., Phys. Lett. B 86 (1979) 243.

O as large at large distances (low energy)

 \circ as small at small distances (high energy)



- QCD Non-Perturbative Side
- Lattice QCD: Numerical solution of path integrals on space-time lattice
- Successful description of various hadron properties (e.g. mass spectrum in the context of lattice QCD calculations)
- For a large class of problems (e.g. hadron formation from quarks), phenomenological methods and modeling are indispensable

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Y. Kuramashi, Lattice 2007, PACS-CS Collaboration



- QCD Questions (The Frontiers of Nuclear Science Long Range Plane 2007)
 - What is the internal landscape of the nucleons?
 - What does QCD predict for the properties of strongly interacting matter?
 - What governs the transition of quarks and gluons into hadrons?
 - What is the role of gluons and gluon selfinteractions in nucleons and nuclei?
 - What are the phases of strongly interacting matter?
 - What determines the key features of QCD?

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http://www.er.doe.gov/np/nsac/docs/ Nuclear-Science.Low-Res.pdf



Exploring the proton structure and dynamics



Structure and dynamics of proton (mass) (\rightarrow visible universe) originates from QCD-interactions!

What about spin as another fundamental quantum number?

Synergy of experimental progress and theory (Lattice QCD / Phenomenology incl. phenomenological fits / Modeling) critical!



Mass in QCD

Quote from Nobel prize lecture in physics, 2004, given by Professor Frank Wilczek (MIT):

Stated as $m=E/c^2$: Possibility of explaining mass in terms of energy.

Einstein's original paper does not contain the equation E=mc², but rather m=E/c²: "Does the Inertia of a Body Depend Upon its Energy Content? "(A. Einstein, Annalen der Physik, 18 (1905) 639.)"

Modern QCD: Mass of ordinary matter derives almost entirely from energy - the energy of massless gluons and nearly massless quarks, which are the ingredients from which protons, neutrons, and atomic nuclei are made.

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Spin in QCD



- Traditional way to introduce spin in QM textbooks: Stern-Gerlach experiment (1922)
- Concept of spin: Long and tedious battle to understand splitting patters and separations in line spectra
- Anomalous magnetic moment of proton by Stern et al. (1933)

Proposal of self-rotating electron by Goudsmit and Uhlenbeck (1925):



How do we probe the structure and dynamics of matter in ep / pp scattering?



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Fundamental QCD ingredients

• Asymptotic freedom:

 $\alpha_s \rightarrow 0$ at short distances: \Rightarrow perturbative QCD

 α_s large at long distances: \Rightarrow non-perturbative QCD

• Factorization: hard scale Q^2 , m_c , m_b



• Evolution:

Beyond Quark-Parton model,
 Parton densities become
 functions of Q²

 Predict Q² dependence of parton distribution functions (DGLAP evolution equations)

 $\sigma^{ep} = \gamma(x, Q^2) \otimes f_j(x, Q^2) \otimes \hat{\sigma}(x, Q^2)$

non-perturbative part

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Evolution

The presence of QCD related diagrams leads to a modification of F_2 0

$$\frac{F_{2}(x,Q^{2})}{x} = \sum_{i} Q_{i}^{2} \int_{0}^{1} \left(\frac{dy}{y}\right) q(y) \left(\delta\left(1-\frac{x}{y}\right) + \left(\frac{\alpha_{s}}{2\pi}\right) P_{qq}\left(\frac{x}{y}\right) \log \frac{Q^{2}}{\mu^{2}}\right) \text{ Logarithmic violation of scaling}$$

$$q(y) \equiv f_{q}(y) \quad \text{Parton model} \quad \text{Gluon radiation} \quad \text{Splitting function}$$

$$\frac{F_{2}(x,Q^{2})}{x} = \sum_{i} Q_{i}^{2} \int_{0}^{1} \left(\frac{dy}{y}\right) \left(q(y) + \Delta q(y,Q^{2})\right) \delta\left(1-\frac{x}{y}\right) = \sum_{i} Q_{i}^{2} \left(q(x) + \Delta q(x,Q^{2})\right)$$

$$\sum_{i} Q_{i}^{2} \left(q(x) + \Delta q(x,Q^{2})\right) \quad \Delta q(x,Q^{2}) = \left(\frac{\alpha_{s}}{2\pi}\right) \log \left(\frac{Q^{2}}{\mu^{2}}\right) \int_{x}^{1} \left(\frac{dy}{y}\right) q(y) P_{qq}\left(\frac{x}{y}\right)$$

depend on x and Q^2 :

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 $F_2 = \nu W_2$

 $\nu =$

 $\underline{p \cdot q}$



 n_f

 $\Sigma(x, Q^2) = \sum_{i=1} \left[q_i(x, Q^2) + \bar{q}_i(x, Q^2) \right]$

- Evolution of parton distribution functions (1)
 - Consider the change of the quark density $\Delta q(x,Q^2)$ over an interval of $\Delta \log Q^2$
 - General including other types of splitting functions:



$$\frac{d}{d\log Q^2}q(x,Q^2) = \left(\frac{\alpha_s}{2\pi}\right)\int_0^1 \left(\frac{dy}{y}\right)q(y,Q^2)P_{qq}\left(\frac{x}{y}\right)$$

 $g(x,Q^2)$

Probability of finding a parton of type i with momentum fraction x which originated from parton j having momentum fraction y!

Singlet distribution

Gluon distribution

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- Evolution of parton distribution functions (2)
 - Types of splitting functions



Probability of finding a parton of type i with momentum fraction x which originated from parton j having momentum fraction y!

$$\frac{d\Sigma(x,Q^2)}{d\ln Q^2} = \frac{\alpha_s}{2\pi} \int_x^1 \frac{dz}{z} \left[P_{qq}\left(\frac{x}{z}\right) \Sigma(z,Q^2) + P_{qg}\left(\frac{x}{z}\right) g(z,Q^2) \right]$$

 $\frac{dg(x,Q^2)}{d\ln Q^2} = \frac{\alpha_s}{2\pi} \int_x^1 \frac{dz}{z} \left[P_{gq}\left(\frac{x}{z}\right) \Sigma(z,Q^2) + P_{gg}\left(\frac{x}{z}\right) g(z,Q^2) \right]$

DGLAP evolution equations:

G. Altarelli and G. Parisi, Nucl. Phys. B 126 (1977) 298; V. Gribov and L.N. Lipatov, Soc. J. Nucl. Phys. 15 (1972)
438; L.N. Lipatov, Soc. J. Nucl. Phys. 20 (1975) 96; Y.L. Dokshitzer, Soc. Phys. JETP 46 (1977) 641.

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Global fits

- Determine F₂^{QCD} in terms of parton distribution functions
- Evolve F₂^{QCD} through parton distribution functions based on evolution equations
- Minimize χ^2 in terms of F_2^{QCD} and F_2^{data} by adjusting parameters in $xf_i(x,Q^2)$
- Net result: QCD prediction for xf_i(x,Q²) and therefore F₂(x,Q²)
- Various global pdf analysis:
 - GRV
 - CTEQ
 - MRST

NEPPSSR 2009 Craigville, MA, August 14, 2009 $xf_i(x, Q_0^2) = A_i x^{-\lambda_i} (1-x)^{\eta_i} F(x)$



Low x: λ_i

High x: η_i



Factorization

 Unpolarized proton structure:



 f_1, f_2

- Three step process:
 - Partons (quarks/gluons) in initial state: Long distance (non-perturbative QCD domain)

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- ⇒ Parton (quarks/gluons) distribution functions
- Hard interaction: Small distances (high energies) (perturbative QCD domain)
 - \Rightarrow Cross-section prediction (LO,NLO,NNLO)
- Quarks in final state: Long distance (non-perturbative QCD domain):
 - ⇒ Quarks fragment into observable hadrons described by fragmentation functions

 \boldsymbol{q}

e

Û

Overview

• Collider programs:

- Hadron-Hadron: Tevatron/ RHIC
- Electron-Hadron: HERA
- Electron/Positron: LEP

O QCD topics

- QCD factorization
- D Parton-distribution functions / Fragmentation functions
- Strong-coupling constant
- Jet algorithms
- QCD matrix elements in LO, NLO, NNLO
- Multi-leg final states
- Low-x physics
- Soft processes: Underlying event / Hadronization / Diffraction

Experimental QCD tests in ep

- \Box Measurement of α_s
- Fragmentation functions
- Extraction of parton distribution functions
- Color/spin dynamics
- Quark-Gluon jet properties
- Event shape variables (Sphericity, thrust, ...)
- Diffraction

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 $E_{e} = 27.5 \, GeV$ $E_p = 920 \text{ GeV}$



Experimental QCD tests in ee

- \Box Measurement of α_s
- Fragmentation functions
- Color/spin dynamics
- Quark-Gluon jet properties
- Event shape variables (Sphericity, thrust, ...)



LEP: Centre-of-mass energy (e^+e^-) up to 205GeV

BNL

- Experimental QCD tests in pp
 - \square Measurement of α_s
 - Fragmentation functions
 - Extraction of parton distribution
 functions
 - Color/spin dynamics
 - Quark-Gluon jet properties
 - Event shape variables (Sphericity, thrust, ...)
 - Diffraction

RHIC: Centre-of-mass energy 200-500GeV (polarized protons)





□ Precision measurements (e.g. F_2) \Rightarrow Precision on quark/gluon structure



- Precision on quark/gluon structure
 - Enormous precision reached over a wide kinematic region
 - Large uncertainties for all distribution functions at large momentum fractions:
 - Impact of W/Z program at
 LHC
 - □ Impact of high-E_T jet production



Cross Section Results - RHIC (Hadrons)



 Good agreement between data and NLO calculations for neutral pion production at forward and central rapidity



Cross Section Results - RHIC (Jets / Photons)



What do we know about the polarized quark and gluon distributions?



D. de Florian et al., Physpl 08.00.04,2094018 (2005).

$$\Delta G(Q^2) = \int_0^1 \Delta g(x, Q^2) dx$$

$$\Delta q_i(Q^2) = \int_0^1 \Delta q_i(x, Q^2) dx$$

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Gluon polarization - Extraction



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Gluon polarization - Correlation Measurements

• Correlation measurements provide access to partonic kinematics through Di-Jet/Hadron production and Photon-Jet production

$$x_{1(2)} = \frac{1}{\sqrt{s}} \left(p_{T_3} e^{\eta_3(-\eta_3)} + p_{T_4} e^{\eta_4(-\eta_4)} \right)$$

- O Di-Jet production / Photon-Jet production
 - Di-Jets: All three (LO) QCD-type processes contribute: gg, qg and qq with relative contribution dependent on topological coverage
 - Photon-Jet: One dominant underlying (LO) process with large partonic aLL at forward rapidity
 - Larger cross-section for di-jet production compared to photon related measurements
 - Photon reconstruction more challenging than jet reconstruction
 - \Box Full NLO framework exists \Rightarrow Input to Global analysis









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• GRSV-STD: Higher order QCD analysis of polarized DIS experiments!

 $\Delta G(Q^2) = \int_0^1 \Delta g(x, Q^2) dx$

□ STAR Inclusive Jet production - RUN 6: ALL ⇒ Gluon spin contribution



O RUN 6 results: GRSV-MAX / GRSV-MIN ruled out - A_{LL} result favor a gluon polarization in the measured x-region which falls in-between GRSV-STD and GRSV-ZERO

• Consistent with RUN 5 result (Factor 3-4 improved statistical precision for pt>13GeV/c)

Recent results in high-energy QCD _F



Run 9 STAR Beam-Use Request: Di-Jet projections



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M/√s

0.35

0.40



- Summary
 - Enormous precision reached for unpolarized distribution functions Large uncertainties remain at larger momentum fractions - Impact for LHC program
 - Higher-order QCD calculations needed (Beyond NLO) for precision pdf extraction and as
 - Generally large uncertainties and model dependence on soft processes such as underlying event and hadronization
 - Evidence for a small gluon polarization ⇒ Renaissance of constituent quark model!
- Outlook
 - Electron-Ion Collider: Precision measurement of polarized ep and eA scattering

