

Physics at the Tevatron Collider: Lecture #1

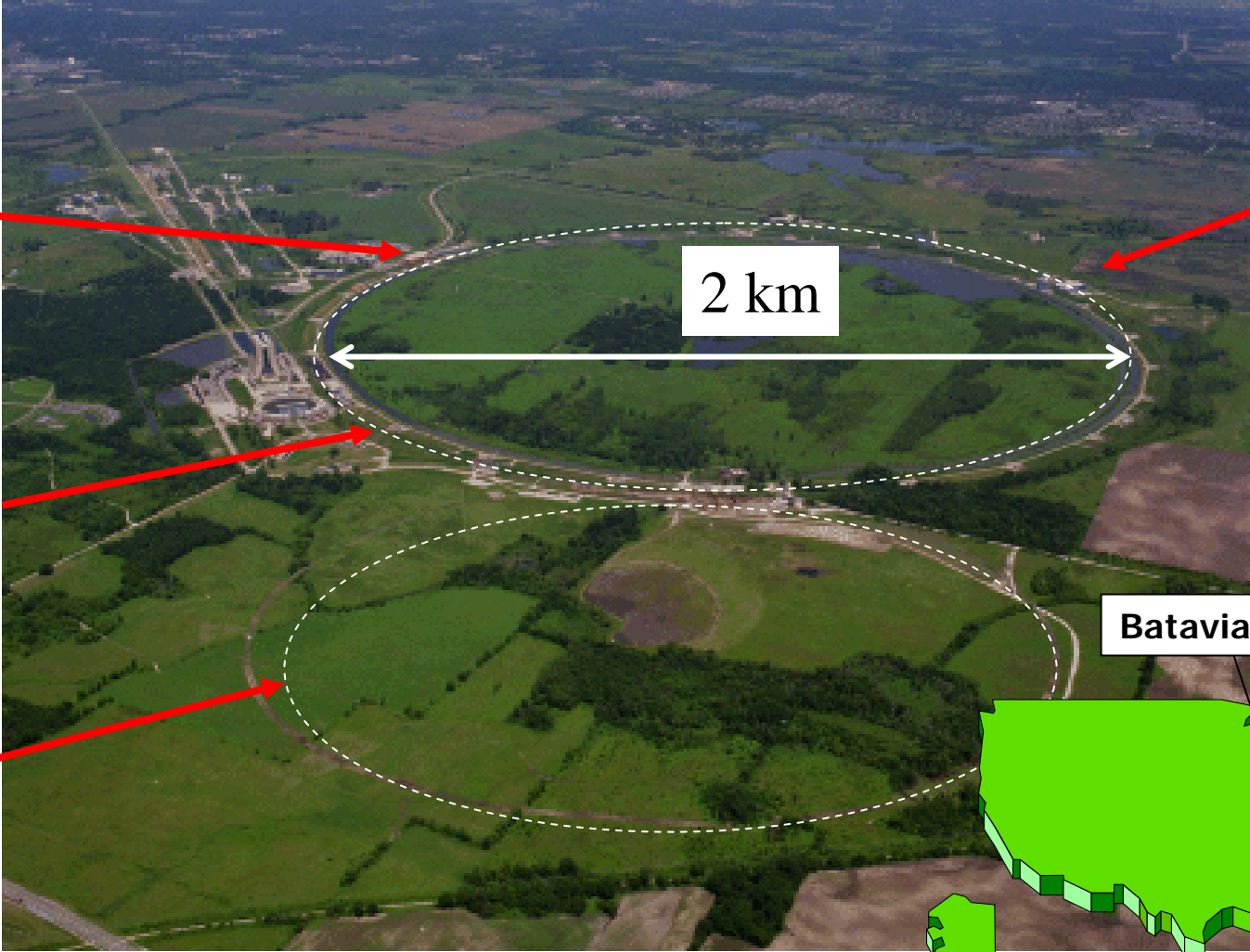
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New England Particle Physics Student Retreat
Falmouth, MA
23 Aug, 2004

Lecture #1

- Quick tour of the Tevatron
- Basics of hadron collisions
 - Partons and parton distribution functions
 - Kinematic variables
 - Cross sections and rates of important processes
- Motivating example: top quark physics
- Generic measurement and identification of objects
 - Jets
 - electrons, photons, hadrons, muons, neutrinos
- CDF and DØ detectors
 - electrons and photons
 - muons
 - missing transverse energy
 - b-tagging
 - triggering

Tevatron Collider: Fermilab



CDF

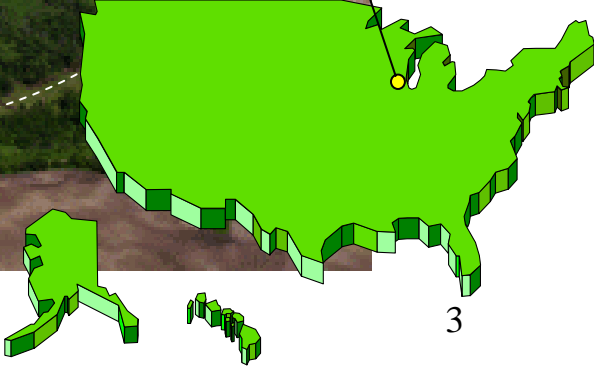
DØ

2 km

Tevatron

Batavia, Illinois

Main
Injector
& Recycler



Tevatron Collider

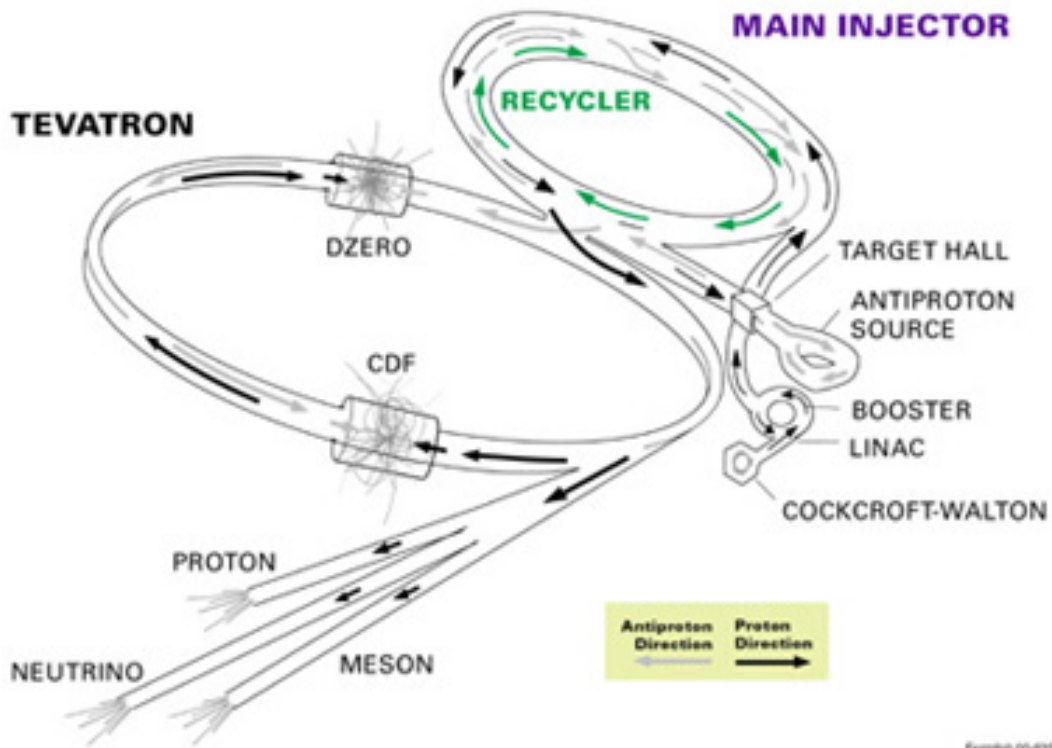
Beam:
900 GeV
protons

“Target”:
900 GeV
antiprotons



Antiprotons

FERMILAB'S ACCELERATOR CHAIN



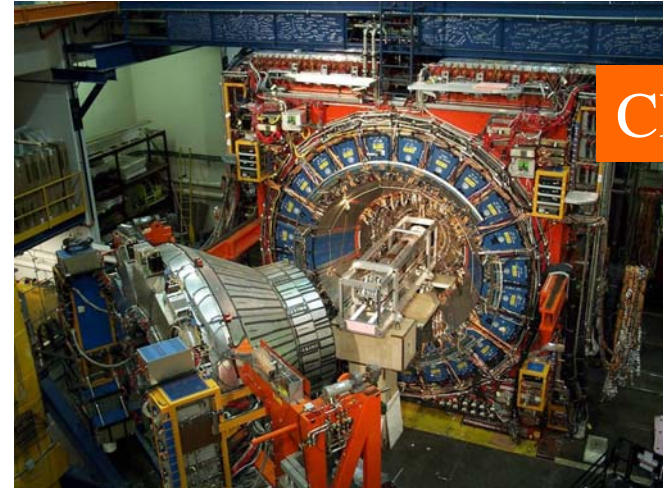
Antiprotons are created in collisions of the proton beam with a nickel target, then collected, cooled, and stored.



$\sim 2 \times 10^{12}$ antiprotons stored

The Detectors

- Much more discussion later
- Important features to keep in mind for now
 - Collisions take place at the (approximate) center of the detectors
 - Detectors must have apertures in the forward and backward directions for the beams to enter and exit
 - Detectors try to measure momentum and or energy of particles produced in collision



CDF



DØ

Tevatron in #'s

Tevatron parameters (2003):

Proton Anti-proton 36x36 bunches

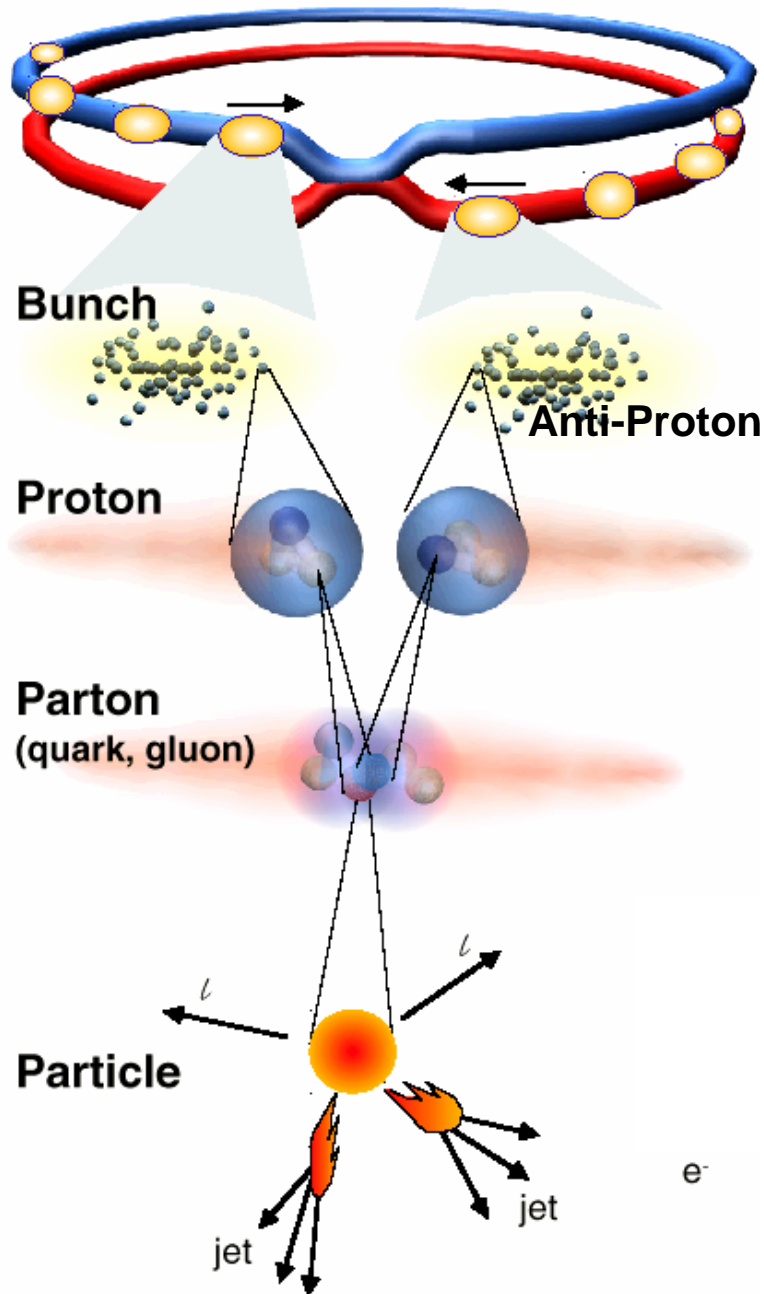
Protons/bunch 10^{11}
Anti-protons/bunch 10^{10}

Beam Energy 0.98 TeV/beam

Luminosity $\sim 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$

bunch crossing time 396 ns

Collision Rate 2.5 MHz

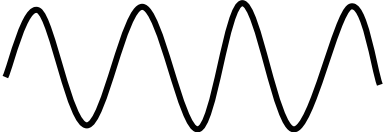


Brief History of the Tevatron Collider

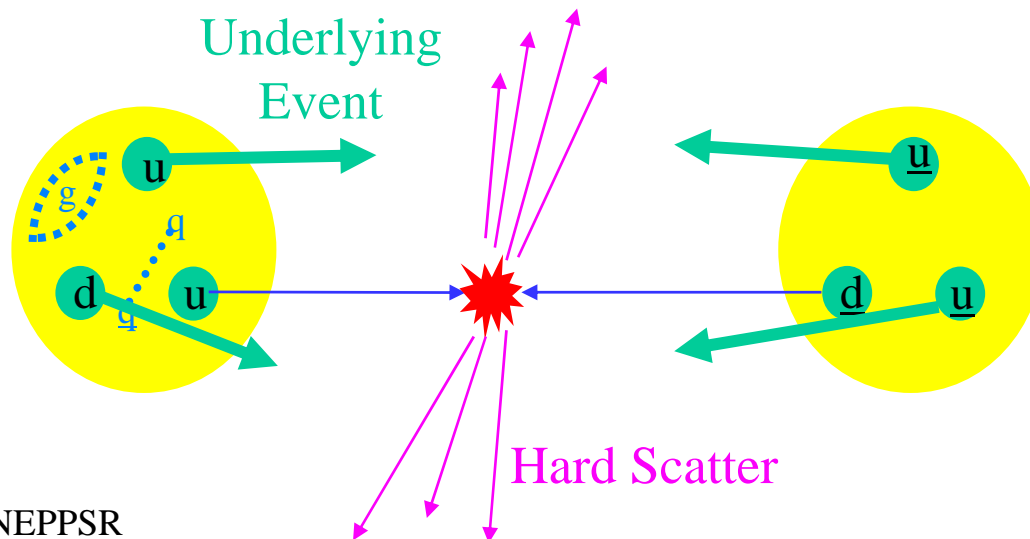
- 1983: Tevatron accelerator began operations
- 1985: First collisions of the Tevatron Collider
- 1986-1989: “Run 0” of the collider
 - Center-of-mass energy: 1.8 TeV
 - Integrated luminosity $\sim 4.5 \text{ pb}^{-1}$
 - CDF detector only
- 1992-1996: “Run I” of the collider
 - 1.8 TeV
 - Integrated luminosity $\sim 100 \text{ pb}^{-1}$
 - Both CDF and DØ detectors
 - Top quark discovery
- 2001-present: “Run II”
 - 1.96 TeV
 - Upgraded CDF and DØ detectors
 - Integrated luminosity $\sim 500 \text{ pb}^{-1}$ so far...
 - Anticipate $4\text{-}8 \text{ fb}^{-1}$ by the end of Run II (~ 2009)

Hadron collisions

(too) Simple minded calculation: deBroglie wavelength of proton


$$\lambda = \frac{hc}{pc} = \frac{1.2\text{GeV} \cdot \text{fm}}{980\text{GeV}} \approx 10^{-18} \text{m}$$

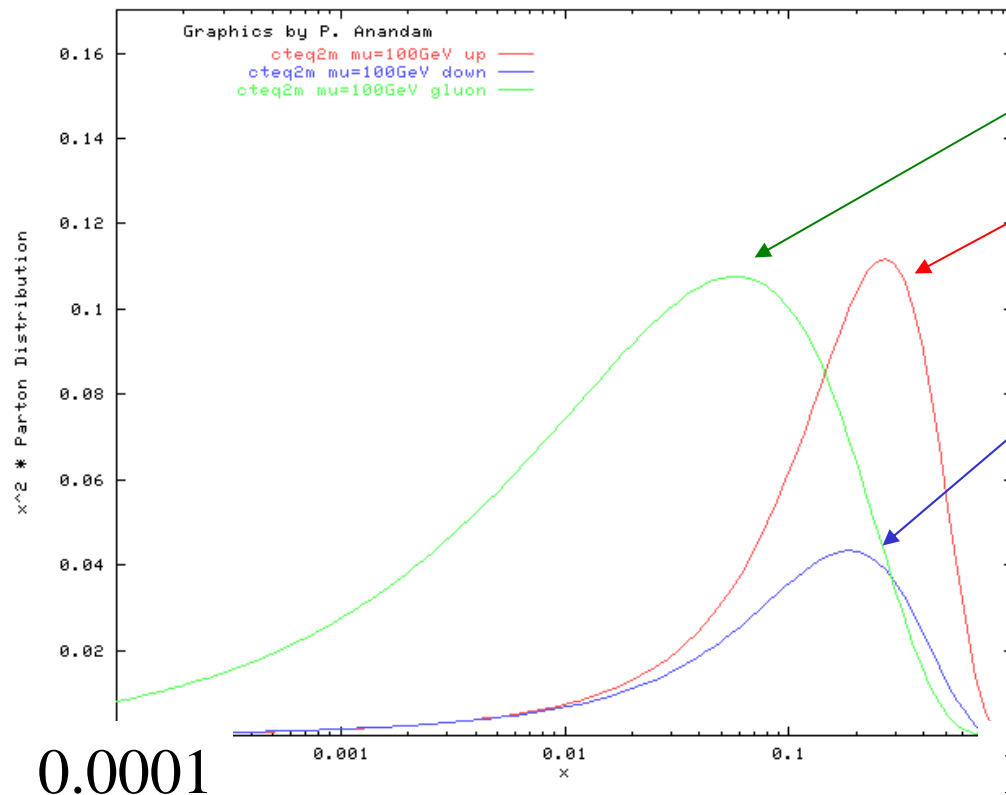
Much smaller than the size of a (anti)proton ($\sim 10^{-15}\text{m}$)
 \Rightarrow hard scatter involves only one parton (q,g) from each



Parton distribution functions

- Probability of finding a parton (quark or gluon) with a fraction x of the (anti)proton's momentum is given by the pdf, $f(x, Q^2)$ for a momentum scale Q

$x^2 * \text{pdf}$

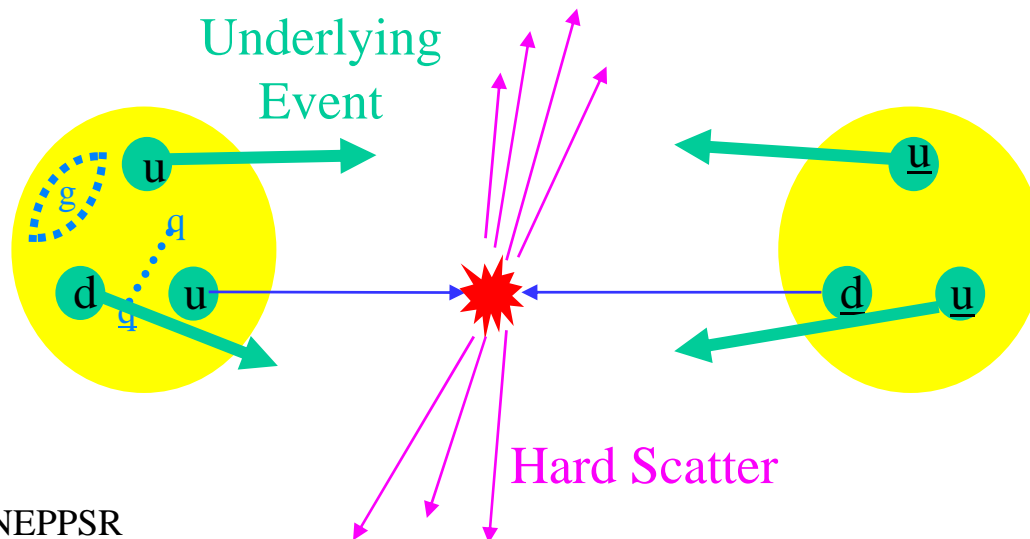


(on a log scale area under $x^2 * \text{pdf}$ shows momentum fraction)

Log scale

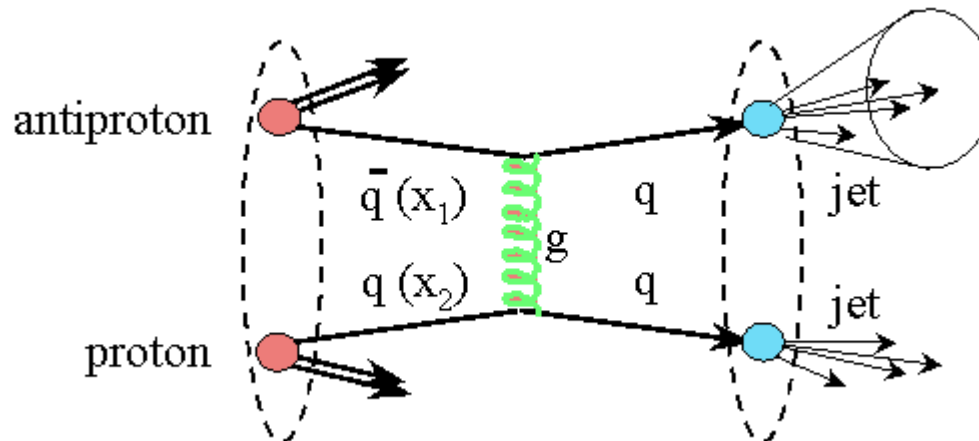
Hadron collisions: experimental consequences

- Energy involved in “hard scatter” is less (typically <10%) than the full proton-antiproton center-of-mass energy
- “hard scatter” system is generally
 - Not at rest along the beam direction $\sum p_z \neq 0$
 - Nearly at rest transverse to the beam direction $\sum \vec{p}_T \approx 0$
- Additional particles & energy are present from the “underlying event”



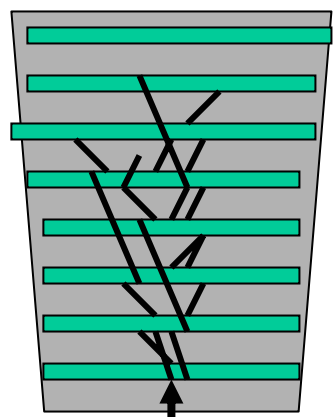
Most common high-pt process: jet production

Quarks are not free, so what emerges is a collimated jet of hadrons along the original quark or gluon direction



As seen by the calorimeter:

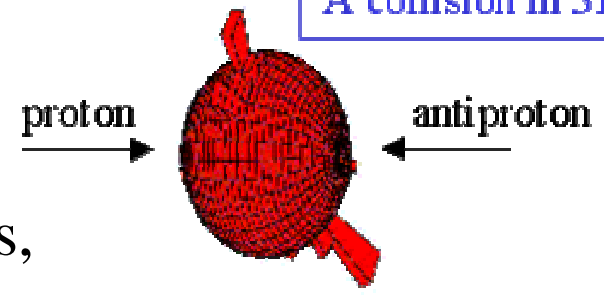
A single calorimeter cell



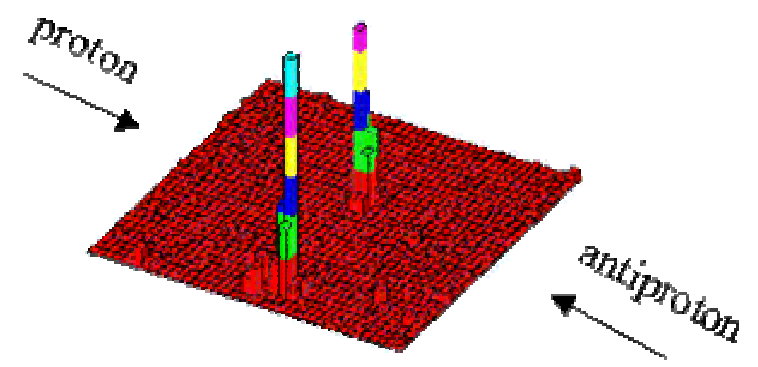
DØ: Uranium plates, with liquid Argon in the gaps

CDF: Iron plates with plastic scintillator in the gaps

A collision in 3D



The same collision in 2D



An event observed in the detector:

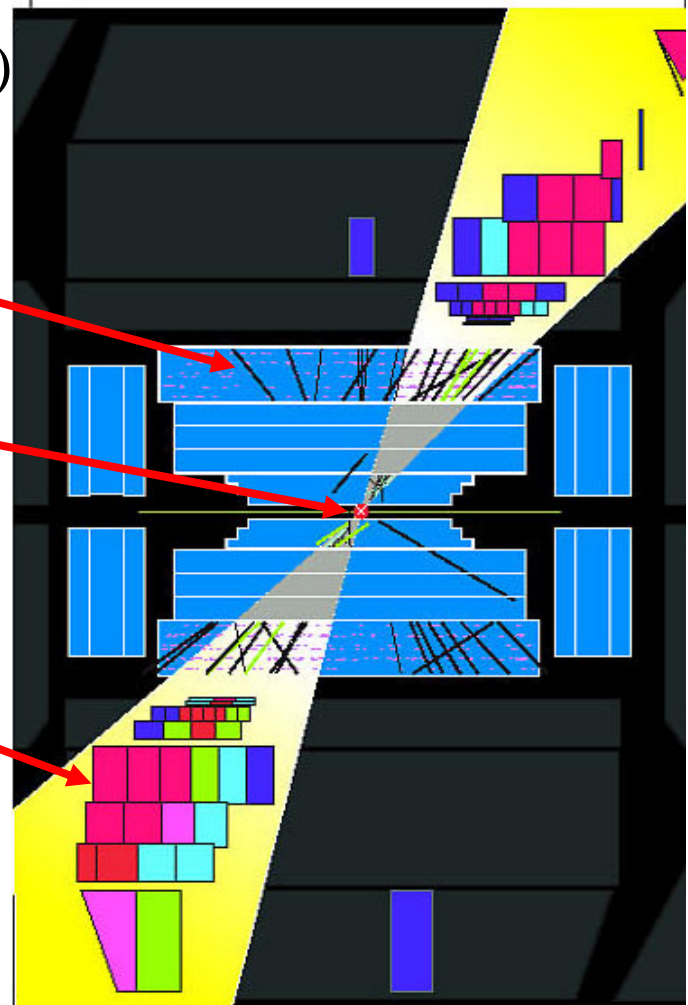
(2-dimensional slice)

Charged tracks

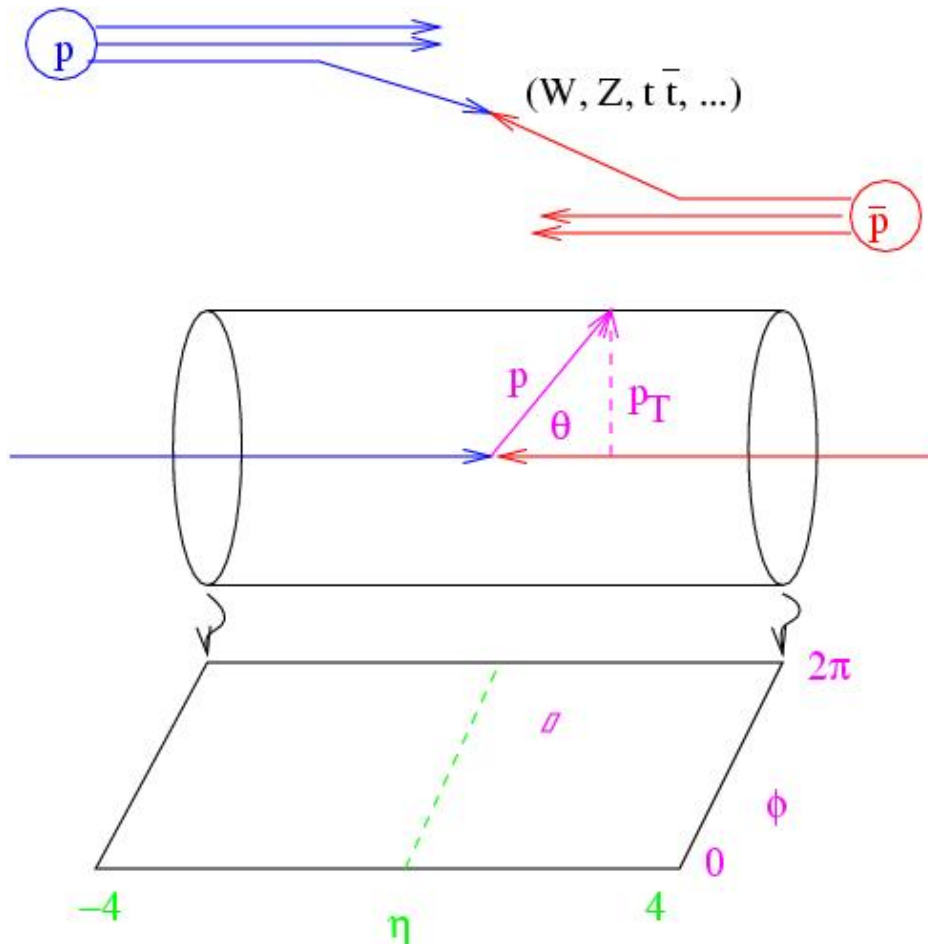
Point of collision

Colors correspond to energy deposited in a “cell” of the calorimeter

Note that energy is concentrated in two narrow cones, or jets.
“Simple cone algorithm”: sum all energy in a cone around jet axis



Coordinate system



Beam axis : z

Azimuth: ϕ

Rapidity:

$$y = \frac{1}{2} \ln \left(\frac{E + p_z}{E - p_z} \right)$$

pseudorapidity:

$$\eta = -\ln(\tan(\theta/2)) \approx y$$

Transverse momentum:

$$p_T = p \sin \theta = \sqrt{p_x^2 + p_y^2}$$

Phase space variables: eta, phi, and p_T

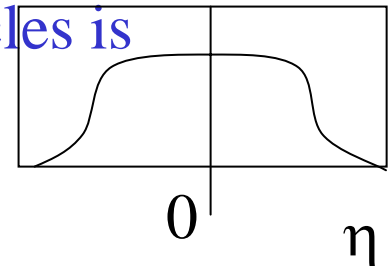
- Recall that the hard scatter system is generally in motion in the z-direction with respect to the laboratory frame
- Under a boost in the z-direction:
 - p_T is invariant
 - ϕ is invariant
 - Rapidity itself is not invariant, but all differences in rapidity are invariant

$$y \rightarrow y - \tanh^{-1} \beta$$

$$y_1 - y_2 \rightarrow y_1 - y_2$$

- At a given p_T , the expected density of particles is (approximately) uniform in eta and phi.

- In contrast to e^+e^- , where $\cos\theta$ is flat

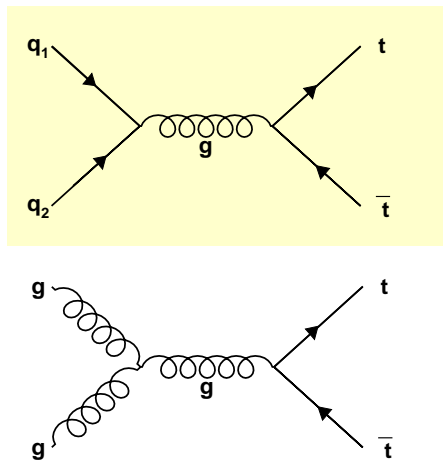


Typical production rates for p-pbar at 2 TeV

Final state	Cross section (pb)	Rate at $L=10^{32} \text{ cm}^{-2}\text{s}^{-1}$
“minimum bias”	4×10^{10}	4 MHz
2 jets	4×10^6	400 Hz
4 jets	1.6×10^5	16 Hz
6 jets	6000	0.6 Hz
W	30000	3 Hz
Z	9000	0.9 Hz
WZ	3.5	3.5×10^{-4} Hz (1.3/hour)
t tbar	7.5	7.5×10^{-4} Hz (3/hour)

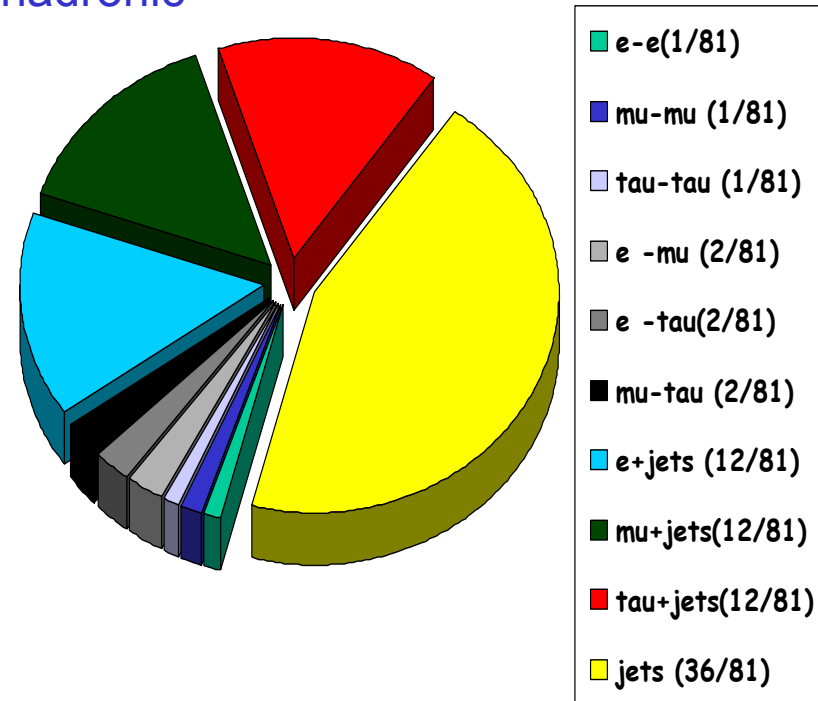
Top quark production and decay

- in proton anti-proton collisions at Tevatron energies, top quarks are primarily **produced in pairs** (Strong interactions)



Single top production
(Electroweak interactions):
not yet observed

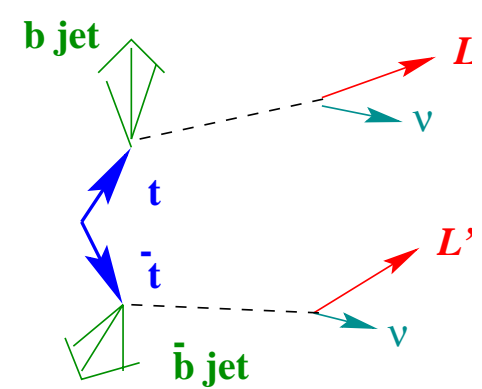
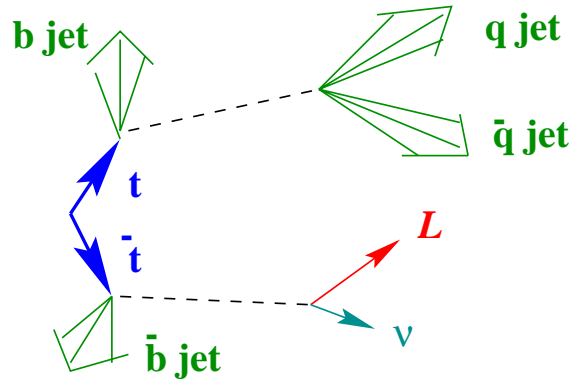
- lifetime very short ($\approx 10^{-24}$ s), $\text{Br}(t \rightarrow Wb) = 100\%$
- Both W 's decay via $W \rightarrow l\nu$ ($l = e$ or μ ; 5%)
dilepton
- One W decays via $W \rightarrow l\nu$ ($l = e$ or μ ; 30%)
lepton+jets
- Both W 's decay via $W \rightarrow qq$ (44%)
all hadronic



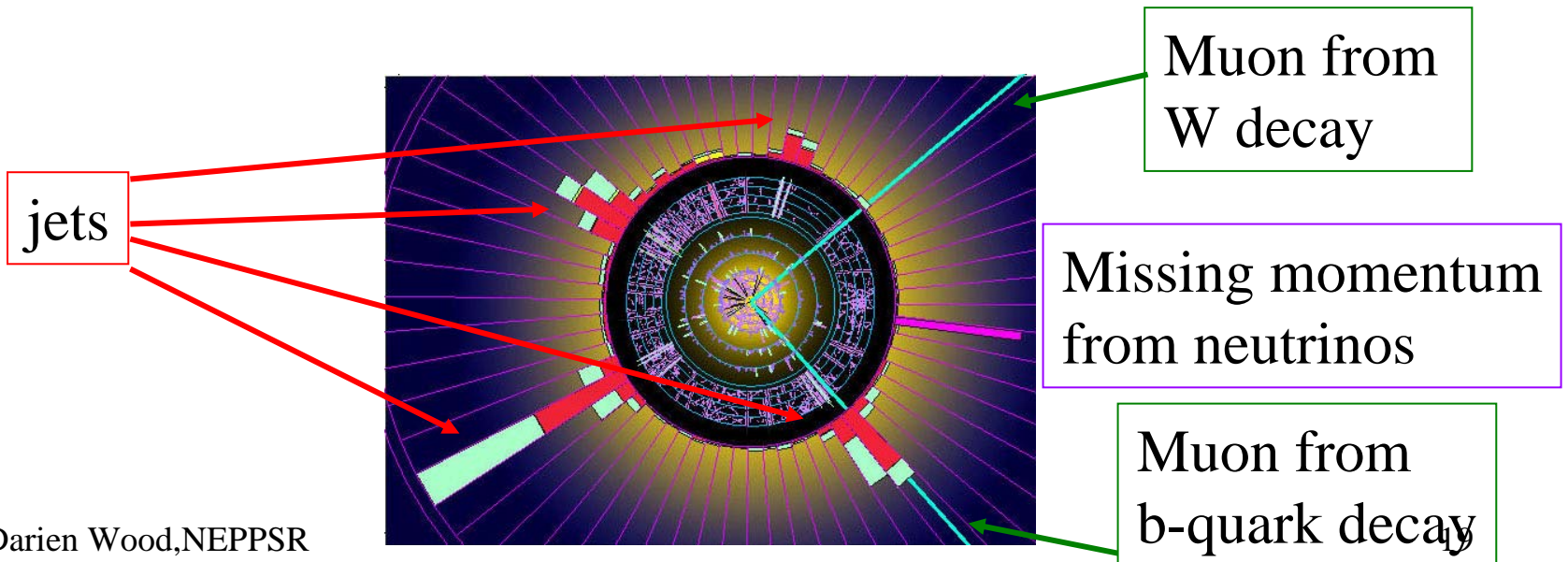
Top quark

lepton+jets

dilepton



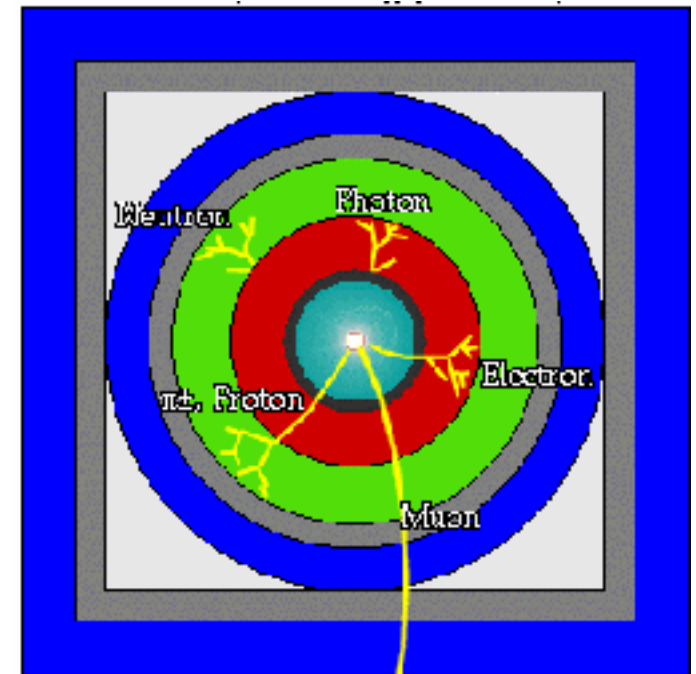
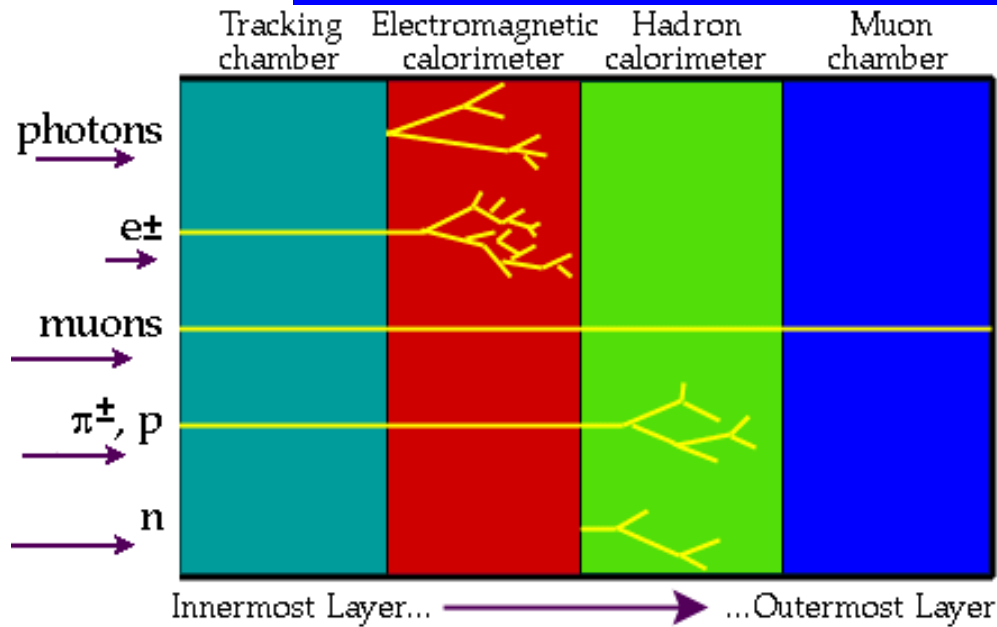
Uses all parts of the detector



Hadron collider detector principles

- Strong interaction dominates
 - Typically ~20 charged tracks per event, mostly pions
 - Not easy to distinguish among various light hadrons (pions, kaons, ...) but often you don't care
 - Isolated leptons are a signature of something interesting and or unusual
 - EM or weak interaction
 - Decay of heavy object
- Basic distinguishable objects: jets, photons, electrons, muons, (taus), neutrinos
- Often useful to be able to separate jets from heavy quarks (b-jets, c-jets) from those from light quarks and gluons

A Generic Detector System



- Tracking chambers ⇒ trajectory of charged particles
- Calorimeters ⇒ measure energy
 - Electromagnetic: e , photon
 - Hadronic: pion, K, proton, neutrons...
- Muon Chambers ⇒ measure muon trajectory
- Magnets ⇒ charged particles bend in magnetic fields. Bend depends on charge and momentum

what's missing?

Calorimeter Wish list

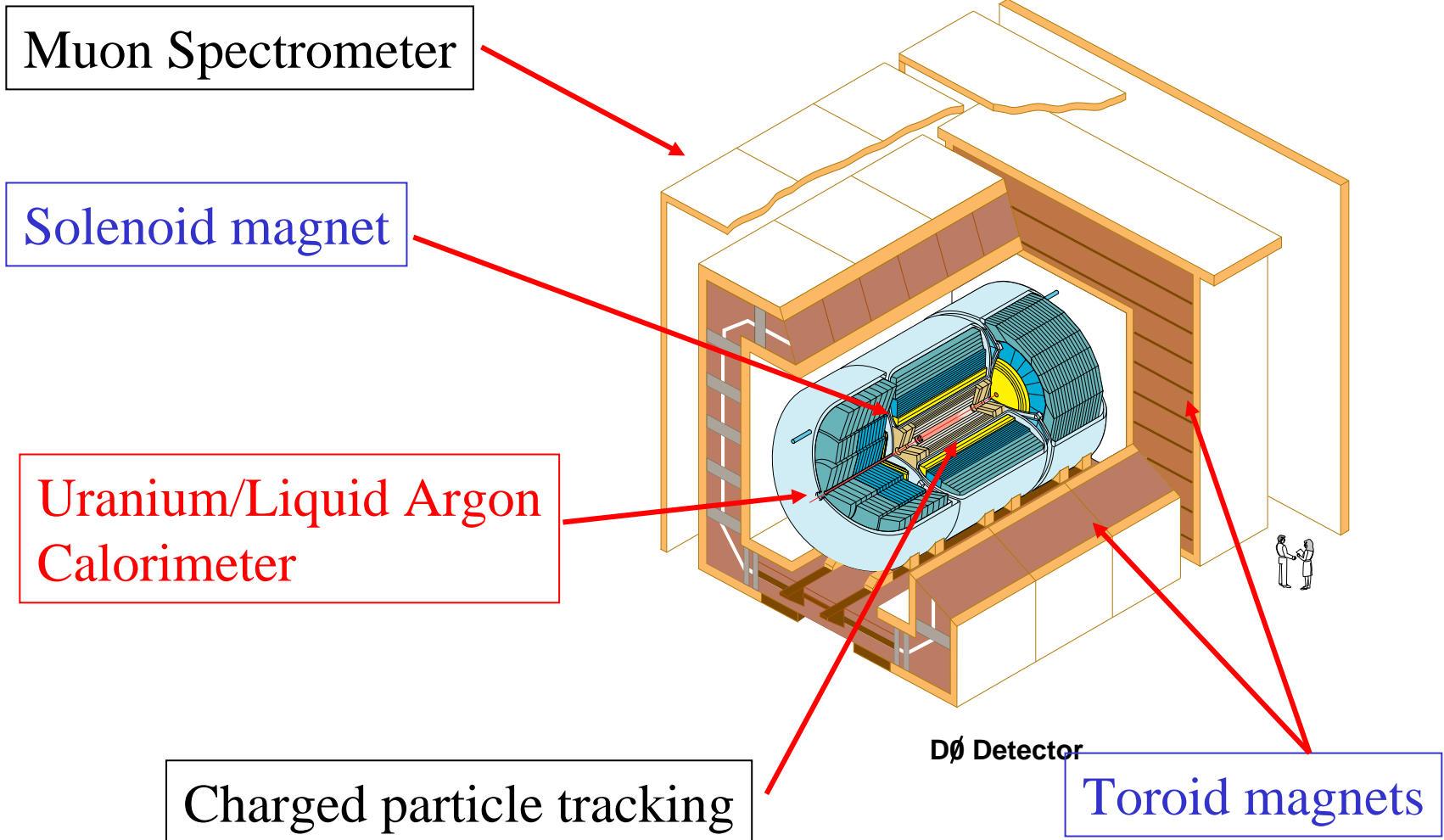
An ideal calorimeter should:

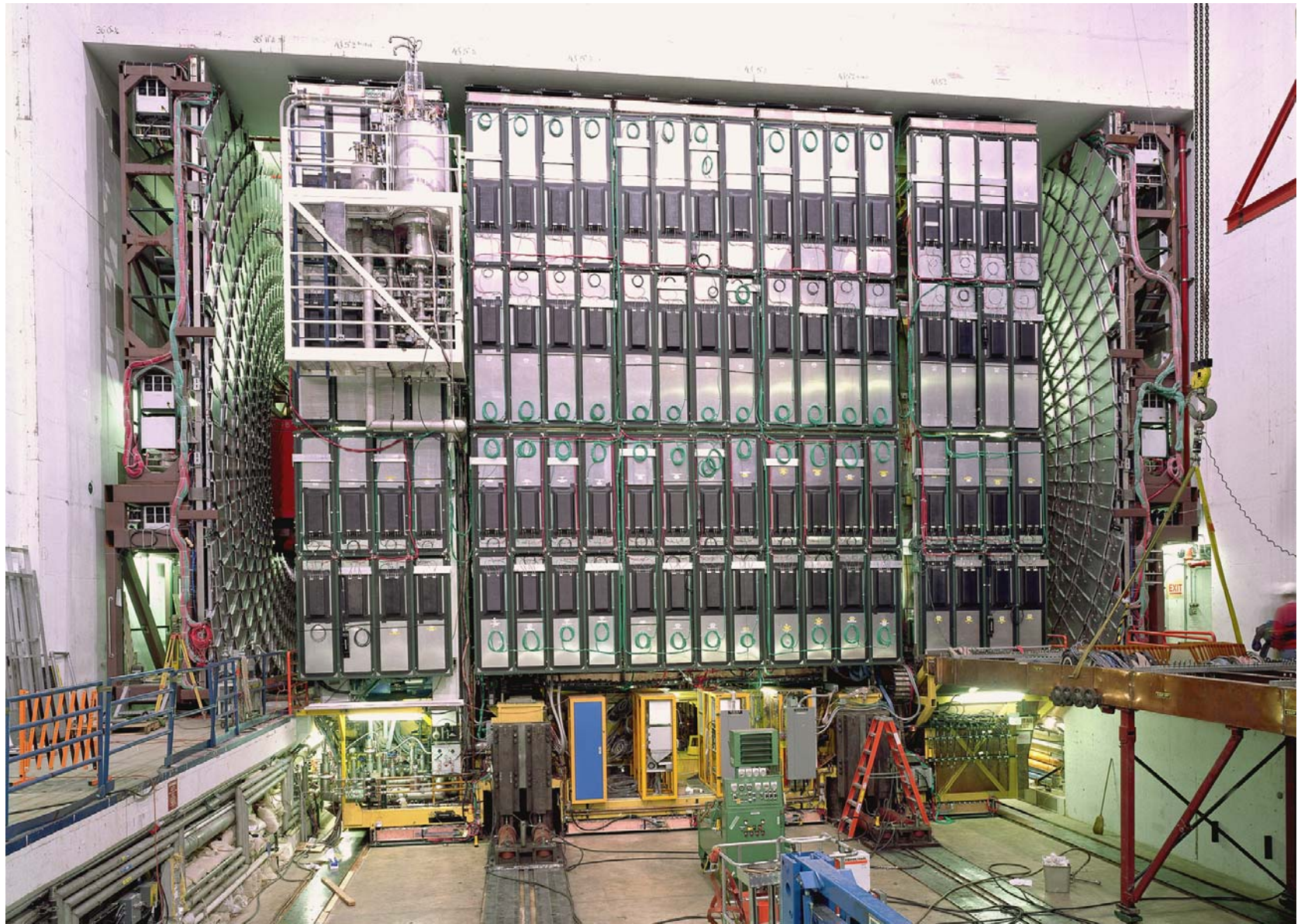
- measure the **energy** of each particle coming into it (except muons and neutrinos) with good **resolution**
- measure the **position** of the energy deposition, so that it can be associated with a momentum vector
- be able to **distinguish** different types of particle from each other by the way they “shower” inside the calorimeter
 - electrons & photons
 - hadrons (protons, pions, kaons, etc.) or jets of hadrons
 - muons
 - neutrinos

Wish list continued

- completely **contain** all particles (except muons and neutrinos) [“hermeticity”]
- have a constant, stable energy **calibration**
- have the **same** energy response for hadrons and electrons
- have a fine **segmentation** to be able to distinguish nearby particles (also important in being able to examine the transverse size of a shower)

DØ detector





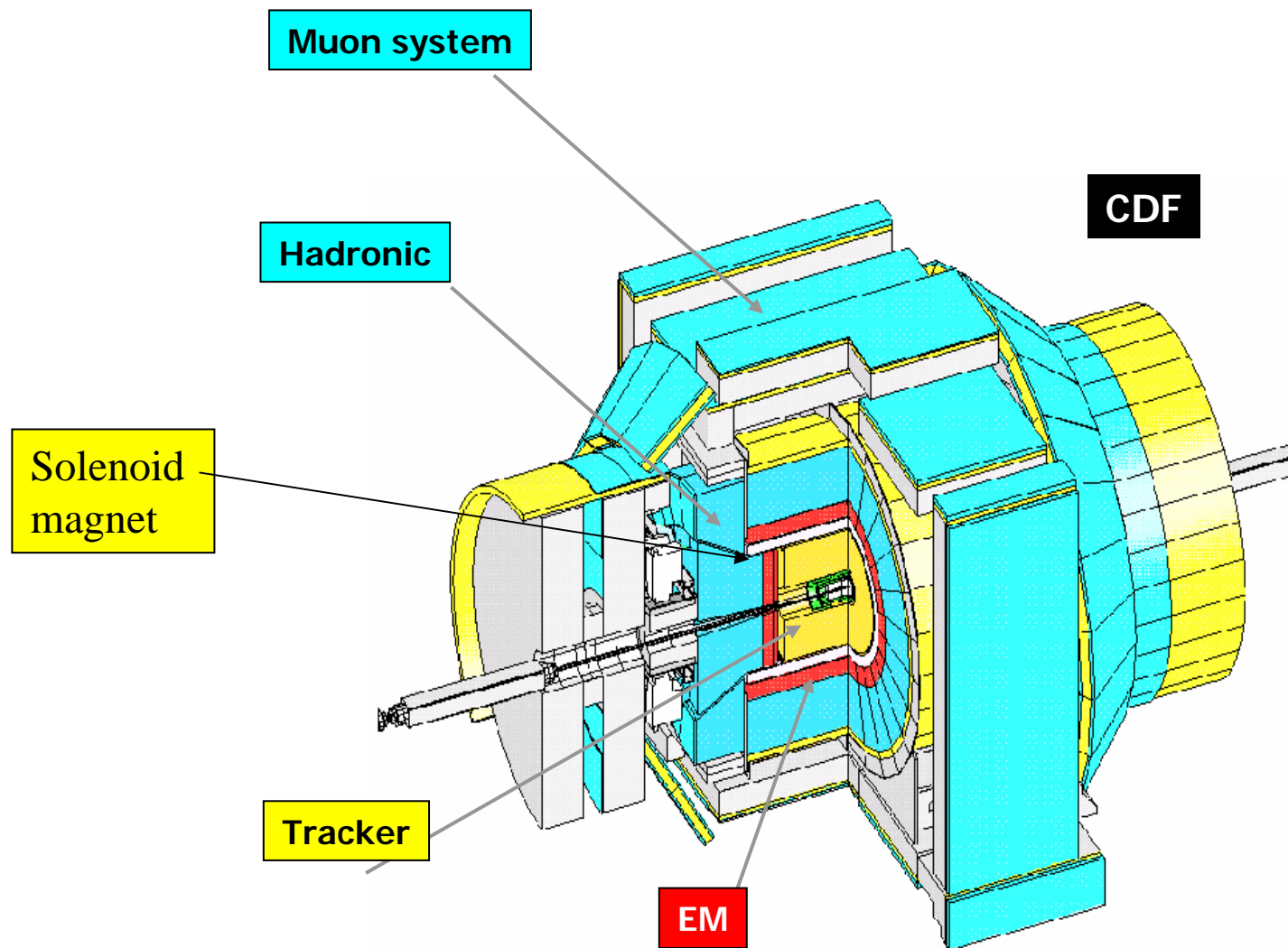
DØ detector, before the closing of the collision hall

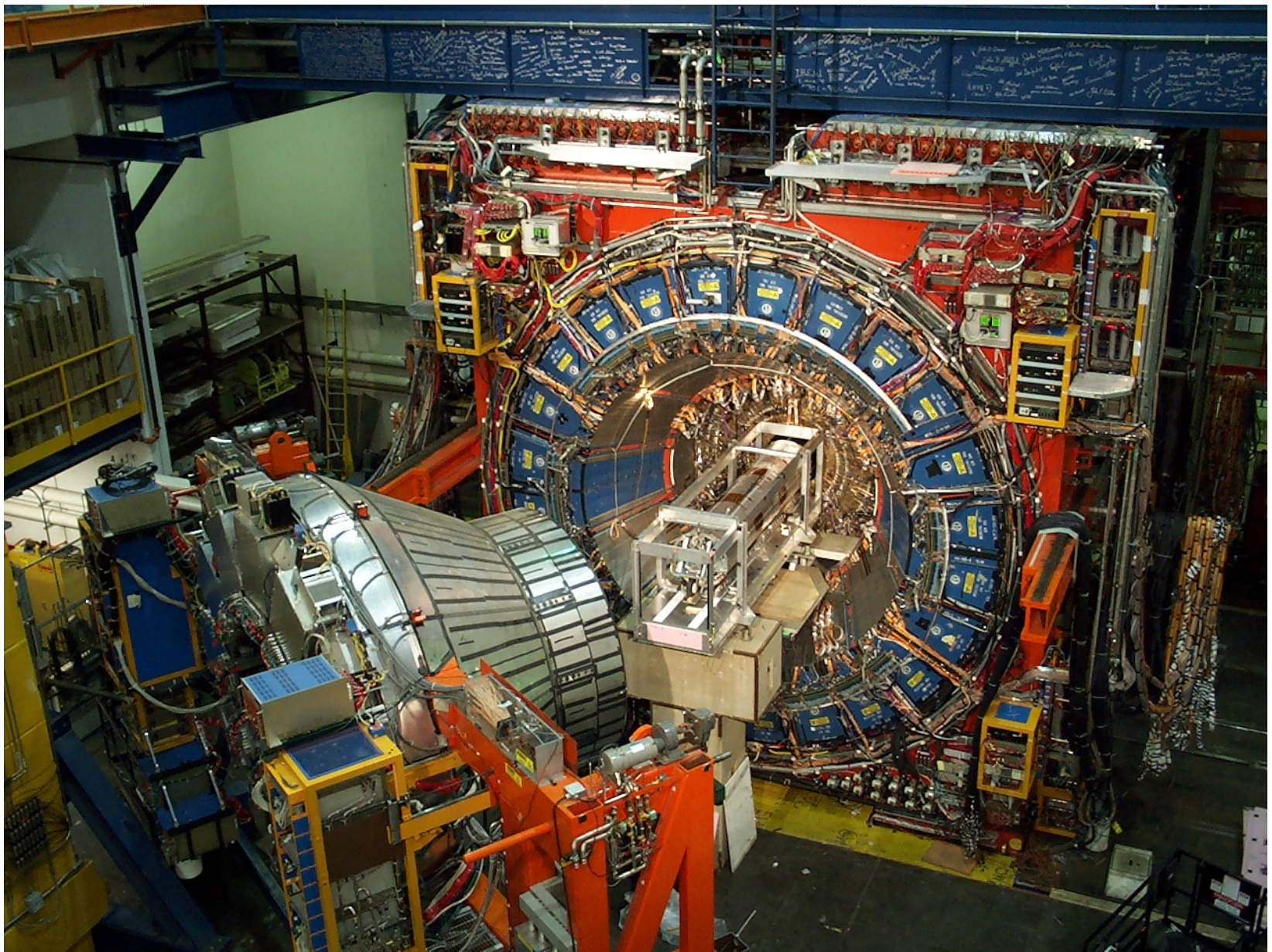
Darien Wood, NEPPSR

04

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CDF Detector



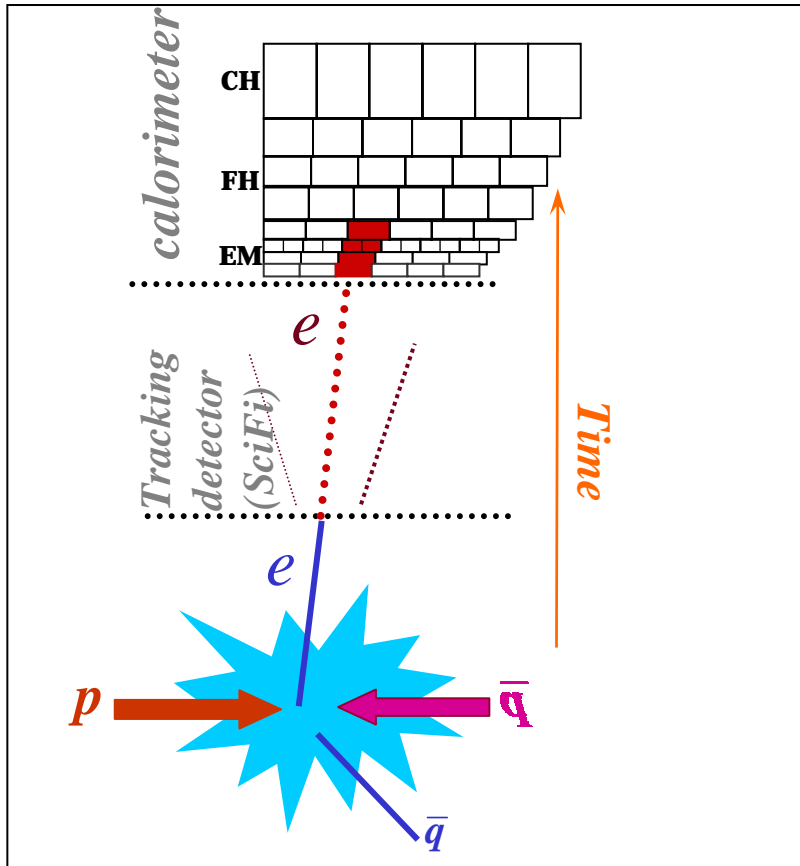


CDF detector, with the plug calorimeter retracted

CDF and DØ

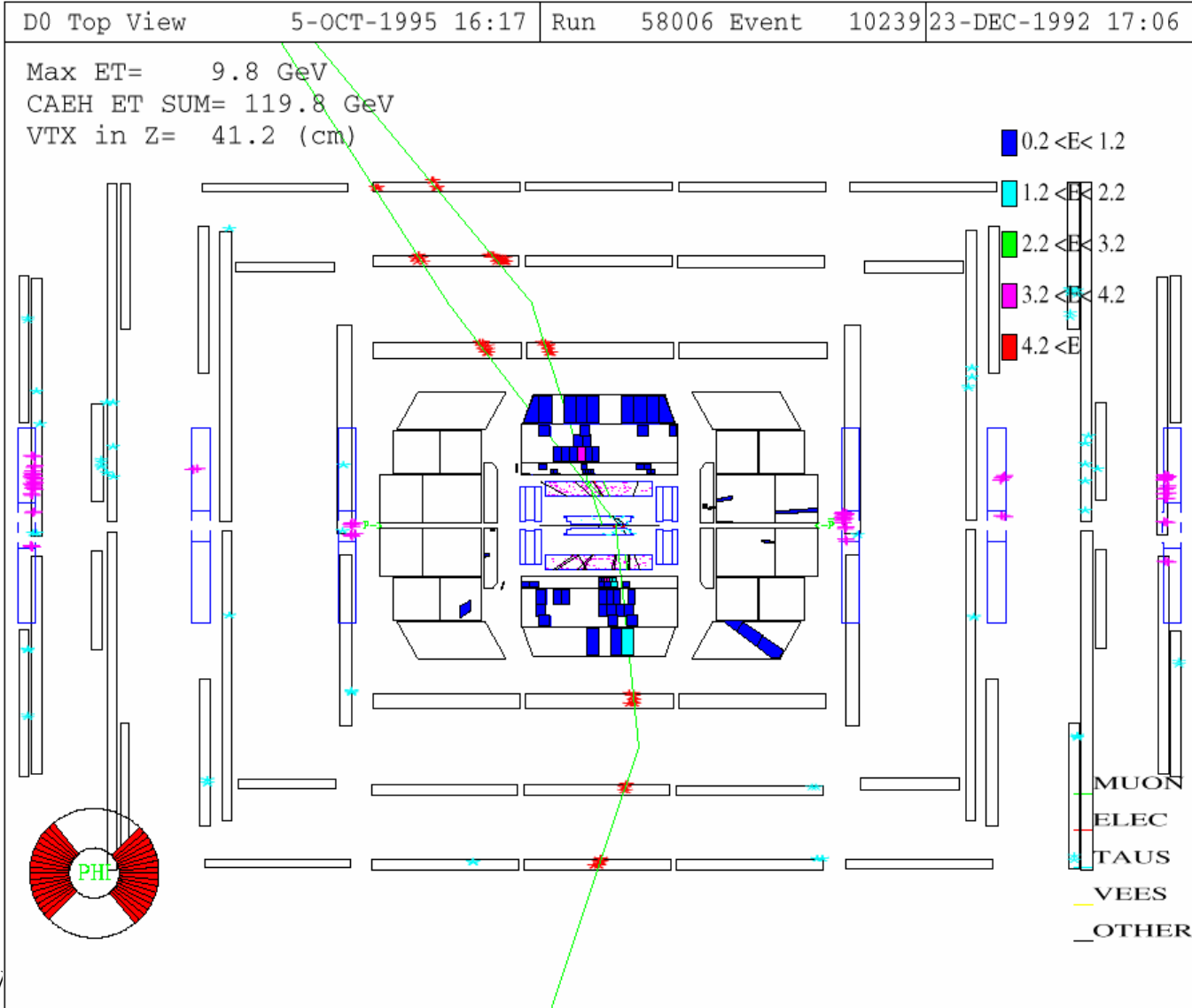
- Common to both detectors
 - Silicon microstrip detector (tracking, vertex)
 - Magnetic tracker
 - Central solenoid magnet
 - Preshower detectors
 - Electromagnetic calorimeter
 - Hadronic calorimeter
 - Iron absorber
 - Muon detectors
- CDF highlights
 - Large radius (1.5m) tracker: good momentum resolution
 - Time-of-flight system
- DØ highlights
 - Hermetic, dense calorimeter
 - More complete muon coverage with magnetized iron toroids

Electrons and photons

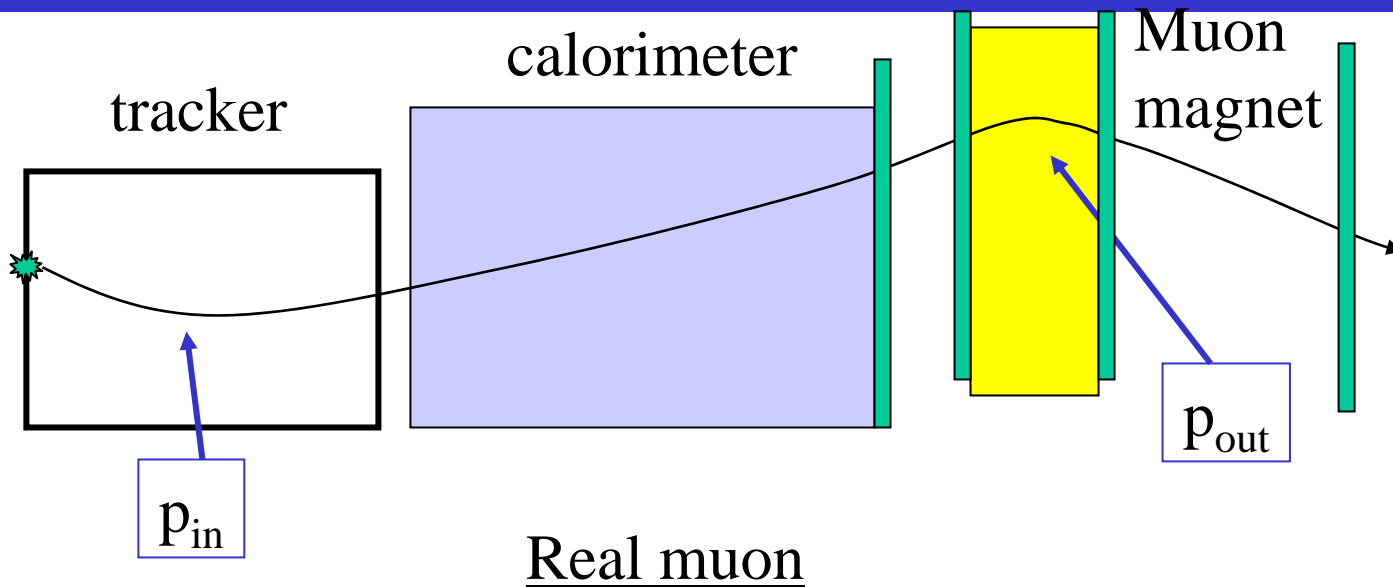


- In the calorimeter
 - Narrow cluster of energy in EM (front) section
 - Little or no energy in Had (back) section
- In the tracker
 - **Photon**: no track (no hits)
 - **Electron**: track with momentum matching energy of calorimeter cluster
- Fake backgrounds
 - Jets with most of their energy in π^0 (‘s)
 - **Photon**: electron with missing track
 - **Electron**: photon conversion

Identifying Muons



Muon Signal

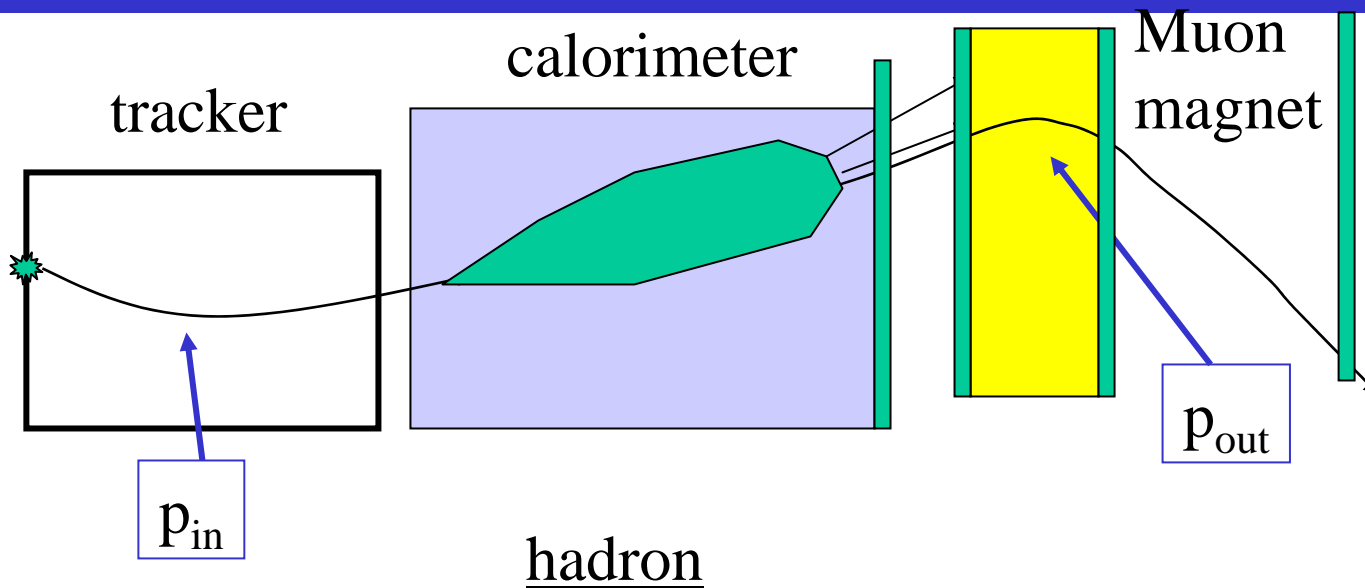


$$p_{in} \approx p_{out} + E_{loss}$$

(muon ID tool)

Better resolution comes from tracker; p_{out} dominated by multiple scattering (or showering)

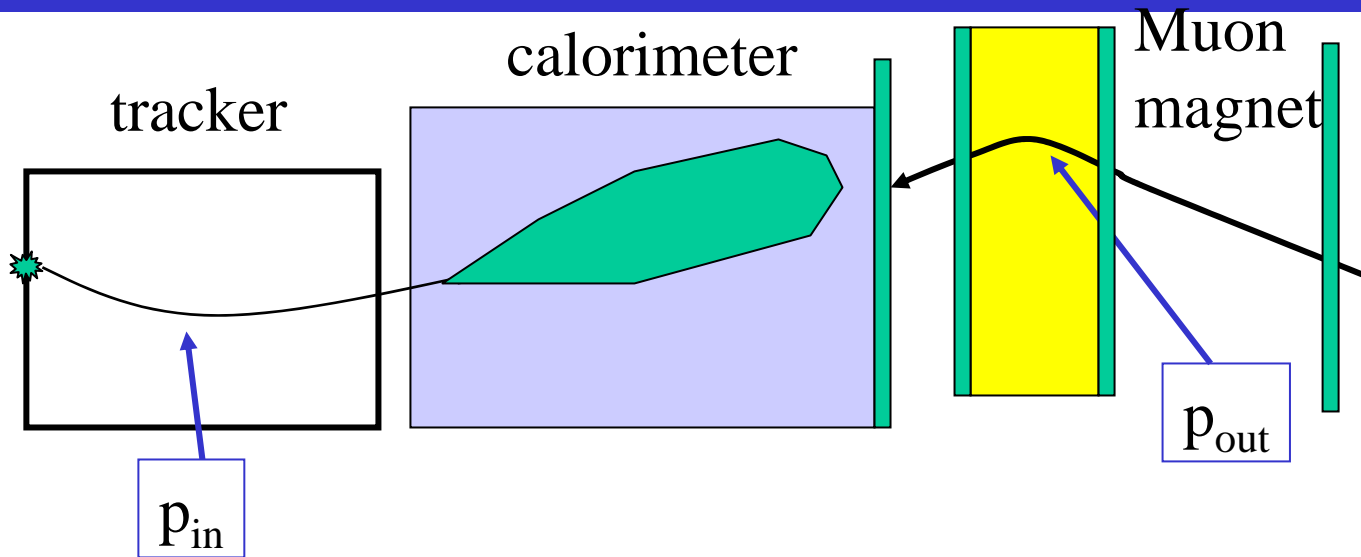
Muon background 1: punchthrough/decay



$$p_{in} \gg p_{out} + E_{loss}$$

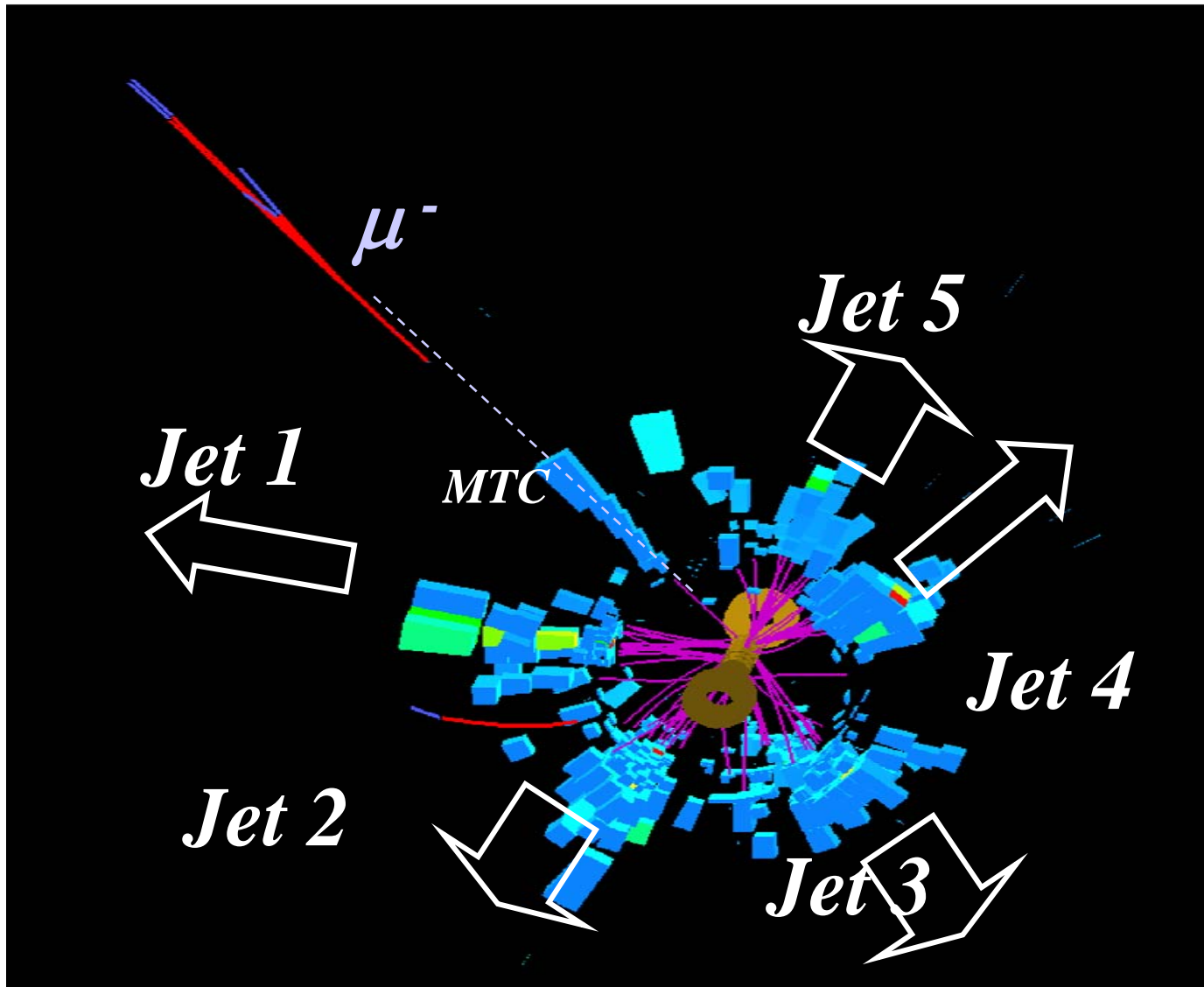
Outer decay/p.t. track points back to parent hadron,
but momenta do not match.

Muon background 2: halo/backscatter



$$p_{in} \neq p_{out} + E_{loss}$$

Good timing (scintillator) can get rid of most of these



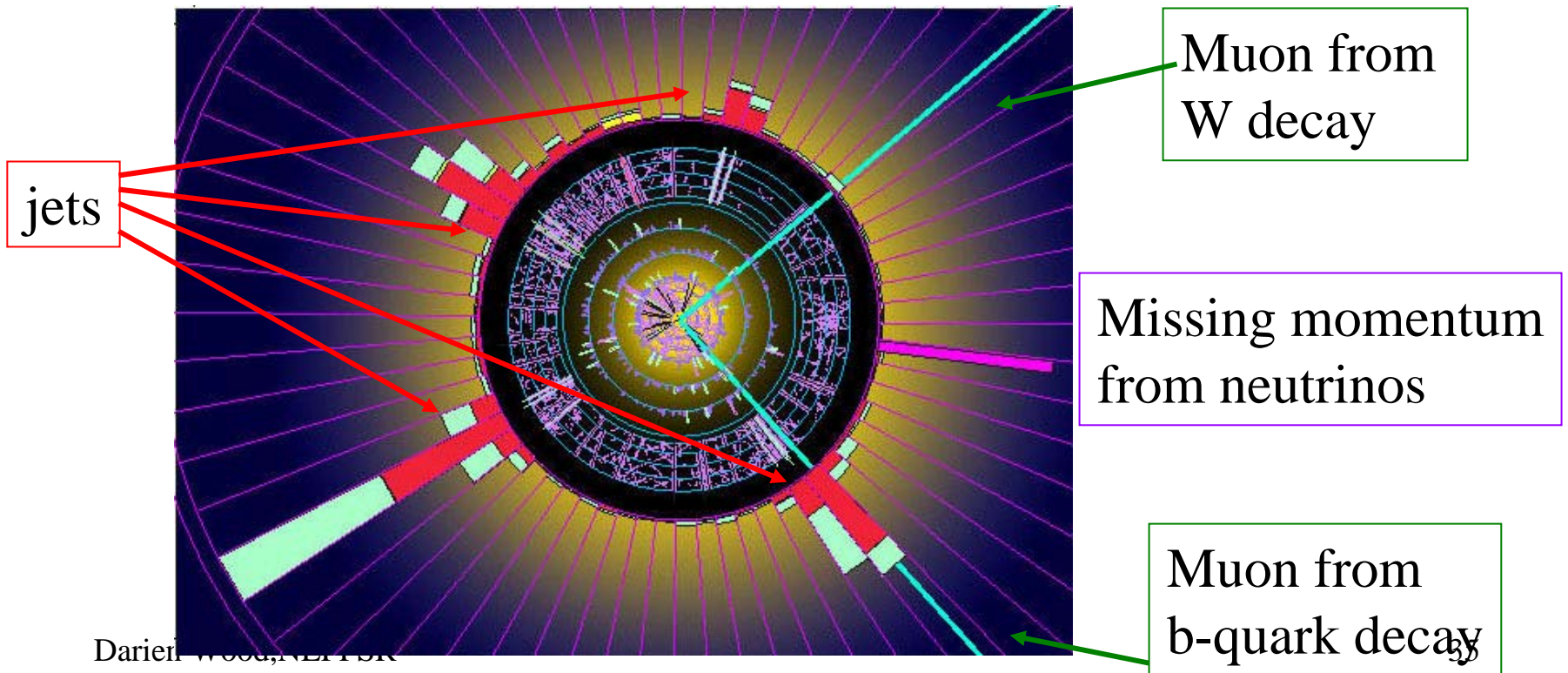
DØ top to μ +jets Candidate Event

Missing Transverse Energy (\cancel{E}_T)

Example: top quark candidate ($t\bar{t} \rightarrow \mu\nu jjjj$)

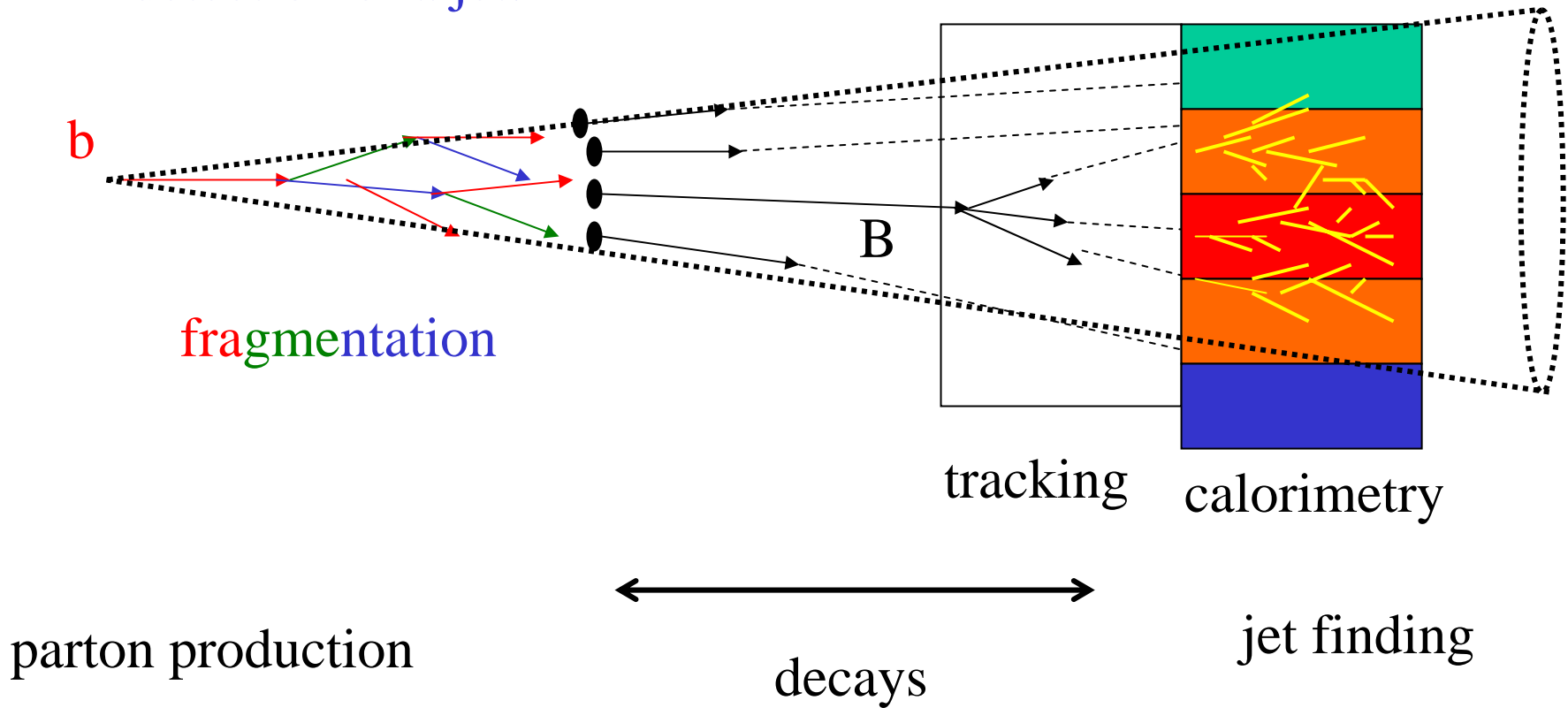
[transverse view]

$$\vec{\cancel{E}}_T = -\sum \vec{E}_T(\text{cal}) - \sum p_T(\text{muons})$$



b-quark jets

- Recall the steps between production of a quark and detection of a jet:

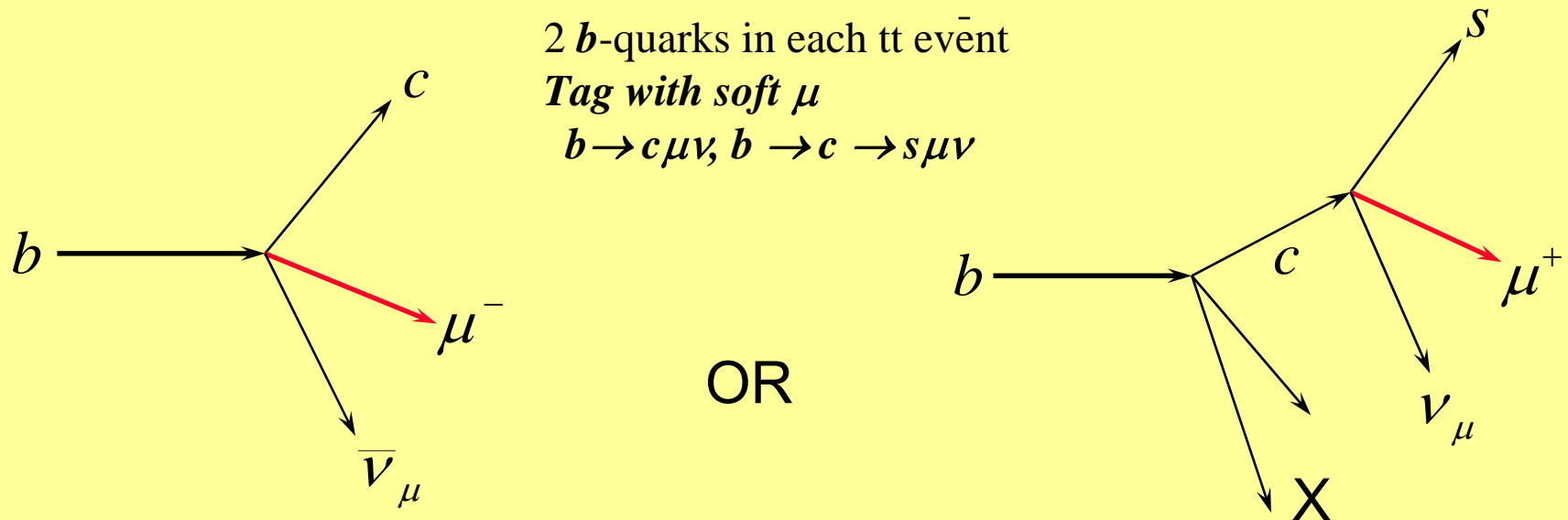


Identifying the b quark

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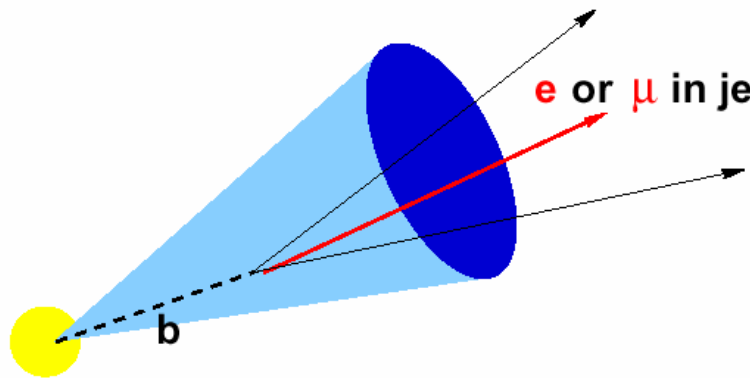
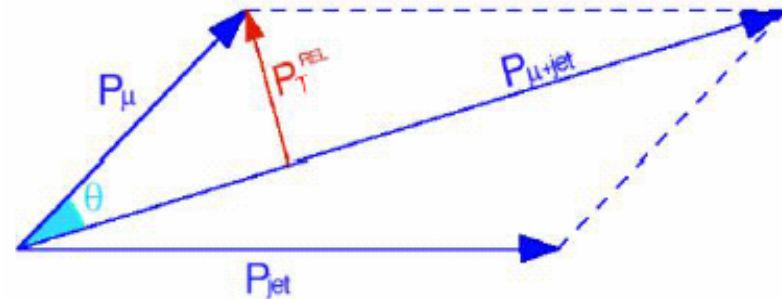
Semileptonic decays of the b-quark

example: $B(b \rightarrow \mu + X) \approx 20\% \Rightarrow$ detect muons in jets



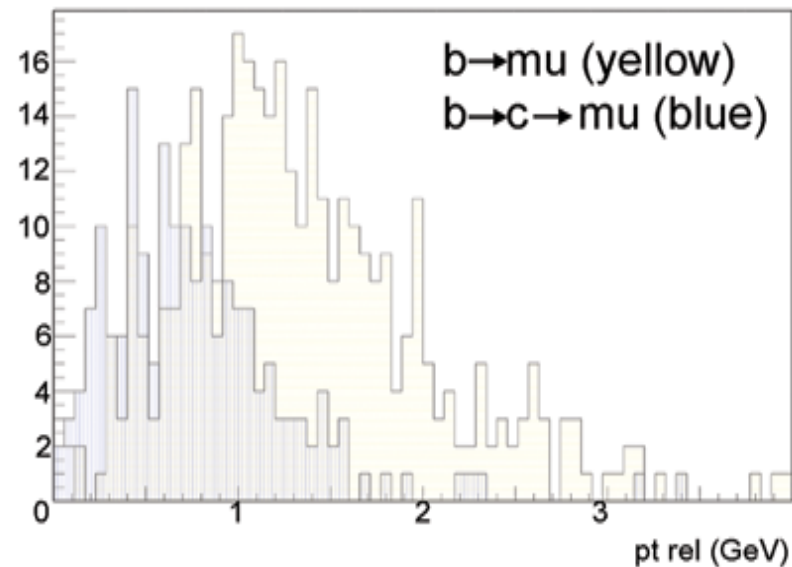
Soft lepton tag

- These leptons have a softer p_T spectrum than leptons from W/Z
- are not isolated

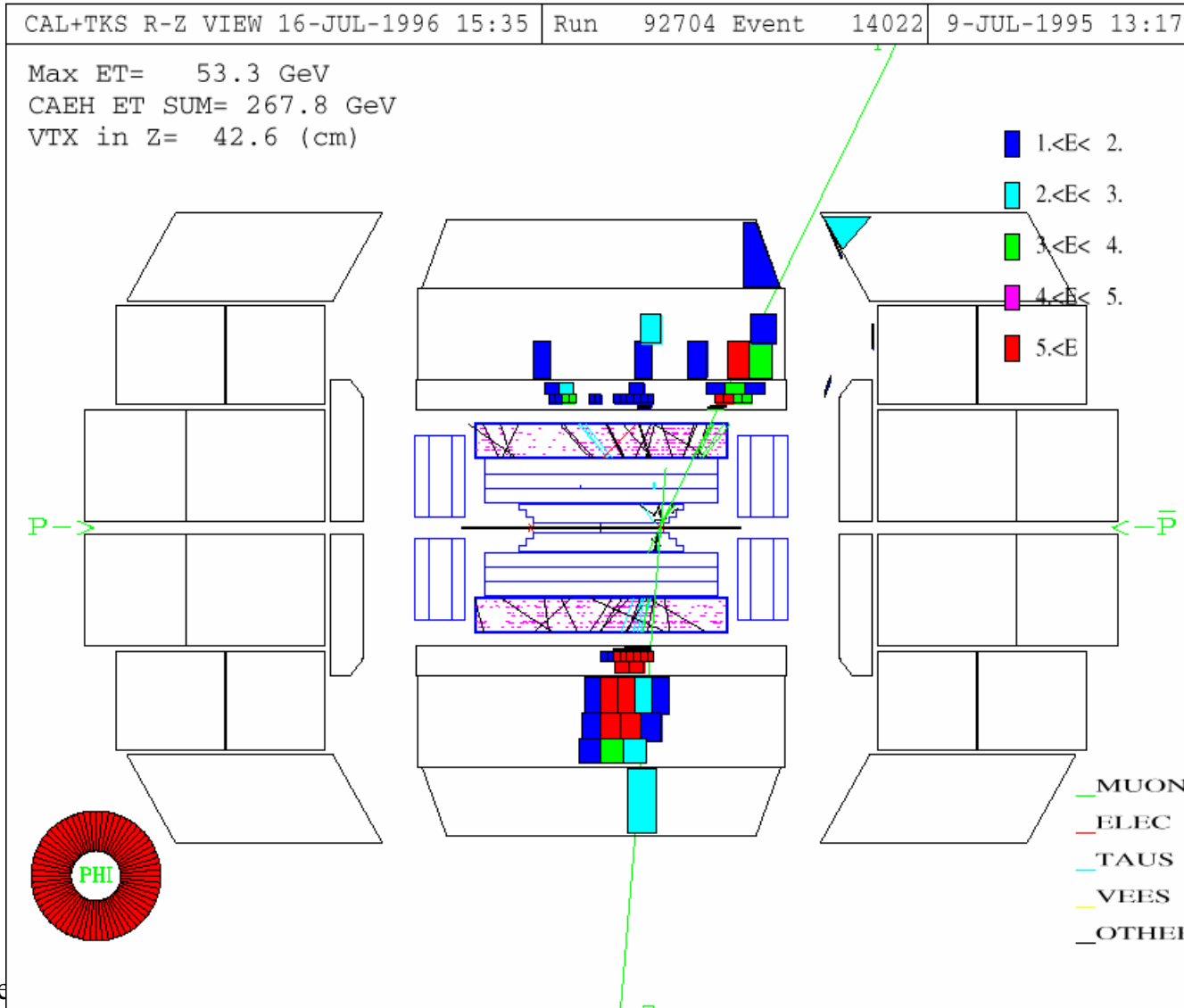


- $b \rightarrow l\nu c$ (BR $\sim 20\%$)
- $b \rightarrow c \rightarrow l\nu s$ (BR $\sim 20\%$)

Comparison of PtRel

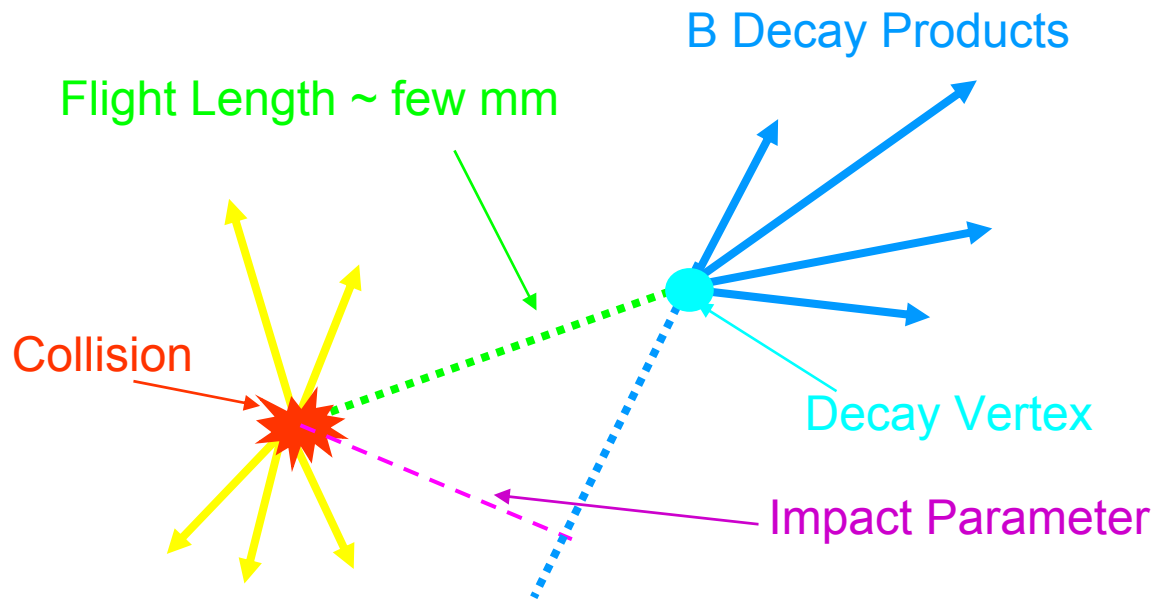


Top Event with b-jet (D0)



Identifying the b quark

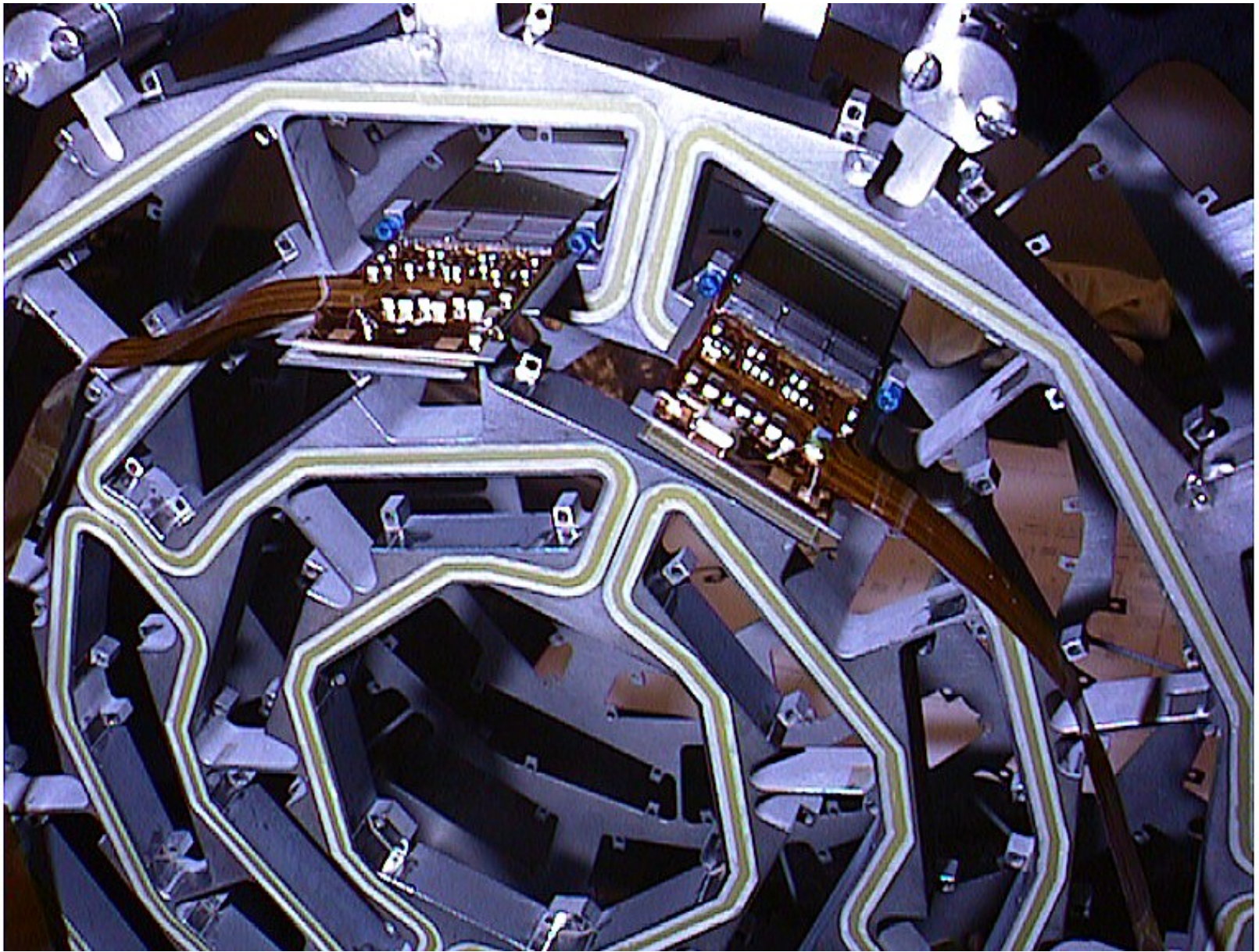
- 2 life time ≈ 1.5 ps $\Rightarrow c\tau \approx 0.5$ mm (short, but not too short)

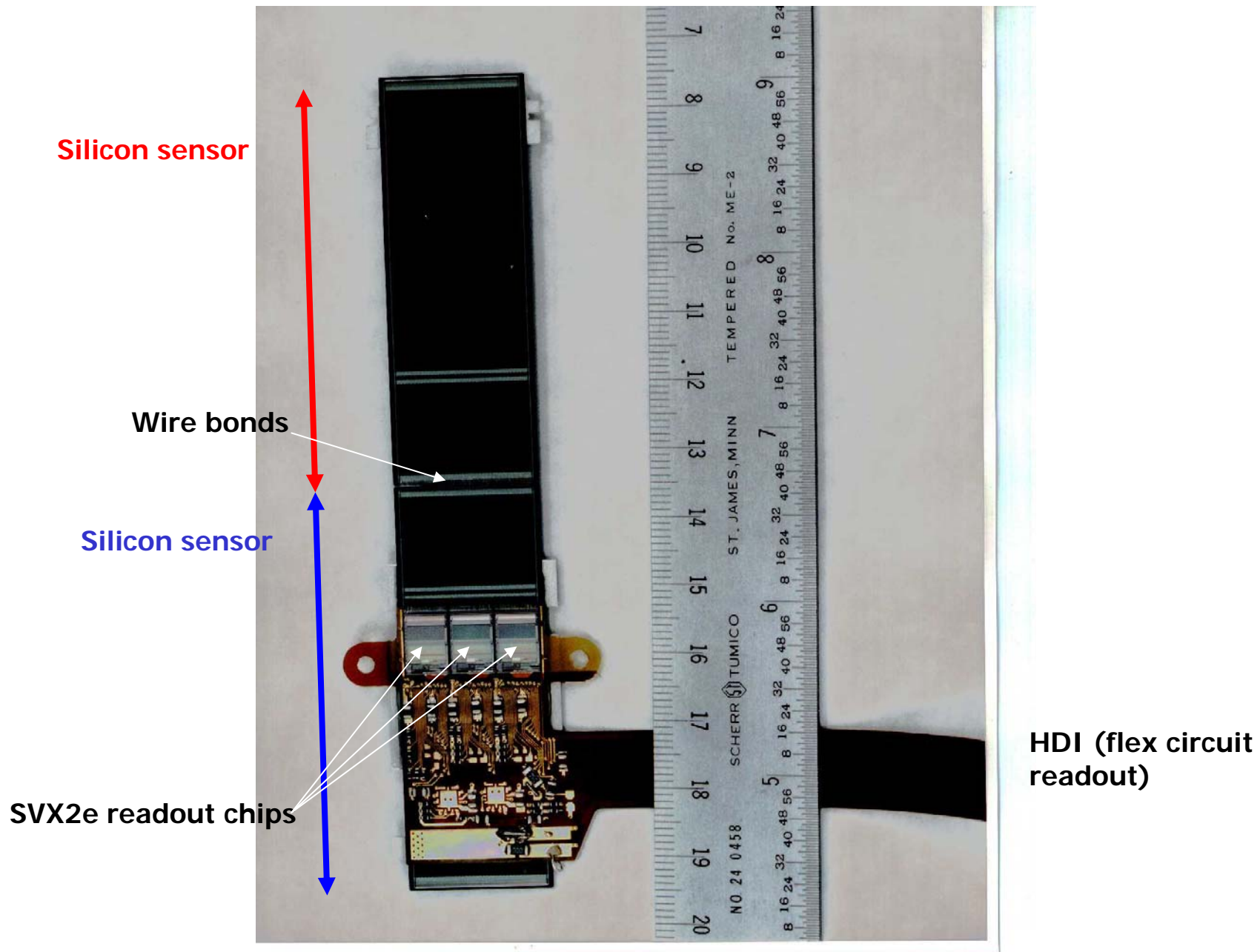


precise tracking close to primary collision point
 \Rightarrow silicon microstrip detectors

Decaying particles: examples

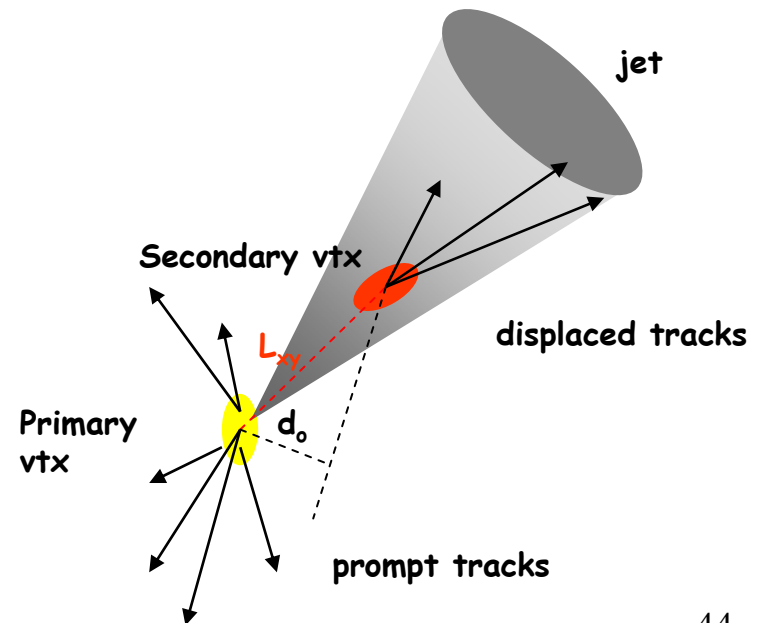
Particles	Lifetime	$c\tau$	Lifetime signature
W,Z,top	$<10^{-23}$ s	~ 0	Decay immediately
$\pi^0(\rightarrow\gamma\gamma)$	8×10^{-16} s	25 nm	Decay length undetectable
τ	2.9×10^{-13} s	87 μm	Inside beam pipe; hard even with SMT
$D^0/D^\pm/D_s$	$0.4\text{-}1.0 \times 10^{-12}$ s	150-350 μm	Inside beam pipe; possible w/ SMT
$B^0/B^\pm/B_s/b$ -baryon	$\sim 1.5 \times 10^{-12}$ s	450 μm	Inside beam pipe; possible w/ SMT
$K_s^0(\rightarrow\pi\pi)$	0.8×10^{-10} s	2.7 cm	decays in outer tracking chamber
K^\pm, π^\pm, μ^\pm	$>10^{-8}$ s	>3 m	reach cal without decaying





Lifetime signature for b-jets

- Long lifetime of B hadrons
 - $c\tau \sim 450\mu\text{m} + \text{boost}$
 - Travel $\sim 3\text{ mm}$ before decay with large charged track multiplicity
- Two ways of looking for lifetime tags
 - A reconstructed vertex, displaced from the primary vertex
 - Presence of track(s) with large impact parameter, d_0

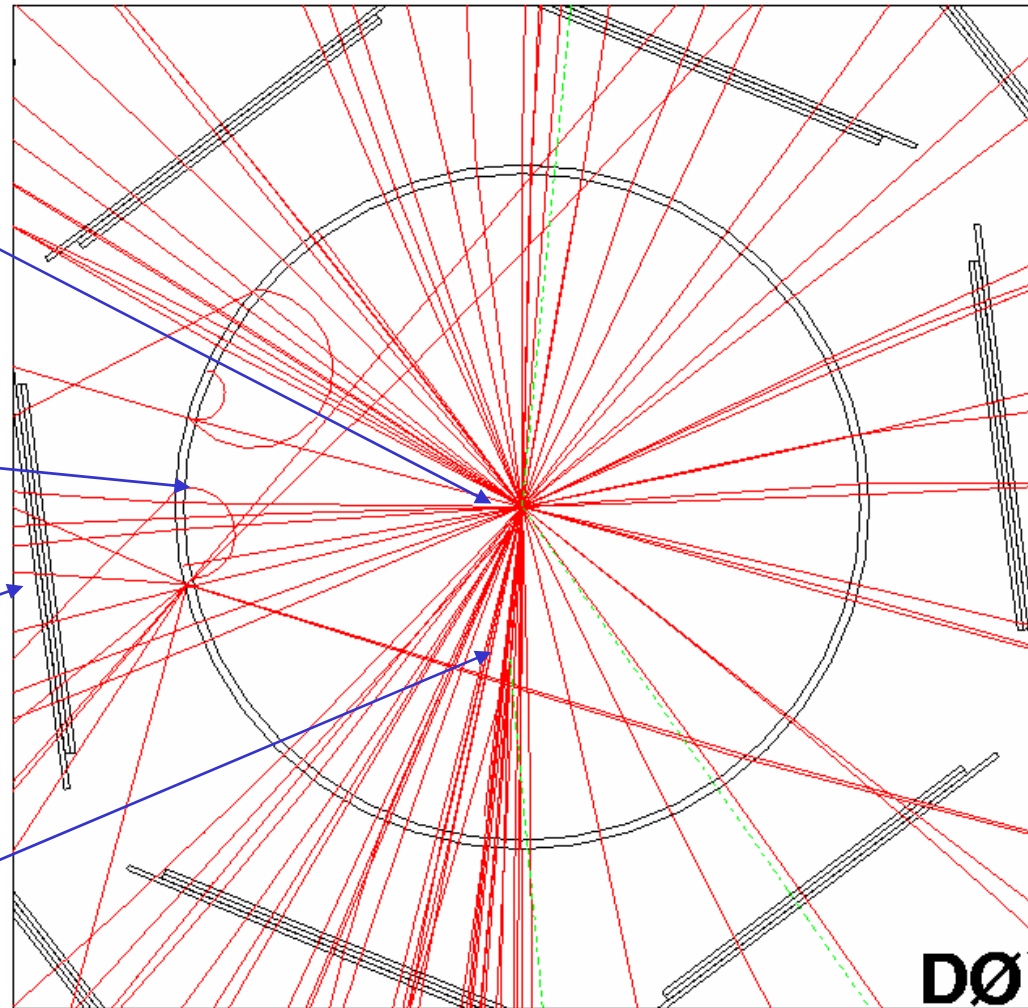


Interaction point
("primary vertex")

Beampipe

Silicon detector

B decay
("secondary
vertex")

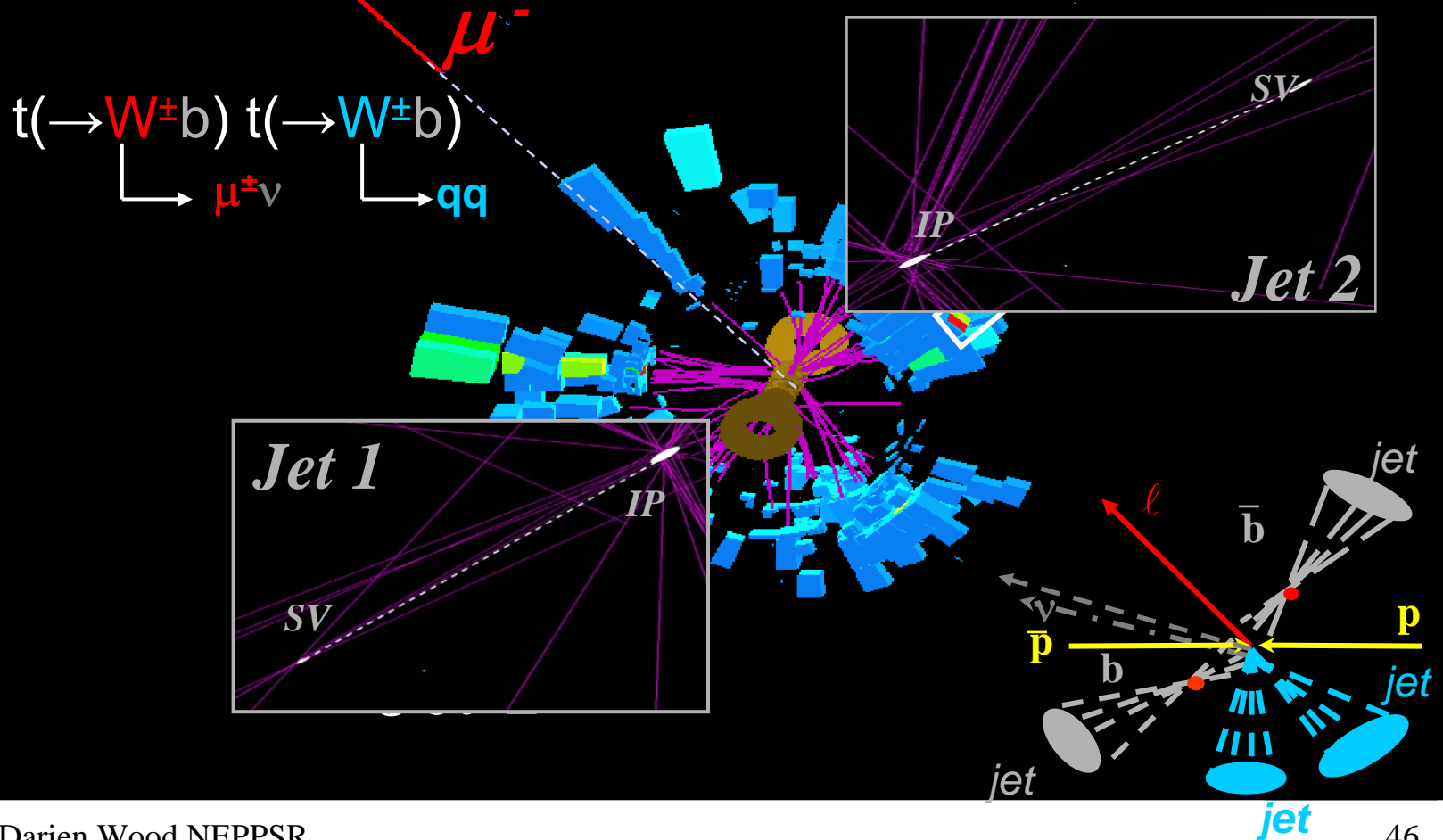


This green track clearly does not originate at the primary vertex

1 inch

t-tbar candidate event with b-tags

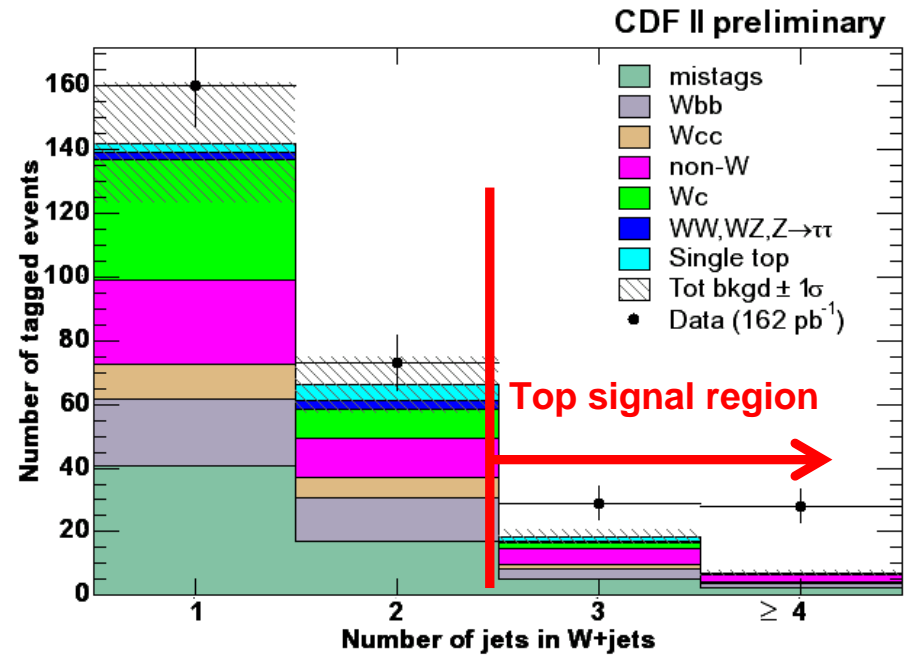
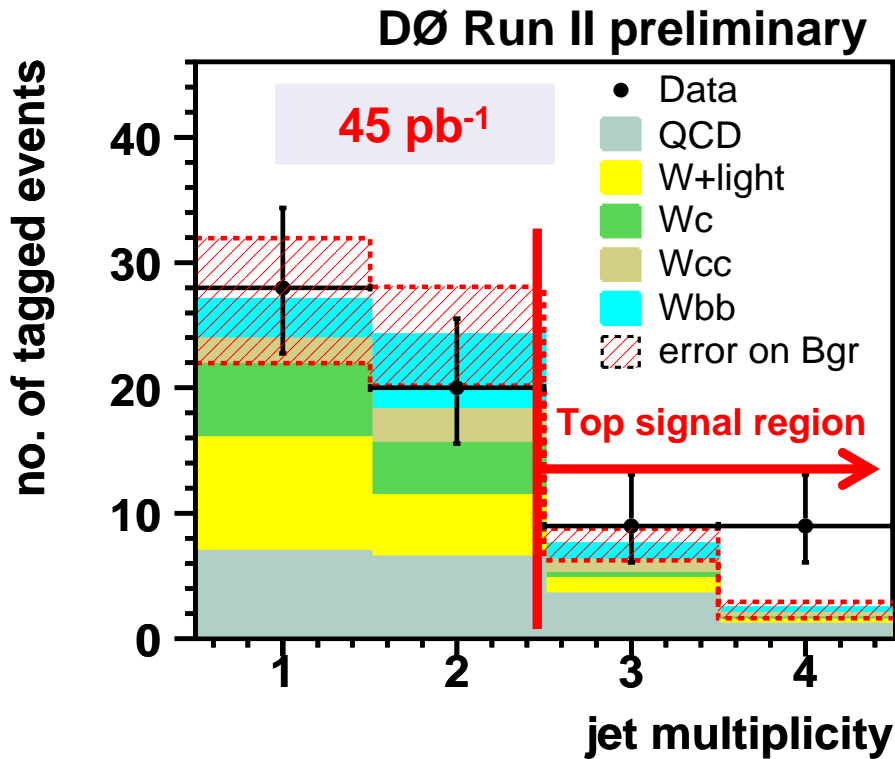
Other processes will also mimic this topology \Rightarrow backgrounds



Background to b-tagging

- soft lepton tag
 - fake leptons (in light-quark/gluon jets)
 - c-quark jets
 - K^\pm, π^\pm decays in flight (in light-quark/gluon jets)
 - chance overlap of light-quark/gluon jet and lepton
- lifetime tag
 - mismeasured tracks in light-quark/gluon jet
 - c-quark jets
 - K_s, Λ decays in light-quark/gluon jets

Top signal and background (Run II)



“lepton+jets” w/ b-tag

One of the best for mass determination

$$t\bar{t} \rightarrow l\nu b\bar{b} q\bar{q}$$

Triggering

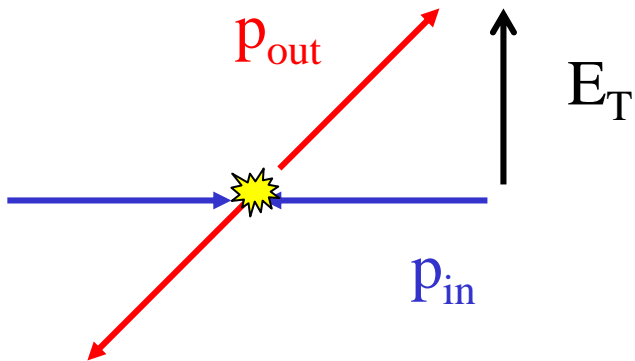
- A few relevant numbers:
 - Bunch collision rate: 2.5 MHz
 - Inelastic interaction rate at $L=10^{32}$: 4 MHz
 - Data size of each event: ~0.5 Megabytes
 - Rate at which data can be recorded: ~50 Hz
- Clearly, most of the events have to be discarded before they are recorded (record 1 crossing in 50,000)
- Events that are discarded are real physics (jet production, low p_T b-quarks, low-mass lepton pairs...)
 - Need to make choices about physics
 - Different from e^+e^- colliders, where generally all real interactions can be recorded
- **Triggers have a big, important job at hadron colliders**

Typical production rates for p-pbar at 2 TeV

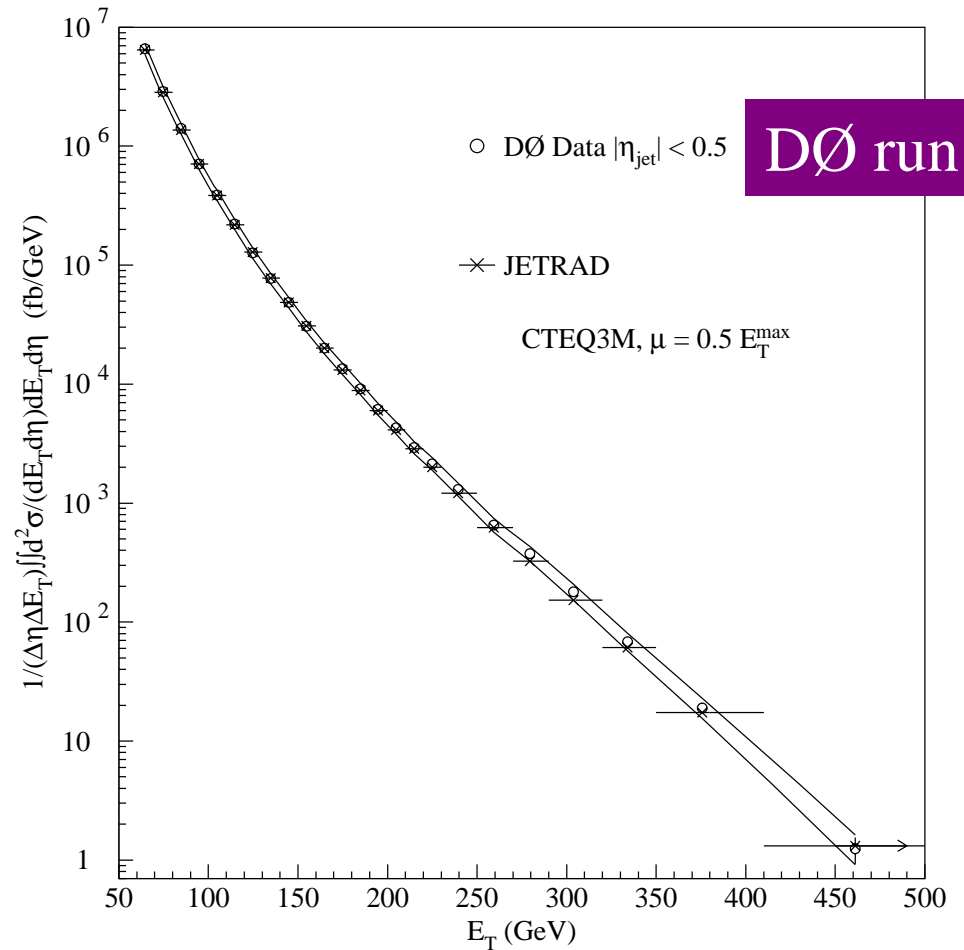
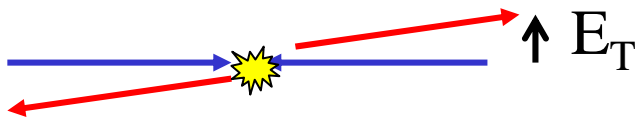
Final state	Cross section (pb)	Rate at $L=10^{32} \text{ cm}^{-2}\text{s}^{-1}$
“minimum bias”	4×10^{10}	4 MHz
2 jets	4×10^6	400 Hz
4 jets	1.6×10^5	16 Hz
6 jets	6000	0.6 Hz
W	30000	3 Hz
Z	9000	0.9 Hz
WZ	3.5	3.5×10^{-4} Hz (1.3/hour)
t tbar	7.5	7.5×10^{-4} Hz (3/hour)

Spectrum of jet transverse energy

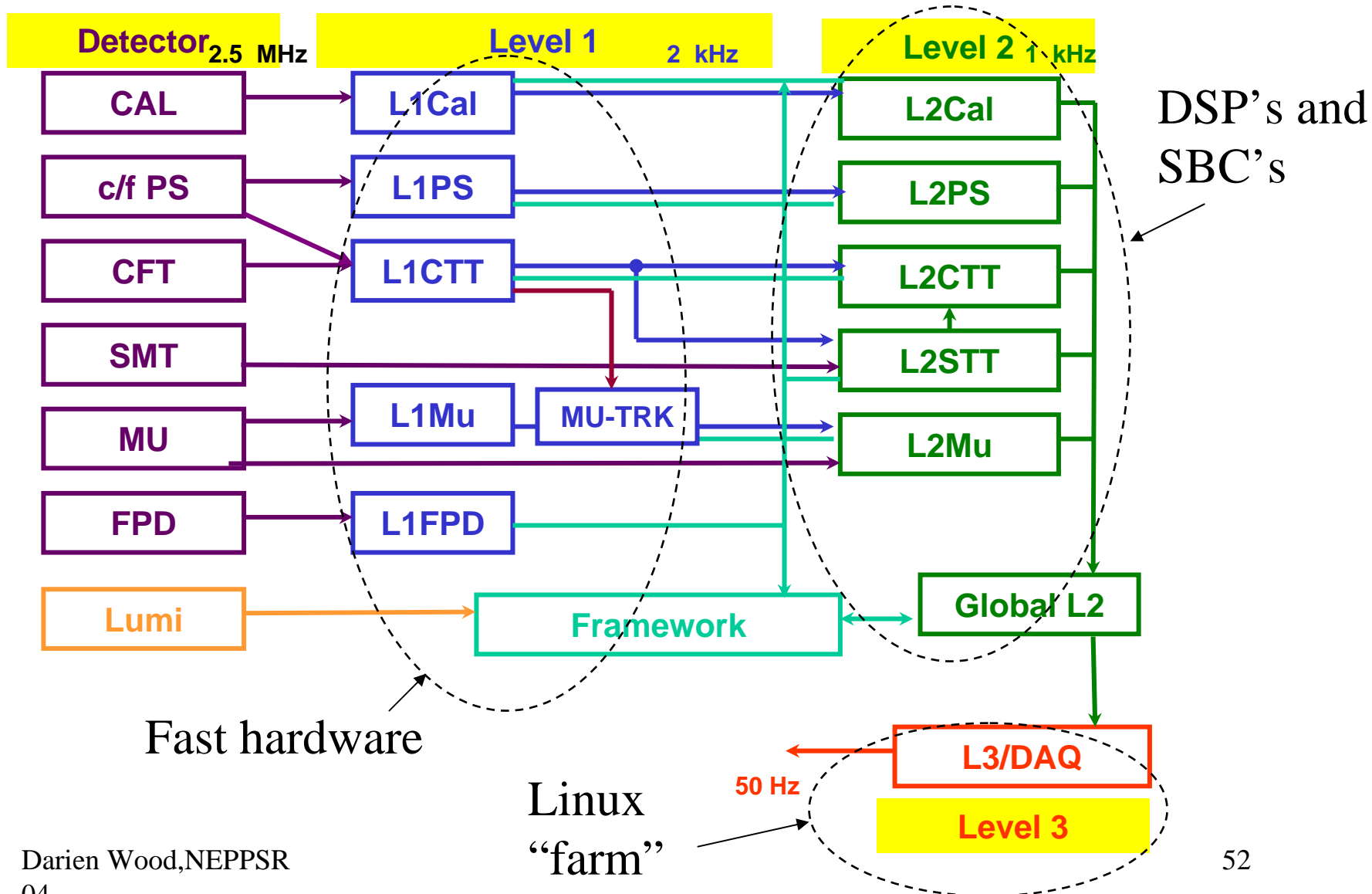
Hard collision:



soft collision:



The DØ Trigger System (current)



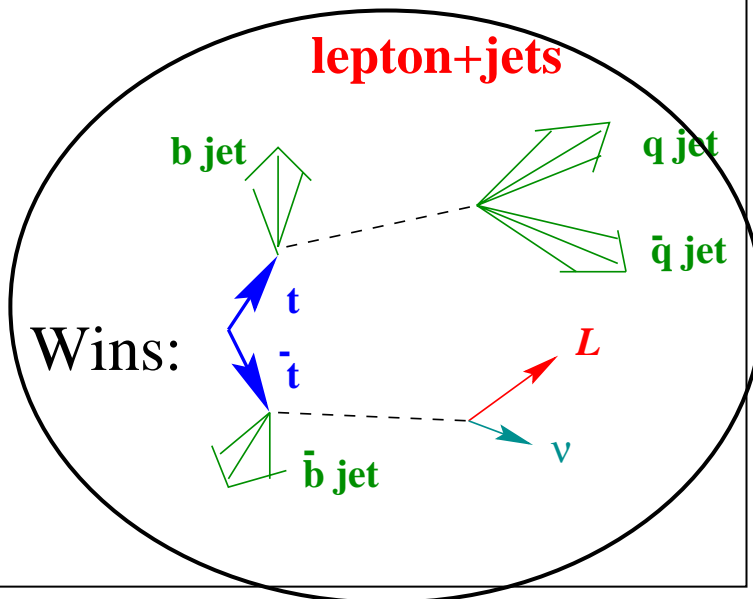
Example of a trigger condition

- Name: MUJ1_JT15HA_TK10 (one of 378 triggers)
- Purpose: trigger on top \rightarrow muon+jets (among others)
- Level 1
 - 2-layer coincidence of muon scintillator and of muon wire chambers
 - At least two “towers” (0.2 eta X 0.2 phi) with E_T above 3 GeV
 - Rate: 70 Hz
- Level 2:
 - Fast reconstruction muon track (no momentum cut)
 - Fast jet cluster with E_T above 8 GeV
 - Rate: 55 Hz
- Level 3:
 - Reconstructed jet above 15 GeV, and sum of missing E_T and scalar $E_T > 50$ GeV
 - Charged track with $p_T > 10$ GeV
 - Rate: ~ 1 Hz

Top quark mass measurement

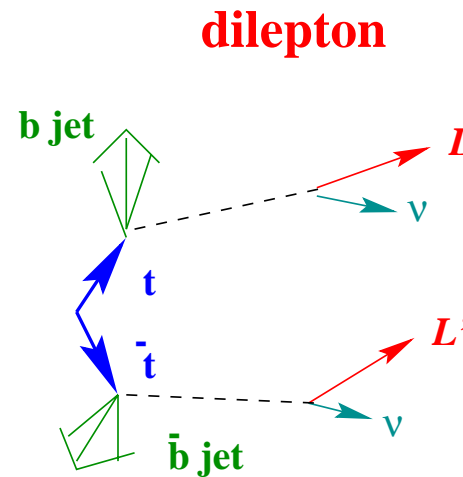
- advantages

- 2nd largest branching ratio
- only one neutrino



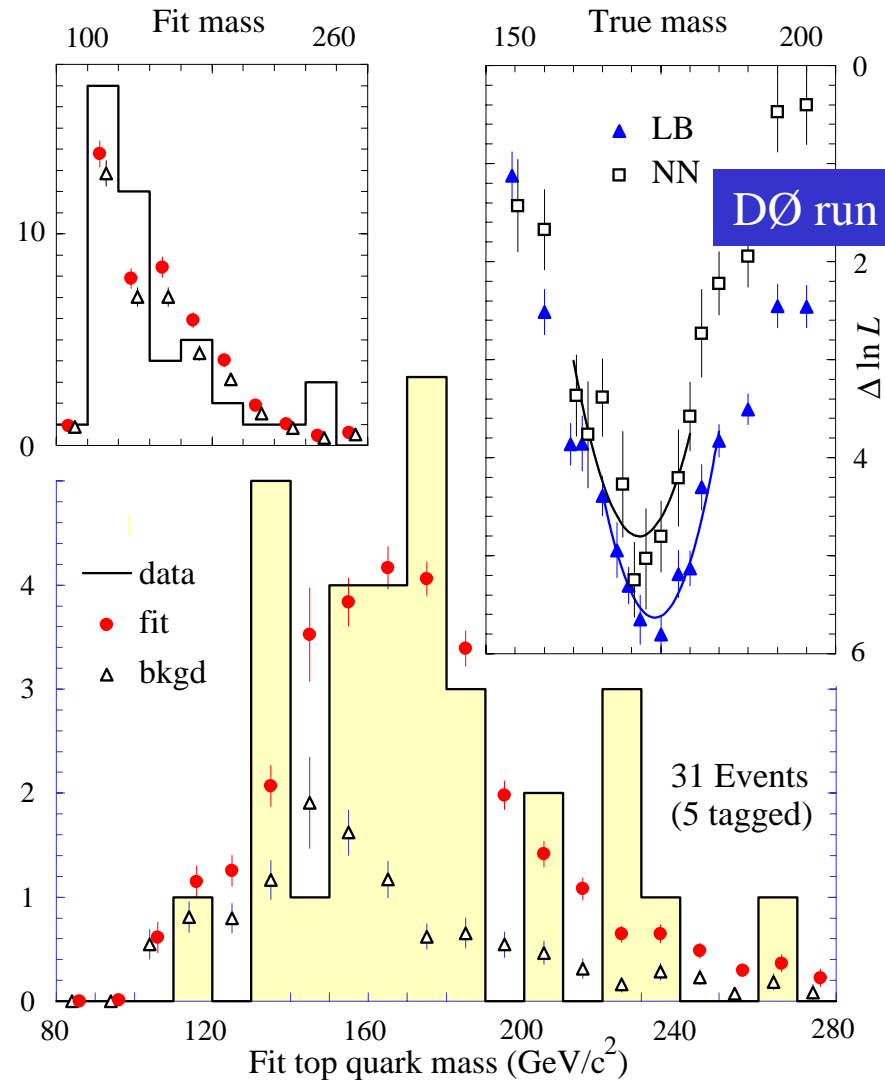
- advantages

- low background
- better energy resolution for leptons than for jets

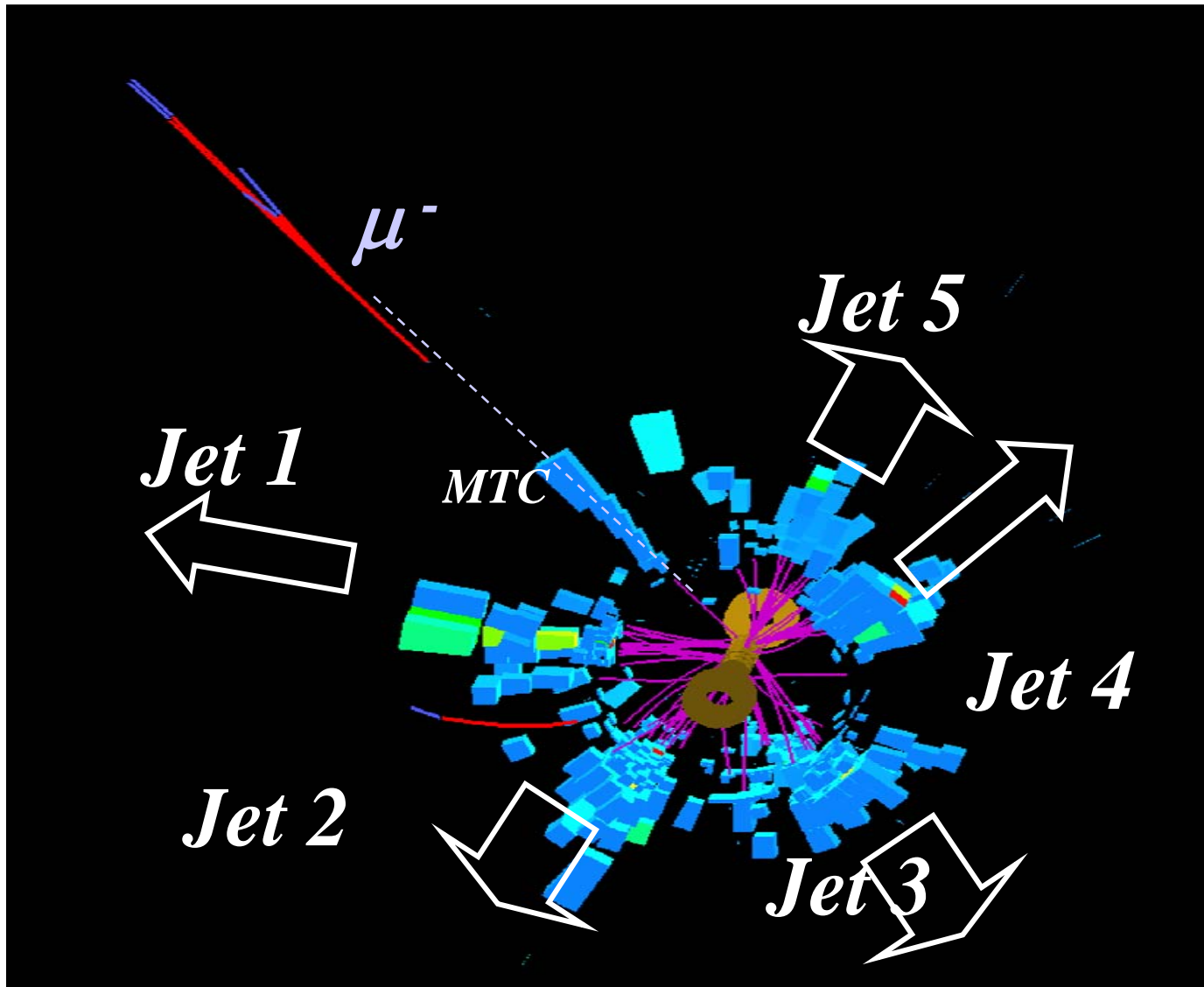


Technique #1: one mass per event

- Lepton+4jets channel: all final state variables (18) are measured except for $p_z(\nu)$
- Three additional constraints:
 - $m(qq) = m_W$
 - $m(l\nu) = m_W$
 - $m(bqq) = m(b\nu)$
- Twice over-constrained fit: (17+3-18)
- For each event, select jet permutation with best fit (lowest chisquared)
- Make distribution of masses, and fit to $\text{signal}(m_t) + \text{bkg}$



Darien Wood, NEPPR $m_t = 173.3 \pm 5.6(\text{stat}) \pm 5.5(\text{sys}) \text{ GeV}$



DØ top to μ +jets Candidate Event

Technique #2: optimized matrix element weighting

Likelihood method using most available information

Measured \rightarrow

$$P(x; \alpha) = \frac{1}{\sigma} \int d^n \sigma(y; \alpha) dq_1 dq_2 f(q_1) f(q_2) W(x, y)$$

to be estimated \rightarrow

resolutions, reconstruction effects \rightarrow

Matrix Element \rightarrow

PDF's \rightarrow

LO ME used, 4 jets required exclusively, additional cut on background probability (to improve purity) \rightarrow **22 events**

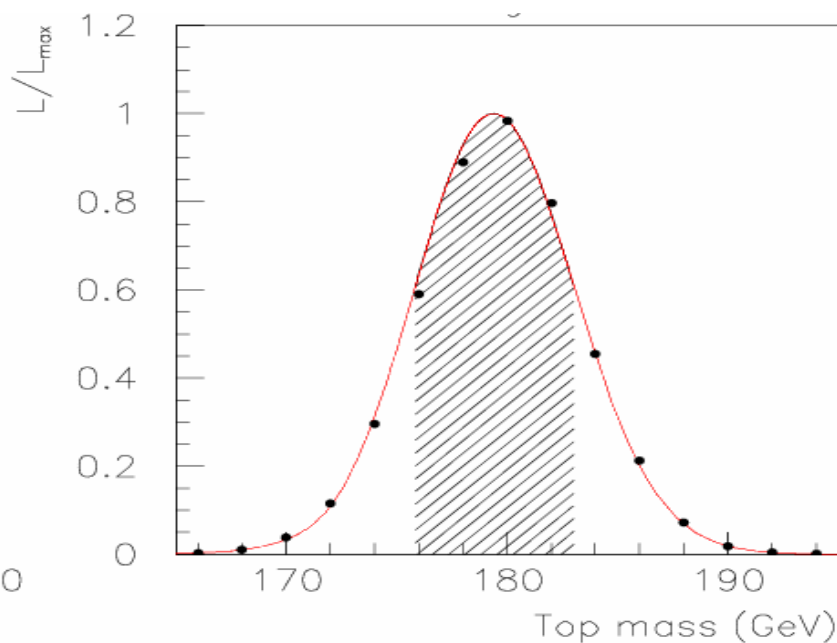
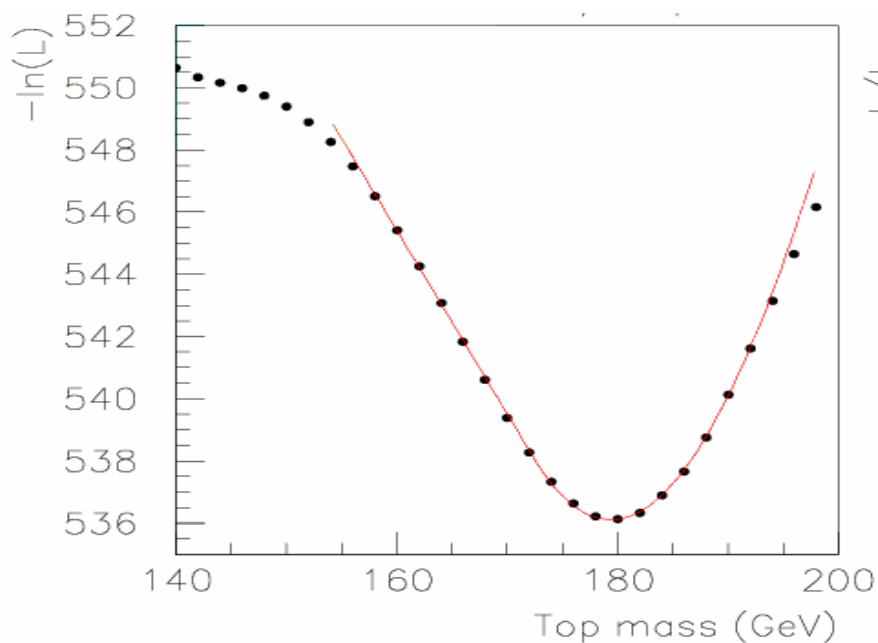
$$-\ln L(\alpha) = -\sum_{i=1}^N \left\{ \ln \left[c_1 P_{\bar{t}t}(x_i; \alpha) + c_2 P_{\text{bkg}}(x_i) \right] \right\}$$

$$+ N \int A(x) \left[c_1 P_{\bar{t}t}(x; \alpha) + c_2 P_{\text{bkg}}(x) \right] dx$$

Acceptance \rightarrow

Likelihood definition:
estimate signal and
background fractions
and m_{top}

Results of matrix element weighting



$$m_{\text{top}} = 180.1 \pm 3.6 (\text{stat}) \pm 3.9 (\text{sys}) \text{ GeV}/c^2$$

large improvement on the statistical uncertainty ($\sim 2.4 \times$ stats)

NEW World average:

$$m_{\text{top}} = 178.0 \pm 4.3 \text{ GeV}/c^2$$

Run 2 mass analyses under way.
With 2 fb^{-1} data: $\delta m_{\text{top}} \approx 1\text{-}2 \text{ GeV}$

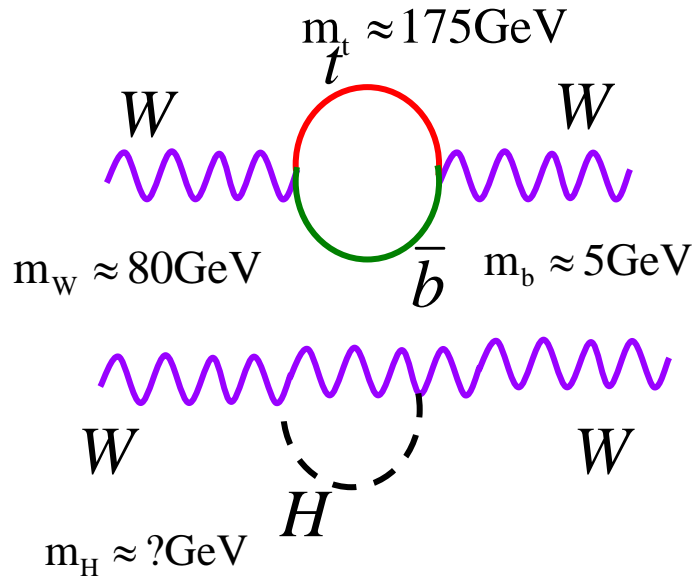
Darien Wood, NEPPSR

04

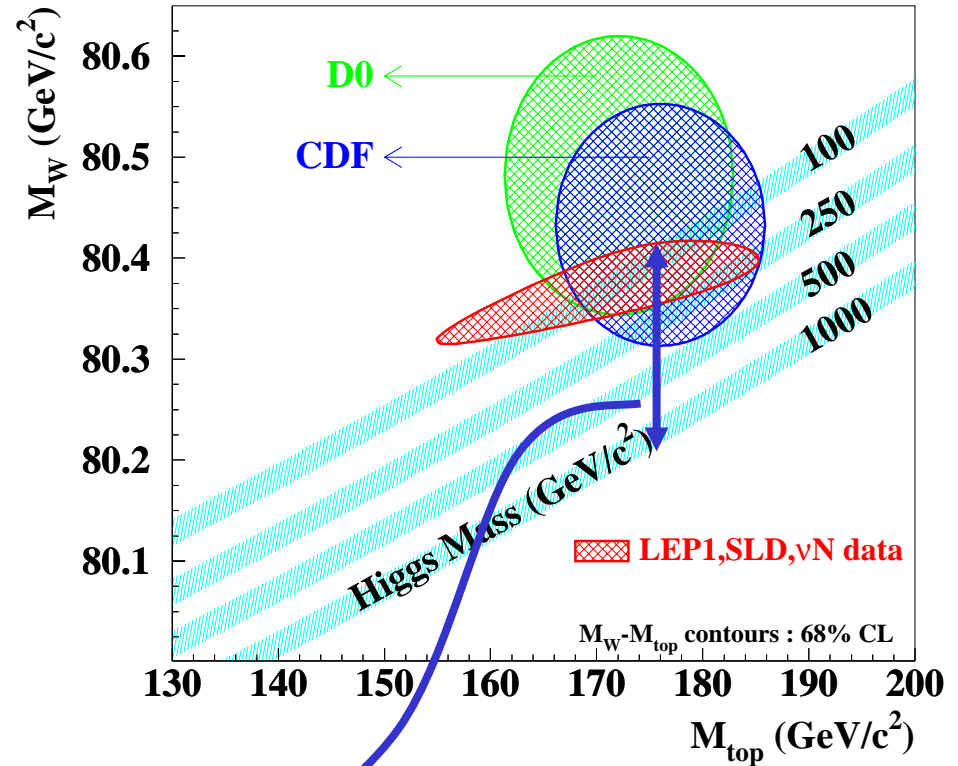
Largest systematic error: Jet energy calibration: 3.3 GeV

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Virtual Effects on W mass



- One piece of evidence missing: the Higgs particle responsible for the mass of all known particles
- Finding the Higgs (or not): verify a prediction or declare clear evidence for new physics

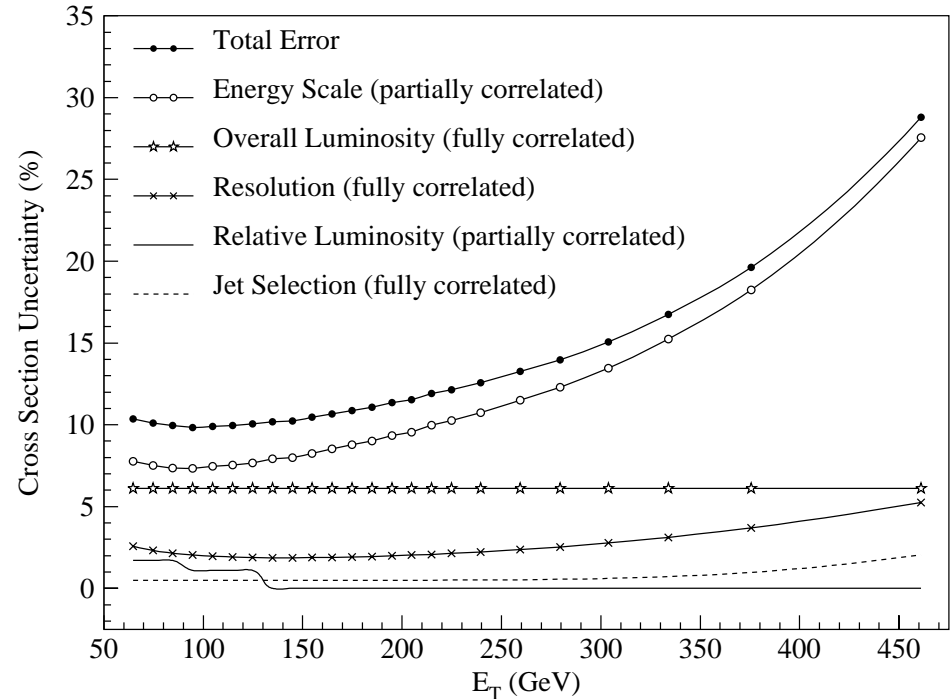
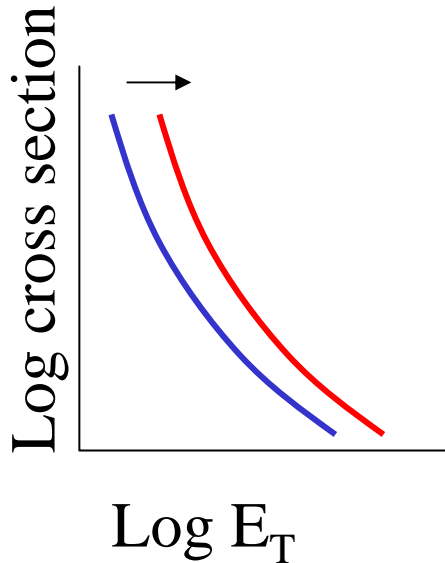


Mass shift from virtual Higgs effects (?)

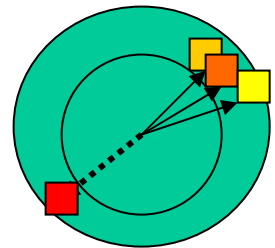
Interesting features of top mass analysis

- small number of events
 - fitting techniques refined to make maximum use of limited statistics
- substantial background
 - fitting has to be robust in the presence of background
- complicated event topology
 - kinematic fitting
 - consideration of multiple permutations of object assignments

Jet energy calibration



A small shift in energy scale calibration gives a large shift in cross section



Jet energy calibrated in situ with photon+jet events